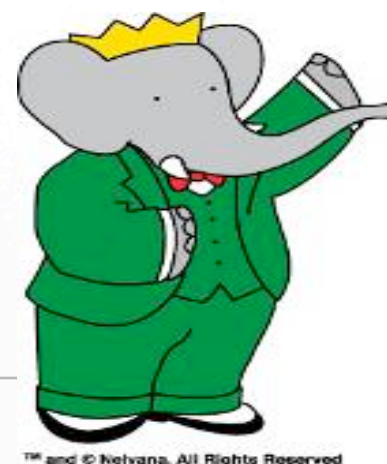




CKM results from the B-Factories



TM and © Nelvana, All Rights Reserved

Youngjoon Kwon

Yonsei Univ. (on behalf of the Belle collab.)

The 6th KEK Topical Conference
“Frontiers in Particle Physics & Cosmology”



Overview

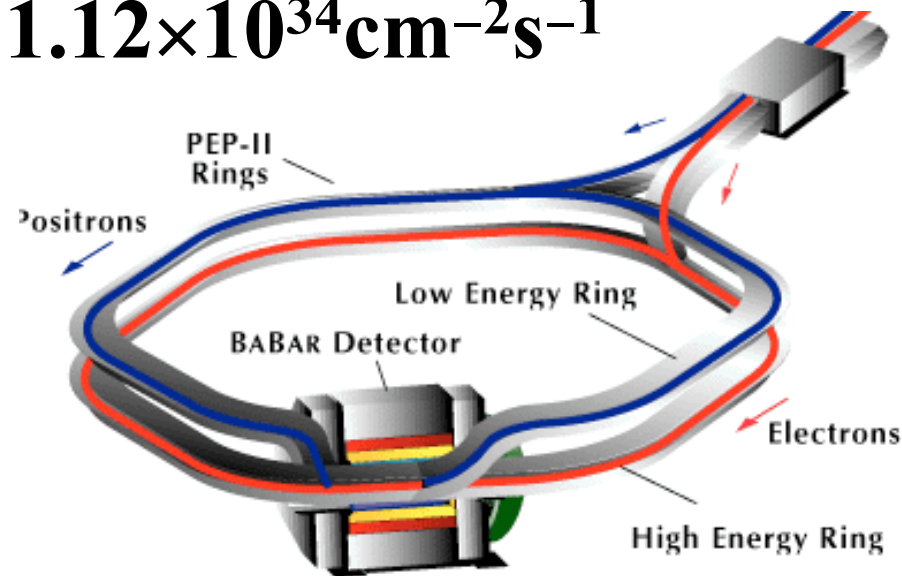
- Intro. (for non-experts)
- Experimental Results
 - The angles
 - The sides
- Summary & Outlook



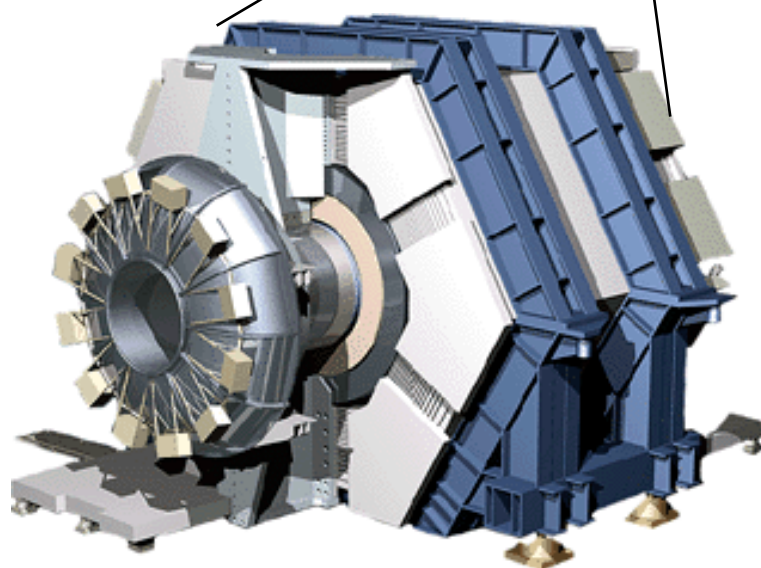
Two asymmetric-energy B factories

PEP-II at SLAC

9GeV (e^-) \times 3.1GeV (e^+)
peak luminosity:
 $1.12 \times 10^{34} \text{cm}^{-2} \text{s}^{-1}$

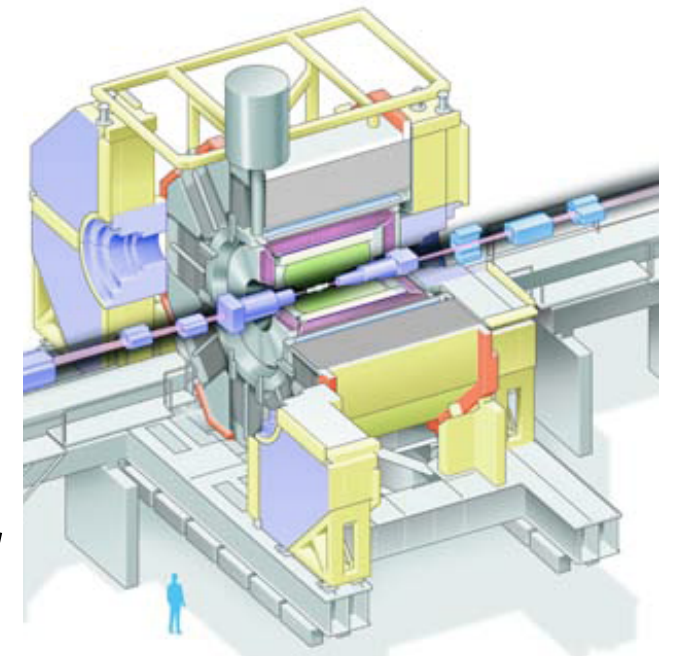


BaBar

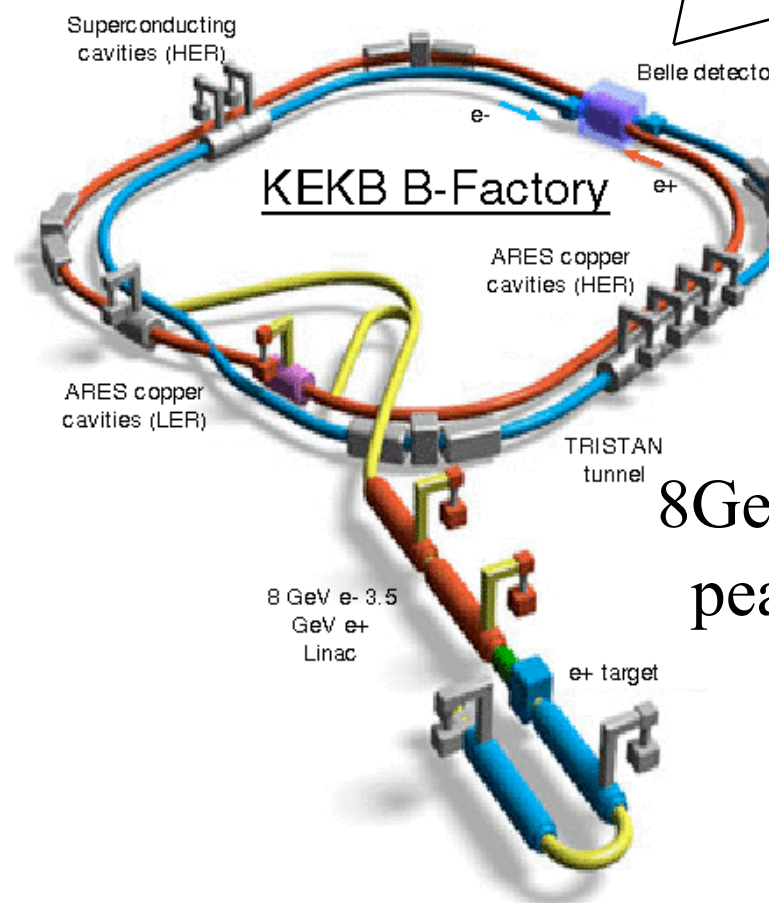


11 nations,
80 institutes,
623 persons

13 countries,
57 institutes,
~400 collaborators



Belle

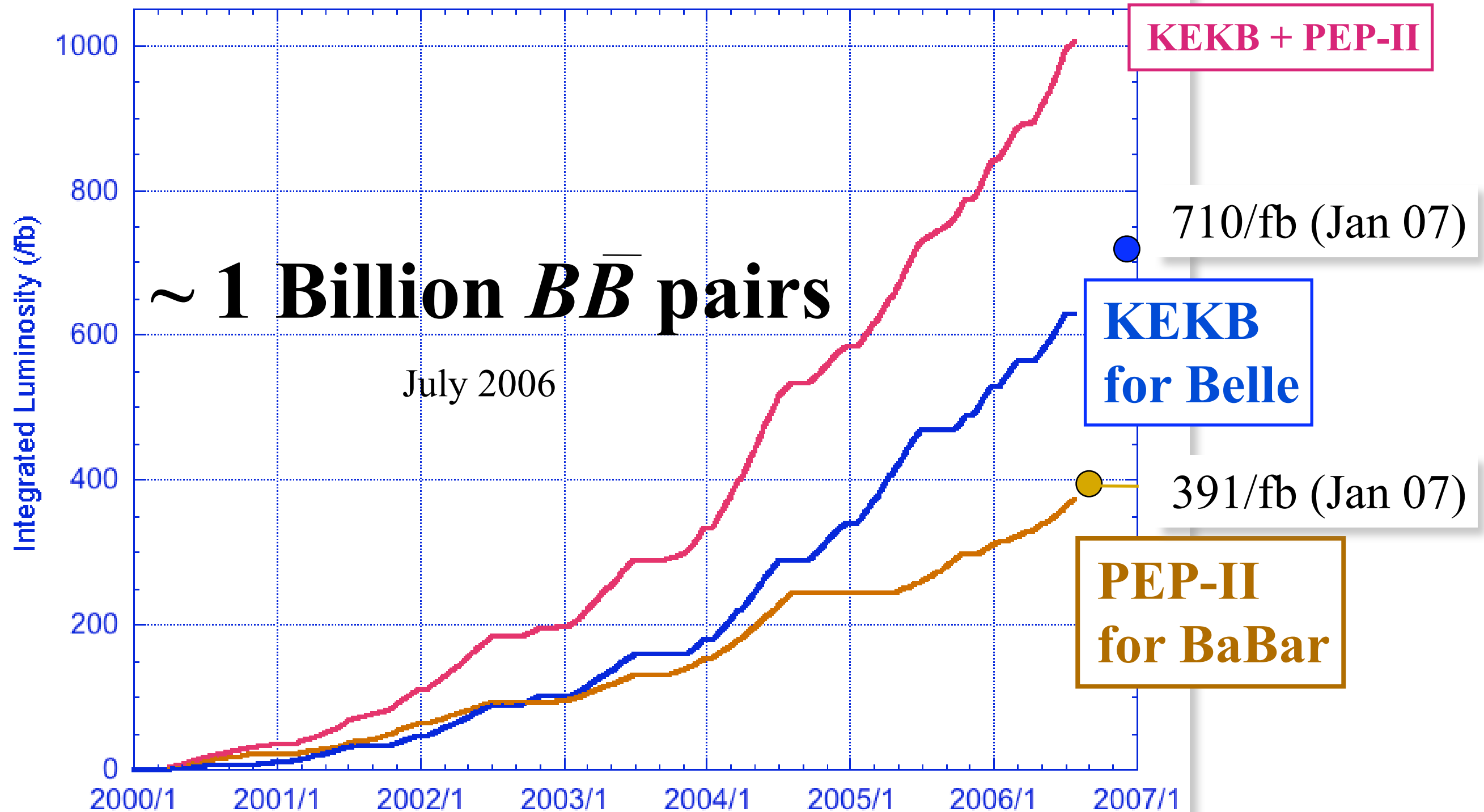


KEKB at KEK

8GeV (e^-) \times 3.5GeV (e^+)
peak luminosity:
 $1.71 \times 10^{34} \text{cm}^{-2} \text{s}^{-1}$
world record

Integrated Luminosity

World Integrated Luminosity (KEKB+PEP-II)



Motivations

- Big bang cosmology suggests a **balanced production of matter & antimatter**, but our current universe is dominated by matter -- at least, the visible part.
- **A. Sakharov's 3 conditions** for matter dominance
 - baryon number non-conservation
 - C and CP violation
 - not in thermal equilibrium
- Lorenz-invariant quantum field theory demands **invariance under CPT**, but nothing more.
- In the Standard Model, we actually have a mechanism for CP violation:
⇒ the **KM mechanism**.

Kobayashi-Maskawa (KM) ansatz



“CPV is due to an irreducible phase in the quark mixing matrix in 3 generations”

Journal of Theoretical Physics, Vol. 49, No. 2, February 1973

***CP*-Violation in the Renormalizable Theory of Weak Interaction**

Makoto KOBAYASHI and Toshihide MASKAWA

Department of Physics, Kyoto University, Kyoto

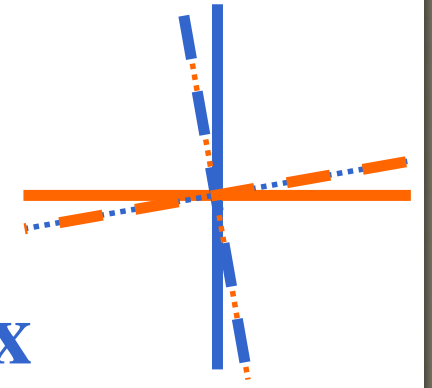
(Received September 1, 1972)

In a framework of the renormalizable theory of weak interaction, problems of *CP*-violation are studied. It is concluded that no realistic models of *CP*-violation exist in the quartet scheme without introducing any other new fields. Some possible models of *CP*-violation are also discussed.

When we apply the renormalizable theory of weak interaction¹⁾ to the hadron system, we have some limitations on the hadron model. It is well known that there exists, in the case of the triplet model, a difficulty of the strangeness changing neutral current and that the quartet model is free from this difficulty. Fur-

First 3rd-gen.
particle (τ)
seen in 1975

Flavor mixing and CKM matrix



- For quarks,
 - weak interaction eigenstates \neq mass eigenstates
 - mixing of quark flavors through a **unitary matrix**

$$\begin{pmatrix} d' \\ s' \\ b' \end{pmatrix} = (V_{\text{CKM}}) \begin{pmatrix} d \\ s \\ b \end{pmatrix} = \begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix} \begin{pmatrix} d \\ s \\ b \end{pmatrix}$$

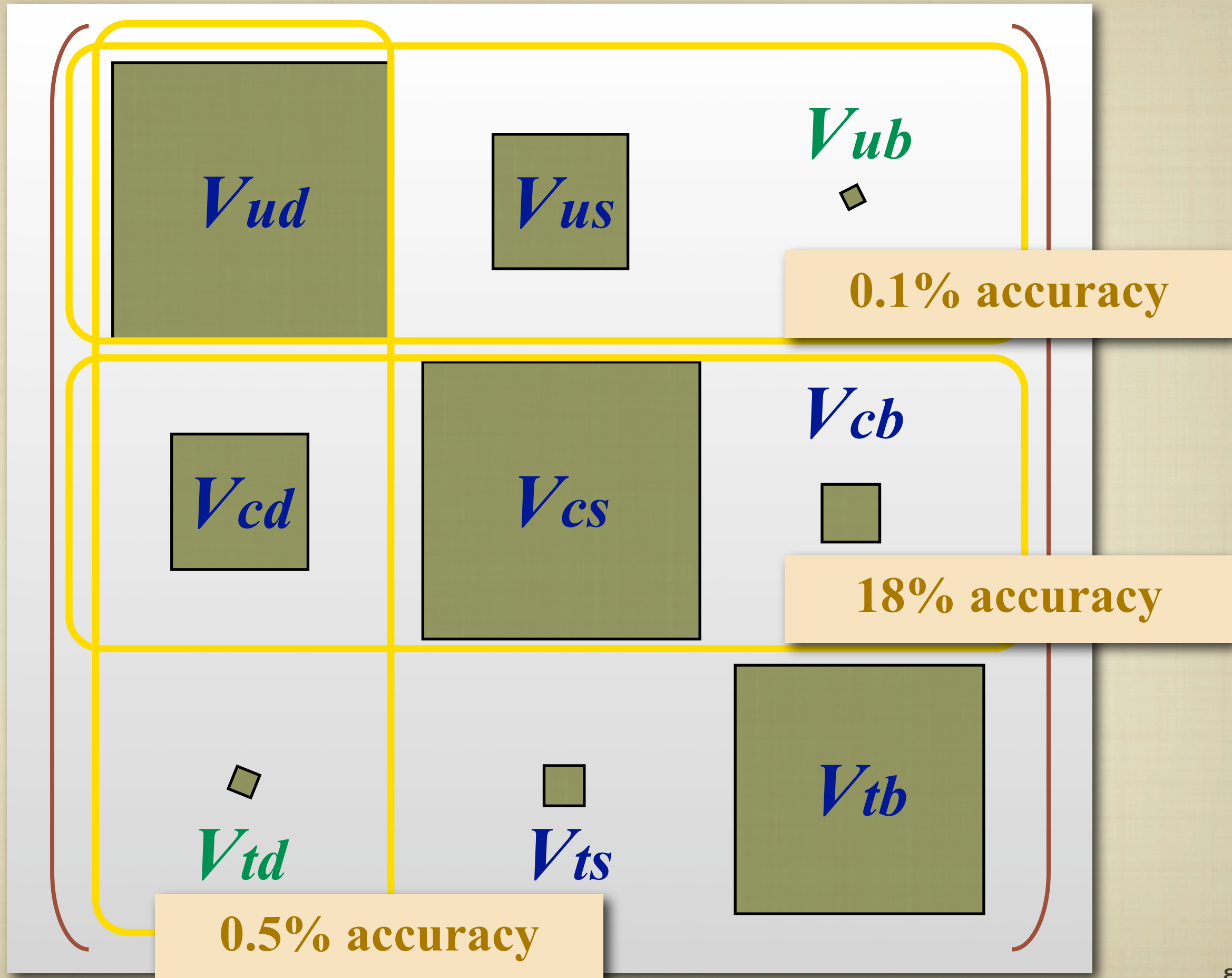
**Wolfenstein
parametrization**

$$V_{\text{CKM}} \approx \begin{pmatrix} 1 - \lambda^2/2 & \lambda & \frac{A\lambda^3(\rho - i\eta)}{A\lambda^2} \\ -\lambda & 1 - \lambda^2/2 & A\lambda^2 \\ \frac{A\lambda^3(1 - \rho - i\eta)}{A\lambda^2} & -A\lambda^2 & 1 \end{pmatrix}$$

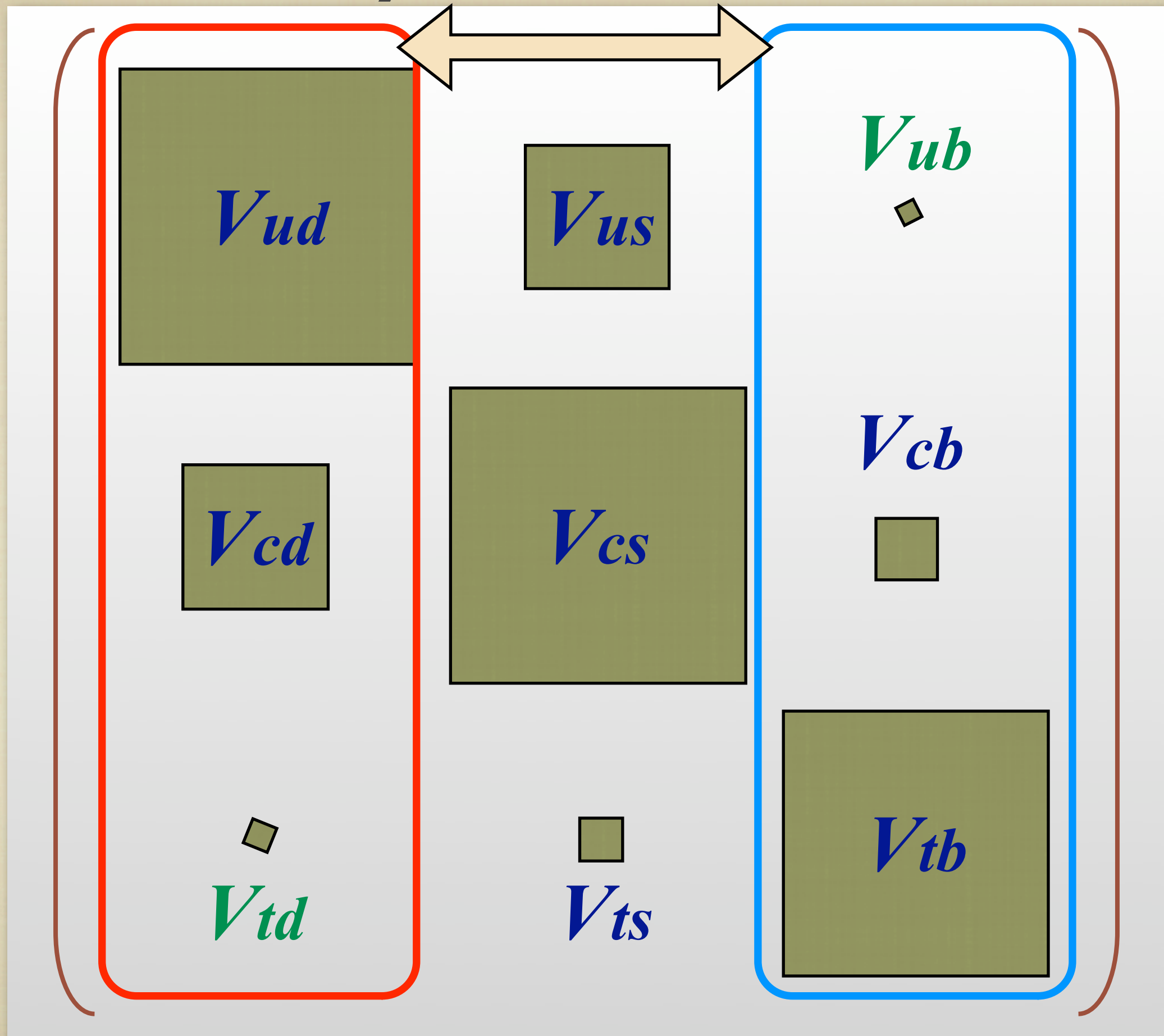
$$|\lambda| \approx O(0.1)$$

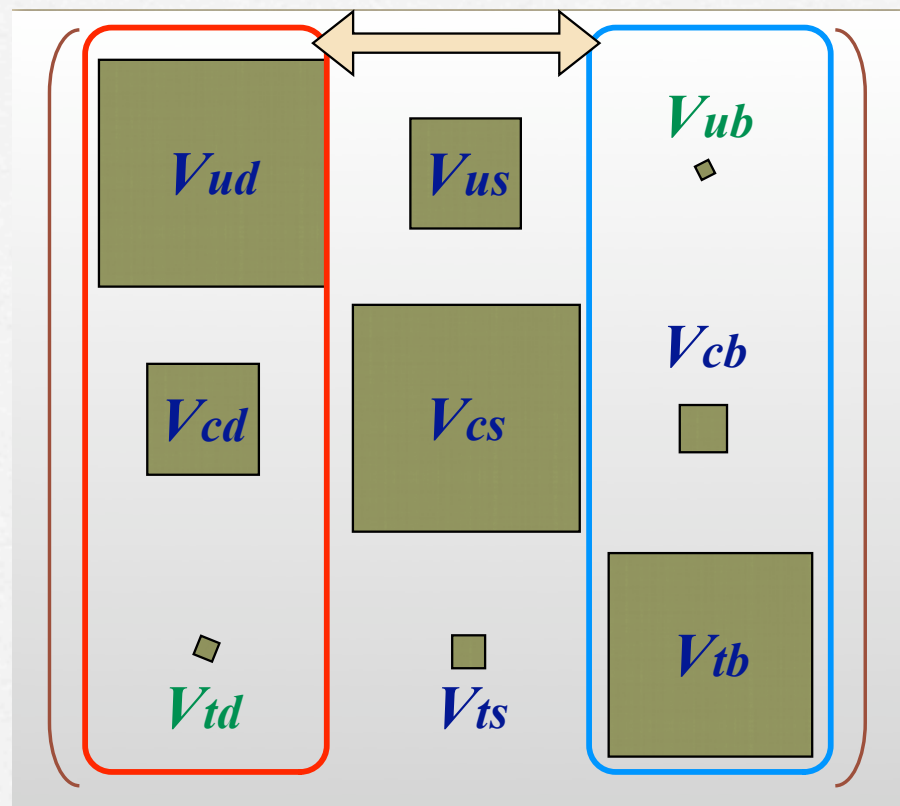
3 real parameters (λ, A, ρ) and 1 phase (η)

Test of Unitarity



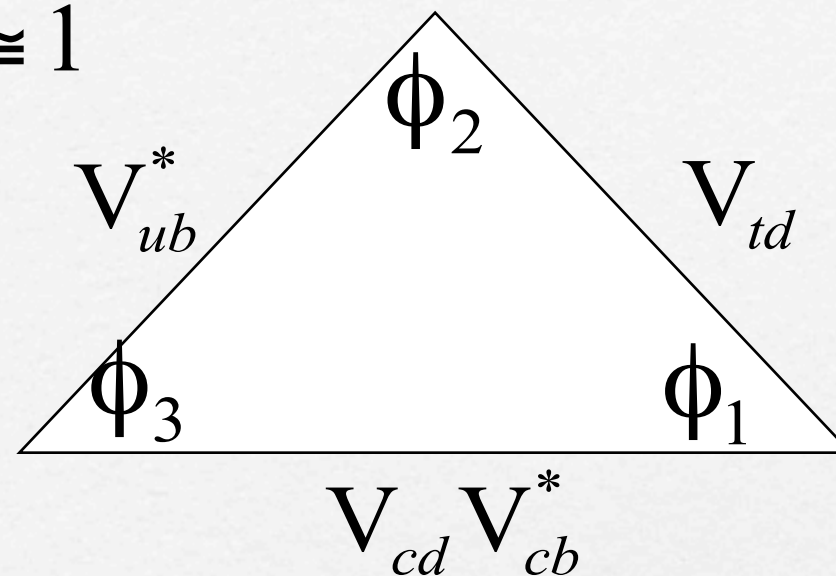
Test of Unitarity





$$V_{ud} V_{ub}^* + V_{cd} V_{cb}^* + V_{td} V_{tb}^* = 0$$

$$V_{ud} \cong V_{tb} \cong 1$$



Unitarity triangle angles

BABAR: β α γ

BELLE: ϕ_1 ϕ_2 ϕ_3

This talk: 易 難 魔

How to measure?

V_{ud}

V_{us}

V_{ub}

V_{cd}

V_{cs}

V_{cb}

$V = |V| \exp(i\phi)$

- $|V|$ from semi-leptonic decay rates
- ϕ from CP asymmetries

just overly simplified guidelines

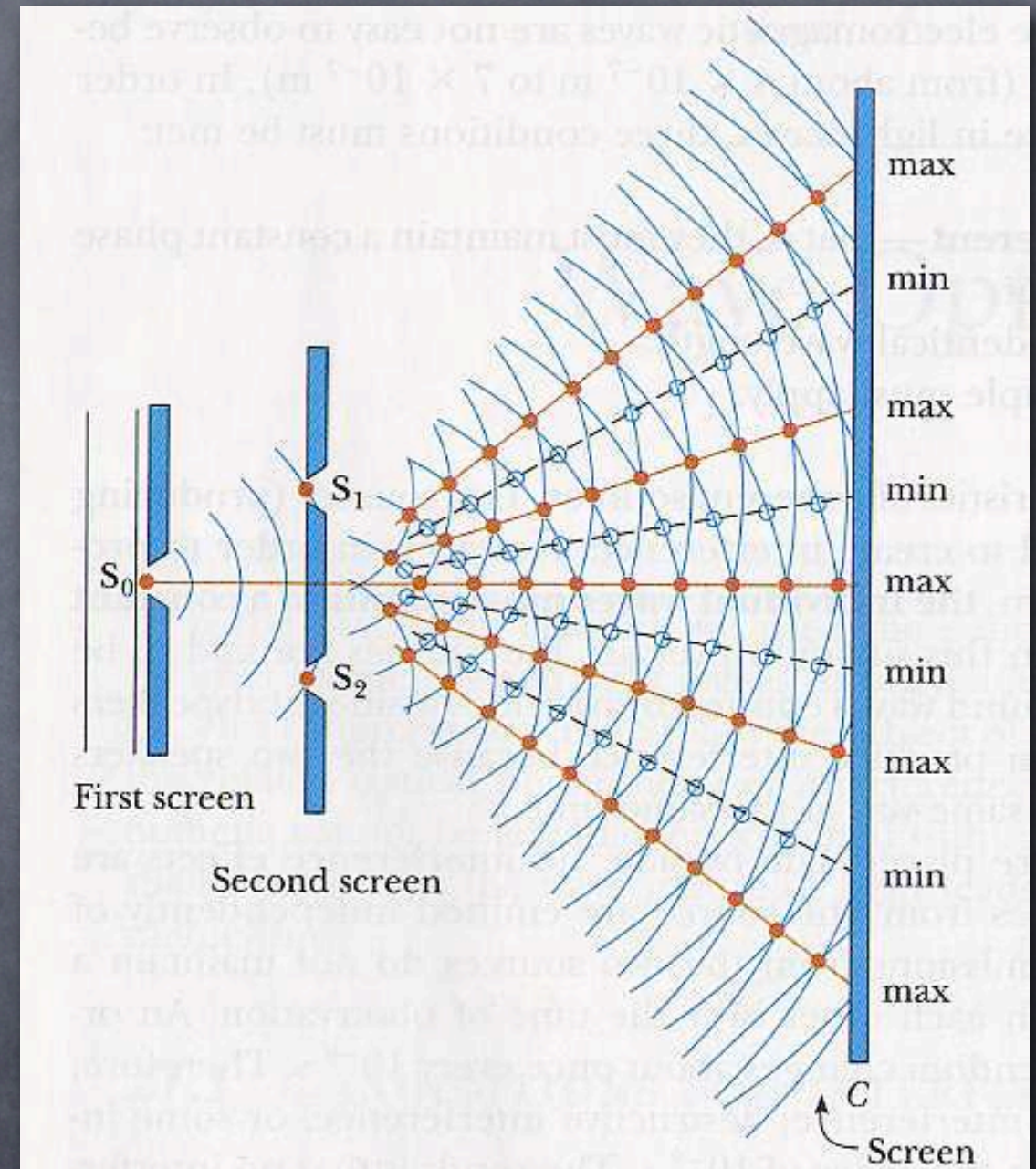
V_{td}

V_{ts}

V_{tb}

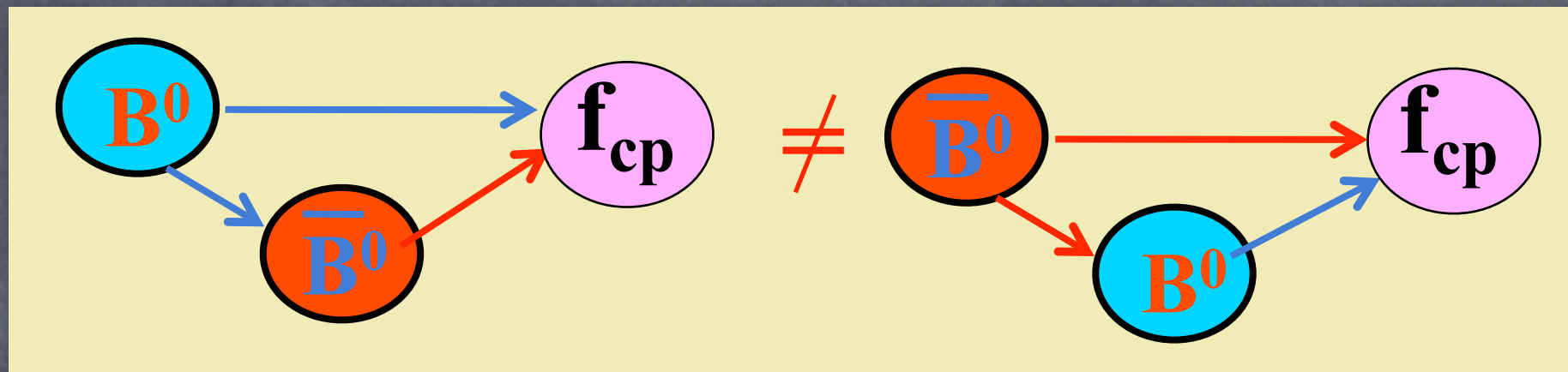
CP violation in the SM

- Interference of 2 amplitudes
 - case (1) : mixing-induced
 - case (2) : b/w tree & penguin

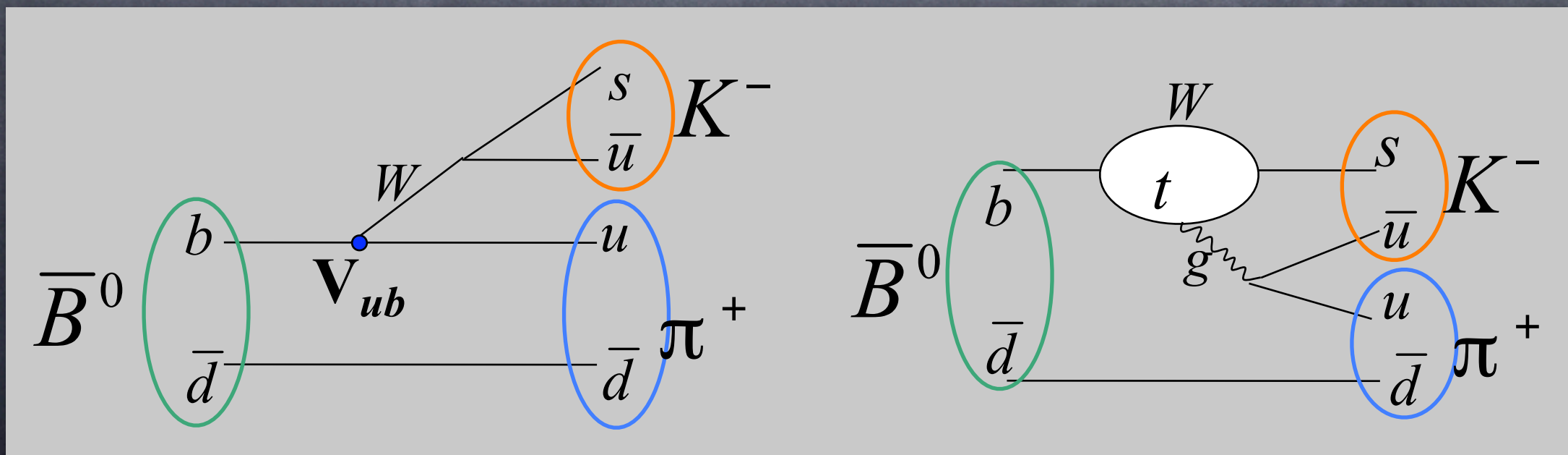


CP violation in the SM

- mixing-induced



- interference of different decay amplitudes

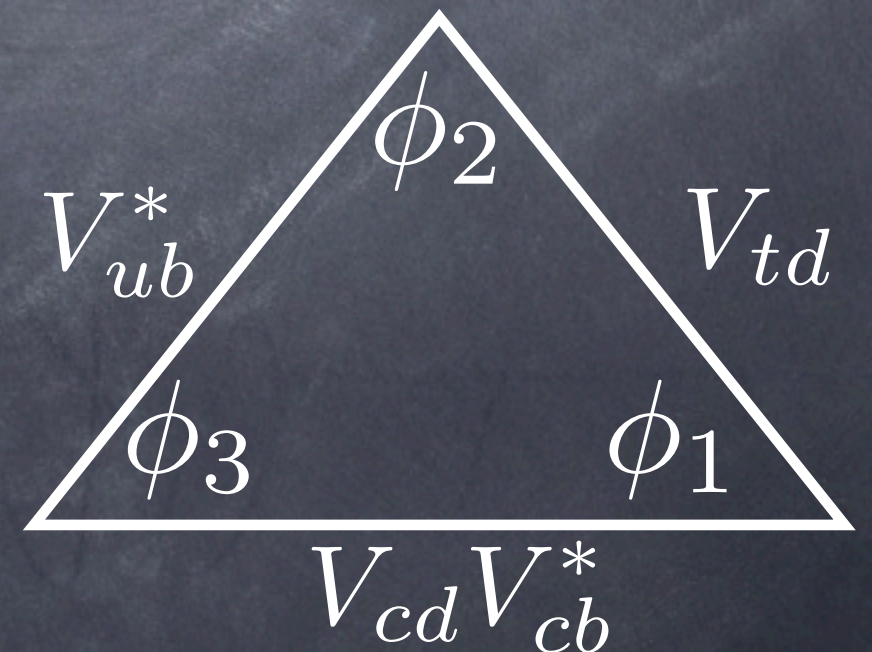


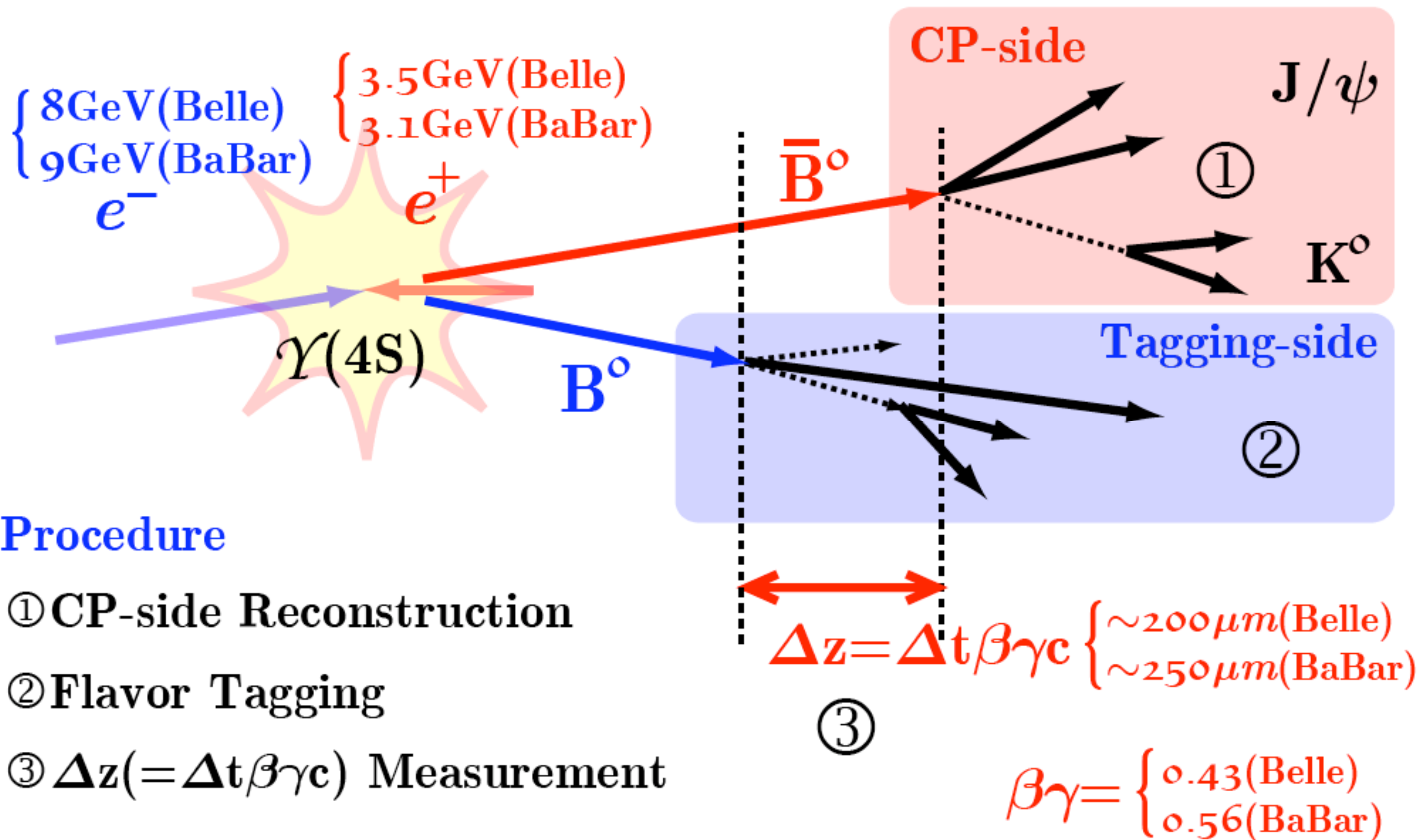


Experimental Results

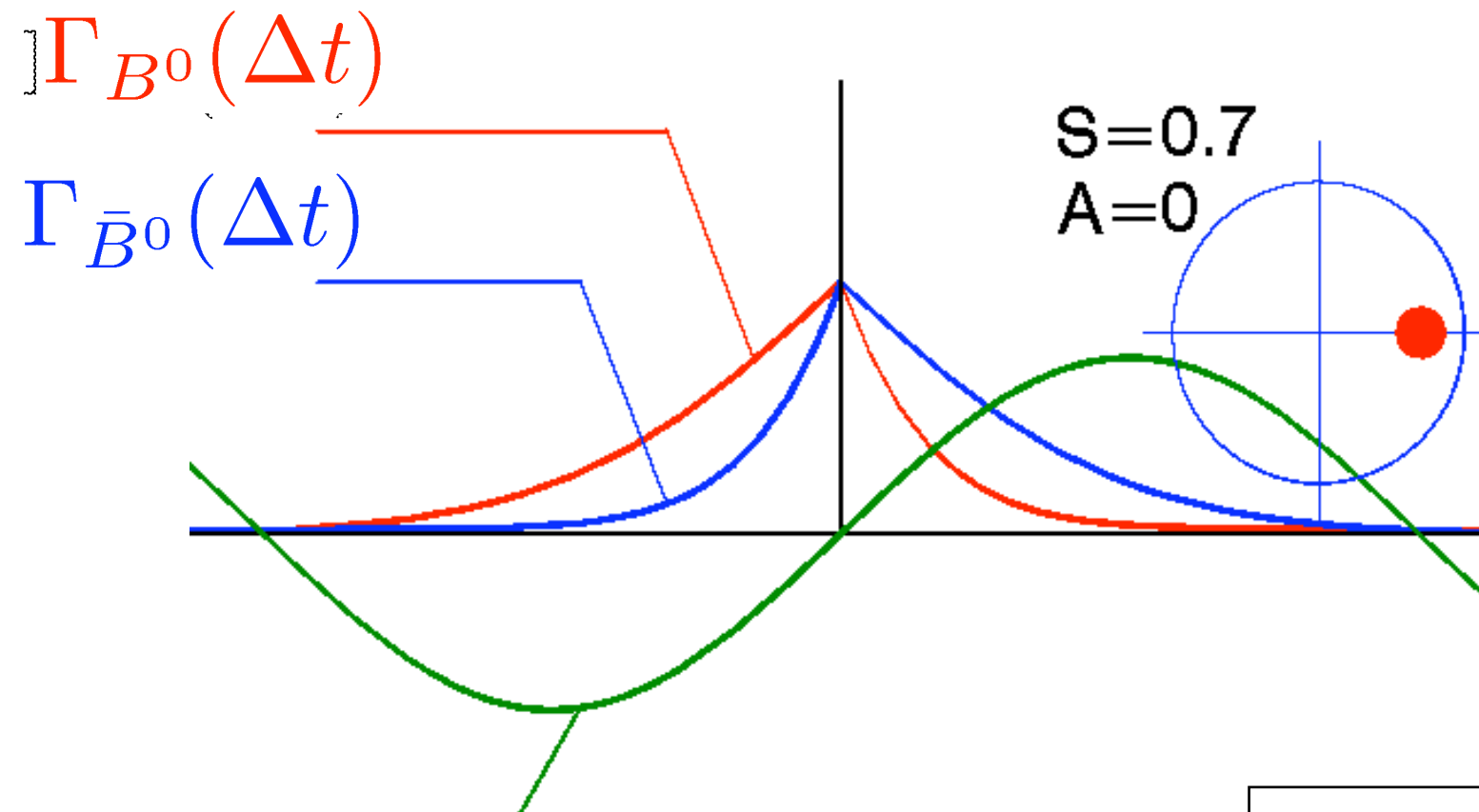
CKM U.T. Angles

- Extract the three angles ϕ_1, ϕ_2, ϕ_3 through time-dependent A_{CP} meas'mt.
- State-of-the-art expt'l technology
 - asymmetric beam energy
 - good vertexing $\Delta I \sim 55\mu\text{m}$
 - Flavor tagging
- Most importantly,
--> need lots of luminosity !





T-dep't CPV in B^0 decays



$$\begin{aligned}
 A_{CP}(\Delta t) &\equiv \frac{\Gamma_{\bar{B}^0}(\Delta t) - \Gamma_{B^0}(\Delta t)}{\Gamma_{\bar{B}^0}(\Delta t) + \Gamma_{B^0}(\Delta t)} \\
 &= \mathcal{S} \sin \Delta m \Delta t + \mathcal{A} \cos \Delta m \Delta t
 \end{aligned}$$

Mixing-induced CPV

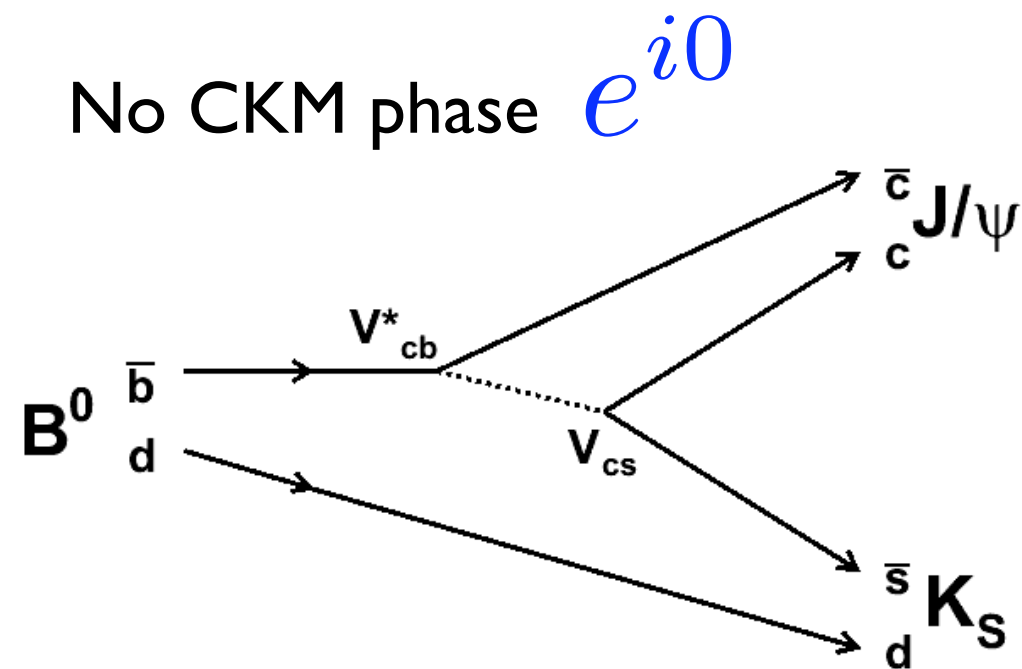
Direct CPV

e.g. for $J/\psi K_S$
 $\mathcal{S} = -\xi_{CP} \sin 2\phi_1 = +\sin 2\phi_1$
 $\mathcal{A} = 0$
 to a good approximation
 (ξ_{CP} : CP eigenvalue)

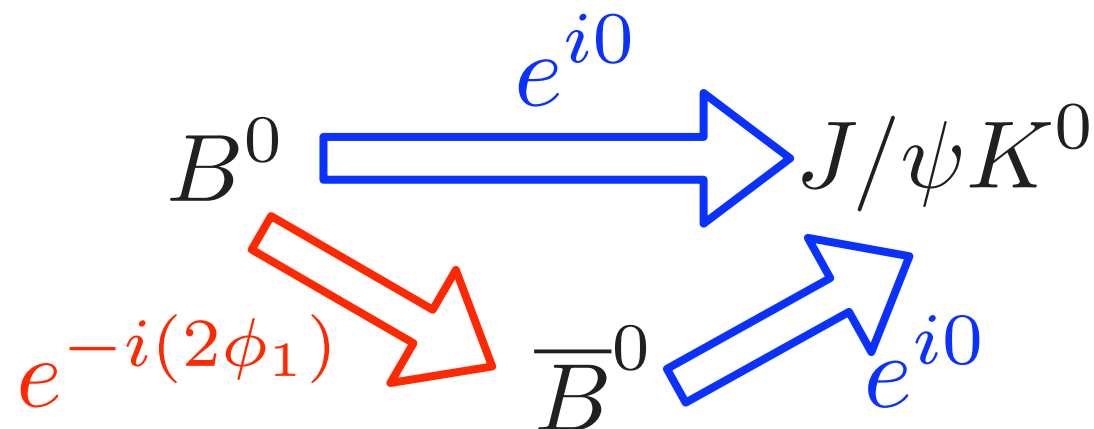
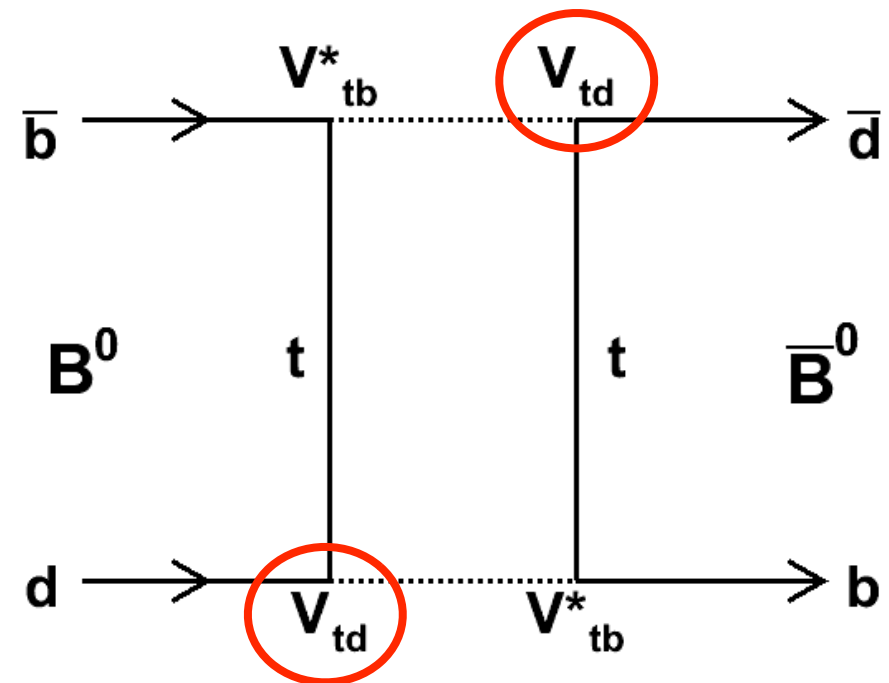
($\mathcal{A} = -C$ a la BaBar)

The Golden mode for ϕ_1

$B^0 \rightarrow J/\psi K^0$; high rate, theoretically clean



Two V_{td} vertices $e^{-i(2\phi_1)}$

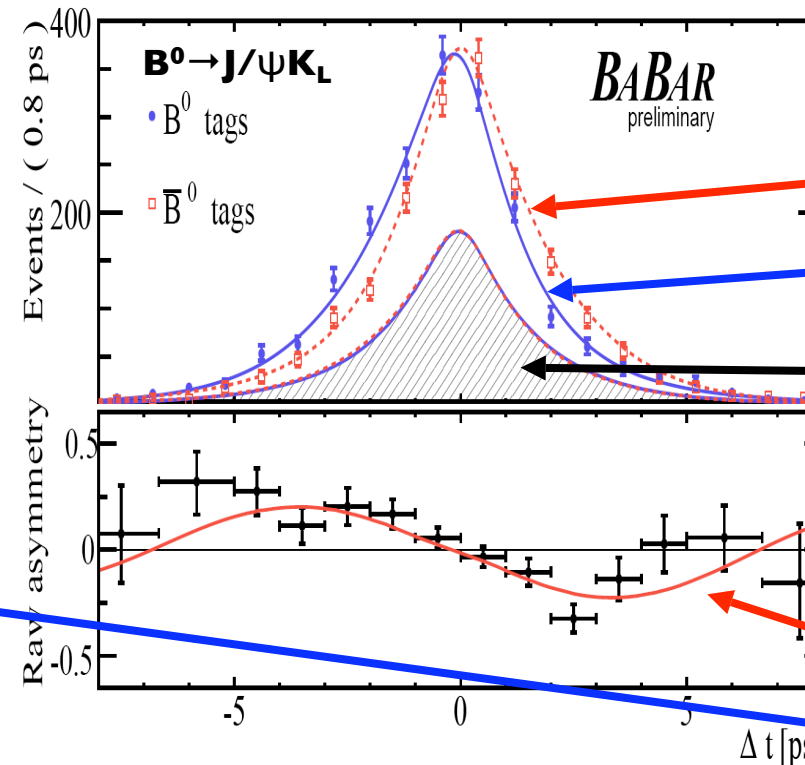
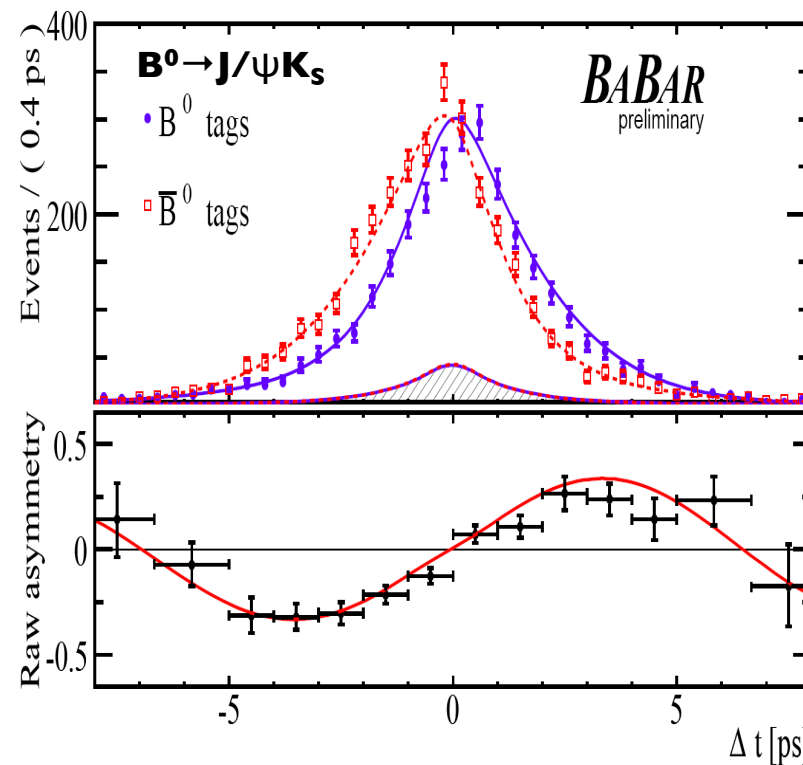


Note: true for **any** B^0 decay **with no phase from decay amplitude**

$A_{CP}(\Delta t)$ time-dependence

BaBar preliminary, hep-ex/0607107

Belle preliminary, hep-ex/0608039



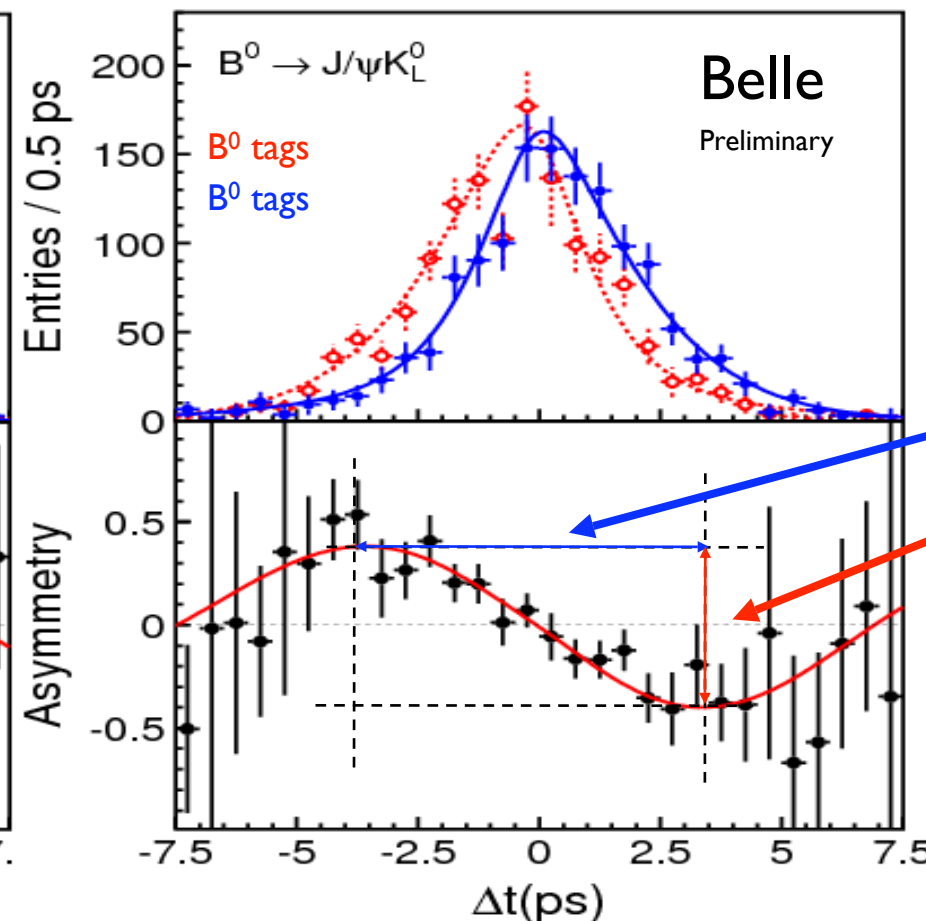
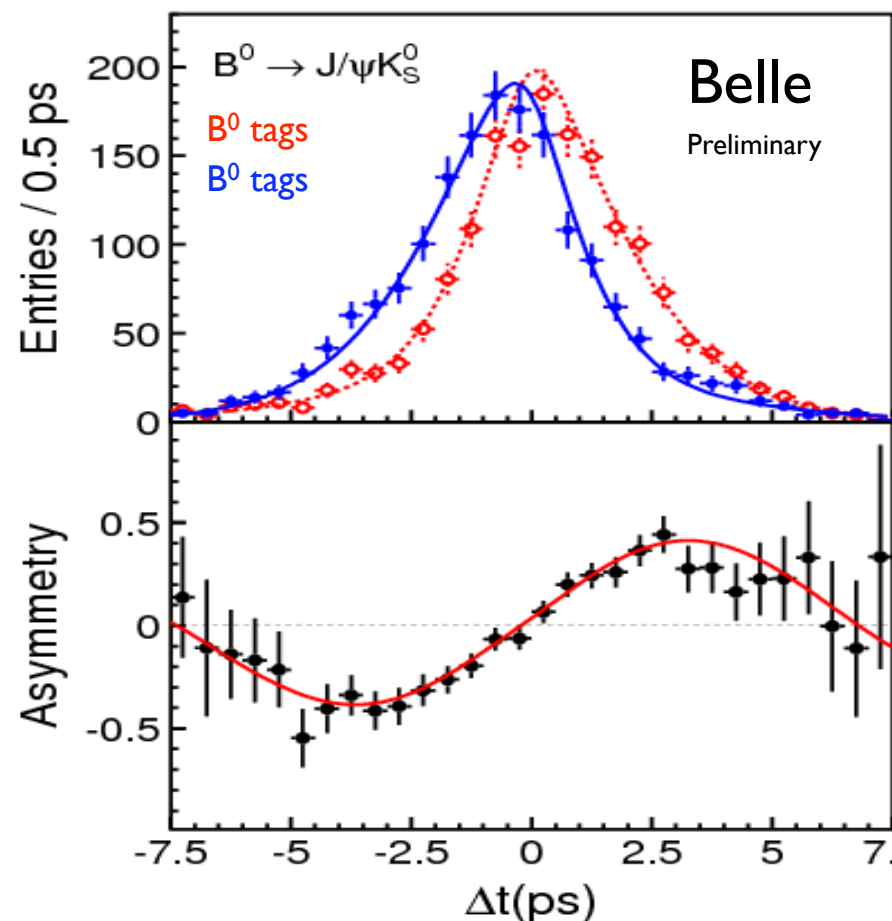
Δt for \bar{B}^0 tag $\approx B^0 \rightarrow J/\psi K$

Δt for B^0 tag $\approx B^0 \rightarrow J/\psi K$

Background

$J/\psi K_L : f_{CP} = +1$

$J/\psi K_S : f_{CP} = -1$

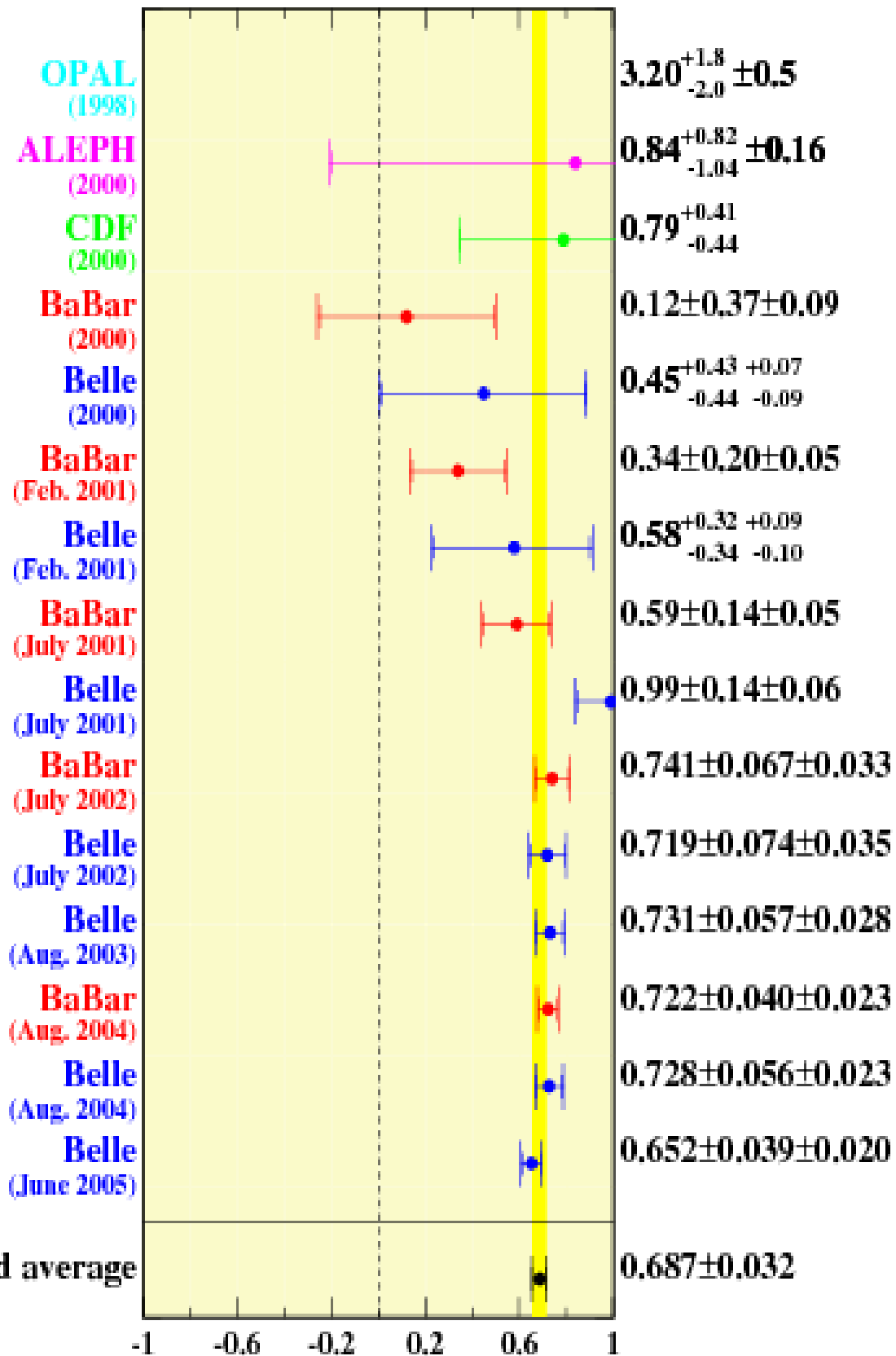


Period = B mixing Δm

Amplitude = $D \sin 2\phi_1$

(Dilution D due to mistags; measured experimentally)

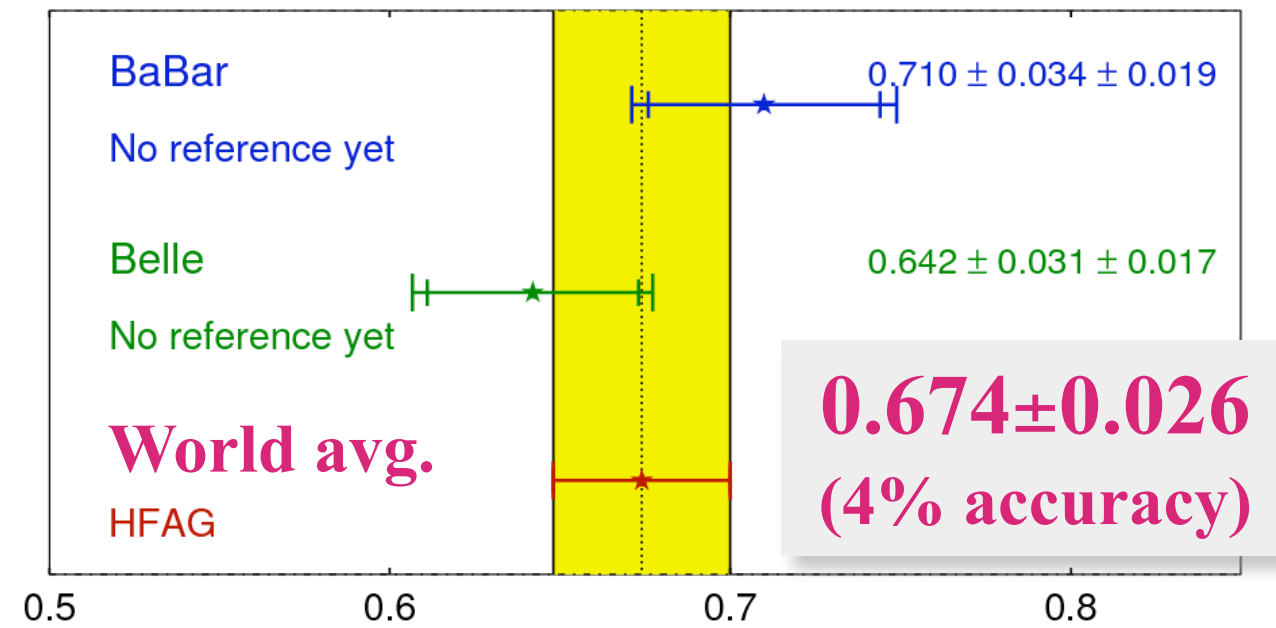
sin2β history (1998-2005)



2006 BaBar + Belle

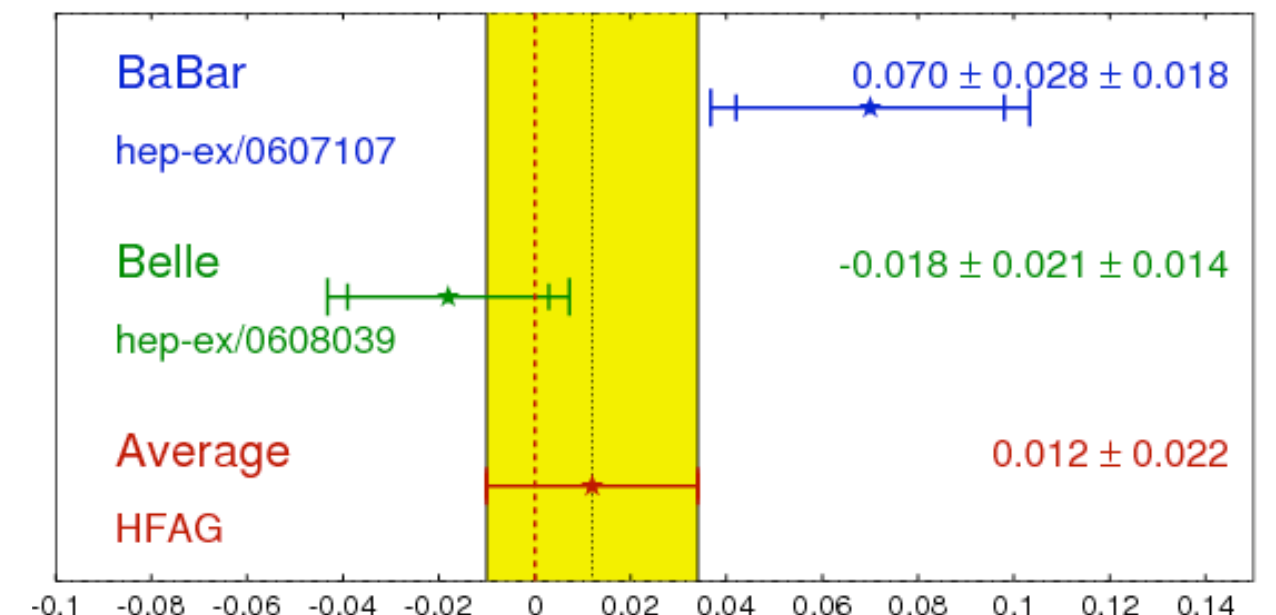
$$S_{CP} = \sin(2\beta) \equiv \sin(2\phi_1)$$

HFAG
ICHEP 2006
PRELIMINARY

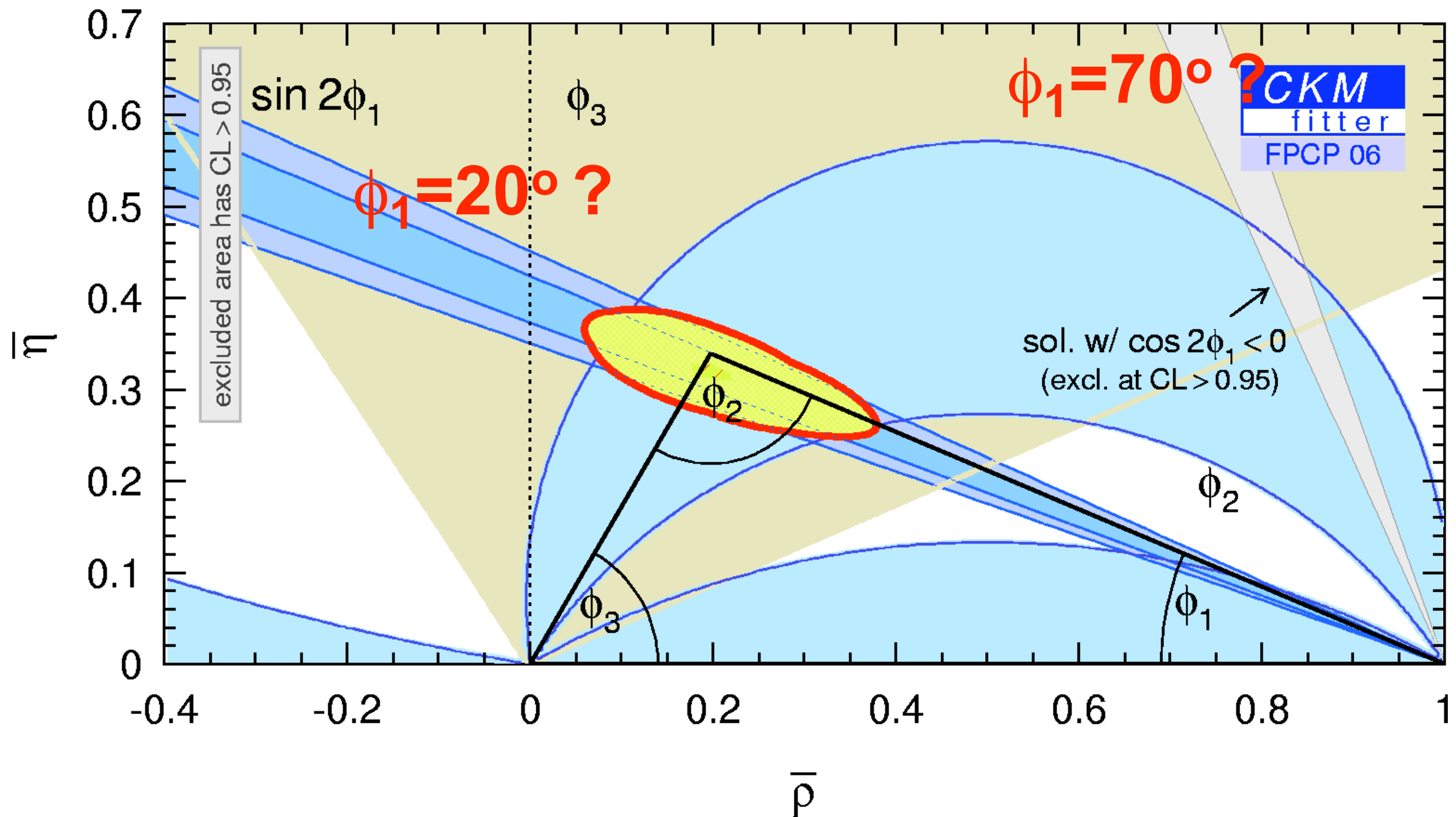


$$b \rightarrow ccs \ C_{CP}$$

HFAG
ICHEP 2006
PRELIMINARY



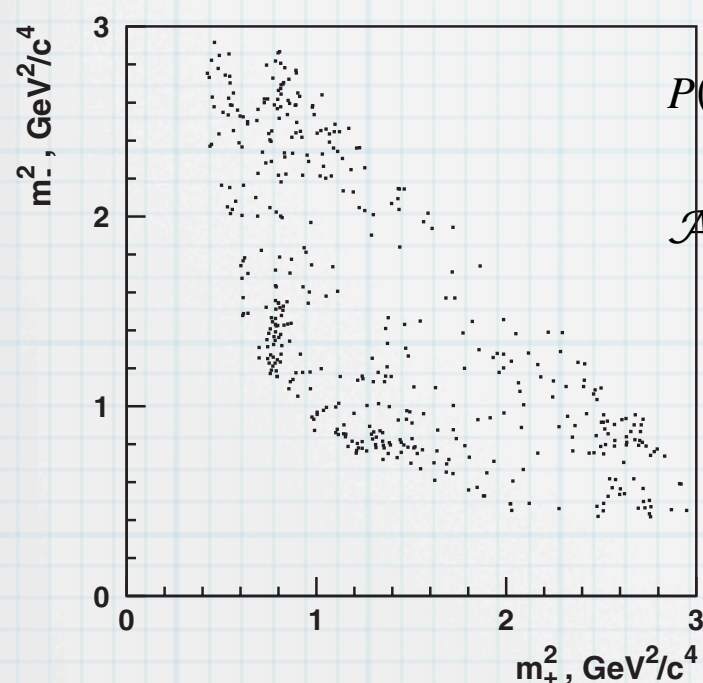
Intrinsic ambiguity still remains ... $2\phi_1 \leftrightarrow \pi - 2\phi_1$





Measurement of the Quark Mixing Parameter $\cos 2\phi_1$ Using Time-Dependent Dalitz Analysis of $\bar{B}^0 \rightarrow D[K_S^0 \pi^+ \pi^-] h^0$

P. Krokovny,⁵ K. Abe,⁵ K. Abe,³⁶ I. Adachi,⁵ H. Aihara,³⁸ D. Anipko,¹ K. Arinstein,¹ Y. Asano,⁴¹ V. Aulchenko,¹ T. Aushev,⁸ S. Bahinipati,³ A.M. Bakich,³³ V. Balagura,⁸ E. Barberio,¹⁵ A. Bay,¹³ U. Bitenc,⁹ I. Bizjak,⁹ A. Bondar,¹ A. Bozek,²¹ M. Brečko,^{5,14,9} T.E. Browder,⁴ V. Chee,²⁰ A. Chen,¹⁸ W.T. Chen,¹⁸ Y. Choi,³² A. Chuvpikar,²⁸ S. Cole,³³



$$P(m_+^2, m_-^2, \Delta t, q_B) = \frac{e^{-|\Delta t|/\tau_{B^0}}}{8\tau_{B^0}} \frac{F(m_+^2, m_-^2)}{2N} (1 + q_B \times \{\mathcal{A}(m_-^2, m_+^2) \cos(\Delta m \Delta t) + \mathcal{S}(m_-^2, m_+^2) \sin(\Delta m \Delta t)\}),$$

$$\mathcal{A} = (|f(m_-^2, m_+^2)|^2 - |f(m_+^2, m_-^2)|^2)/F(m_+^2, m_-^2), \quad \mathcal{S} = \frac{-2\xi_{h^0}(-1)^l \text{Im}\{f(m_-^2, m_+^2)f^*(m_+^2, m_-^2)e^{2i\phi_1}\}}{F(m_+^2, m_-^2)},$$

TABLE II. Fit results for the data. Errors are statistical only.

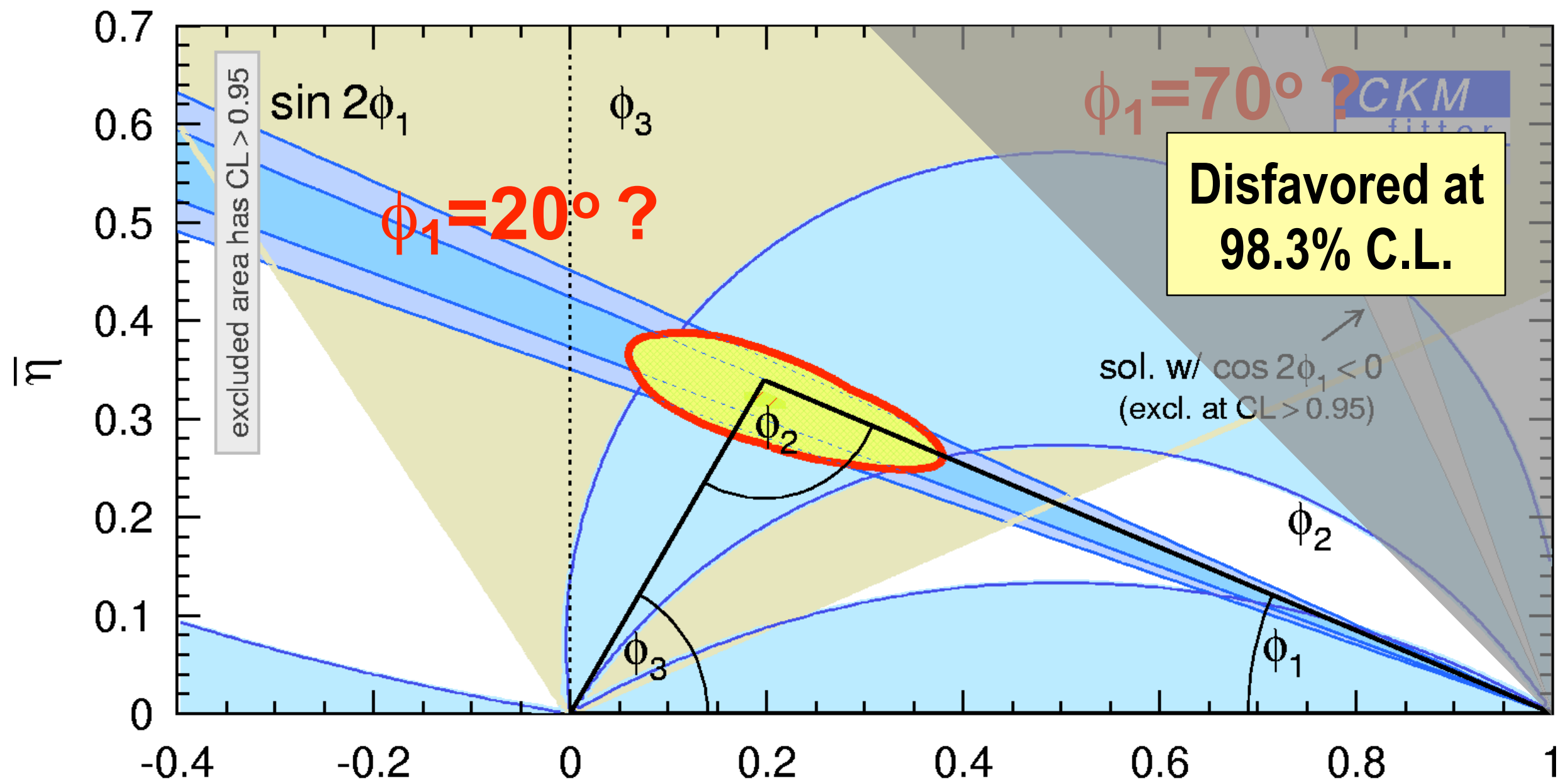
Final state	$\sin 2\phi_1$	$\cos 2\phi_1$
$D\pi^0, D\eta[\gamma\gamma]$	$0.80^{+0.54}_{-0.60}$	$2.07^{+0.78}_{-0.91}$
$D\omega, D\eta[3\pi]$	0.43 ± 0.90	$1.53^{+0.67}_{-0.93}$
$D^*\pi^0, D^*\eta$	1.07 ± 1.14	$3.46^{+1.80}_{-2.01}$
Simultaneous fit	0.78 ± 0.44	$1.87^{+0.40}_{-0.53}$

Process	N_{tot}	Efficiency (%)	N_{sig}	Purity
$D\pi^0$	265	8.7	157 ± 24	59%
$D\omega$	88	4.1	67 ± 10	76%
$D\eta$	101	3.9	58 ± 13	57%
$D^*\pi^0, D^*\eta$	67		43 ± 12	64%
Sum	521		325 ± 31	62%

In summary, we have presented a new method to measure the unitarity triangle angle ϕ_1 using a time-dependent amplitude analysis of the $D \rightarrow K_S^0 \pi^+ \pi^-$ decay produced in the processes $\bar{B}^0 \rightarrow D^{(*)} h^0$. We find $\sin 2\phi_1 = 0.78 \pm 0.44 \pm 0.22$ and $\cos 2\phi_1 = 1.87^{+0.40+0.22}_{-0.53-0.32}$. The sign of $\cos 2\phi_1$ is determined to be positive at 98.3% C.L., favoring the $\phi_1 = 23^\circ$ solution.

Intrinsic ambiguity: resolved at C.L. > 95%

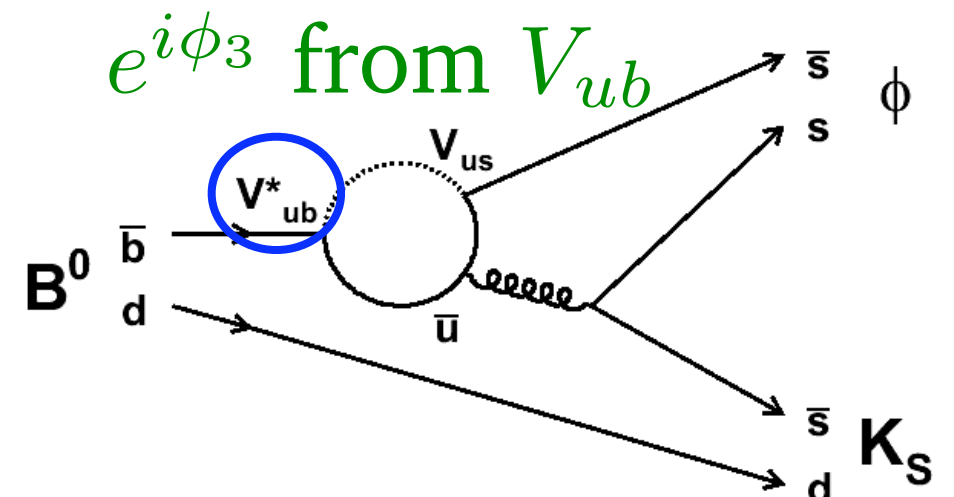
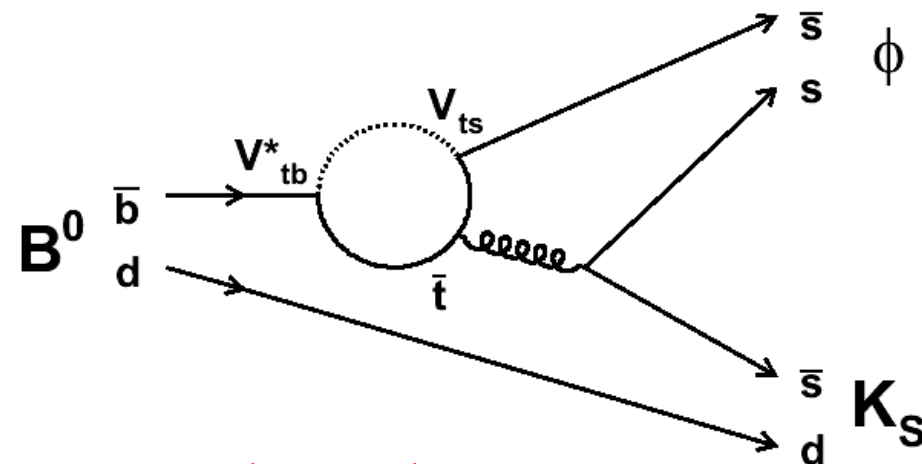
Time-dependent Dalitz analysis of $\bar{B}^0 \rightarrow D[K_S \pi^+ \pi^-] h^0$
PRL97, 081801 (2006)



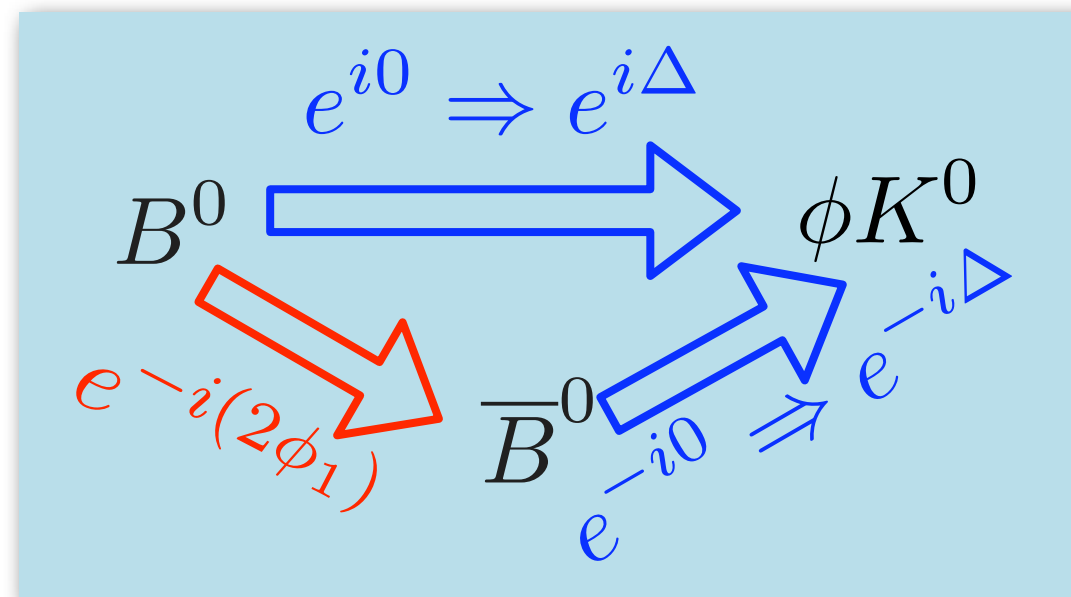
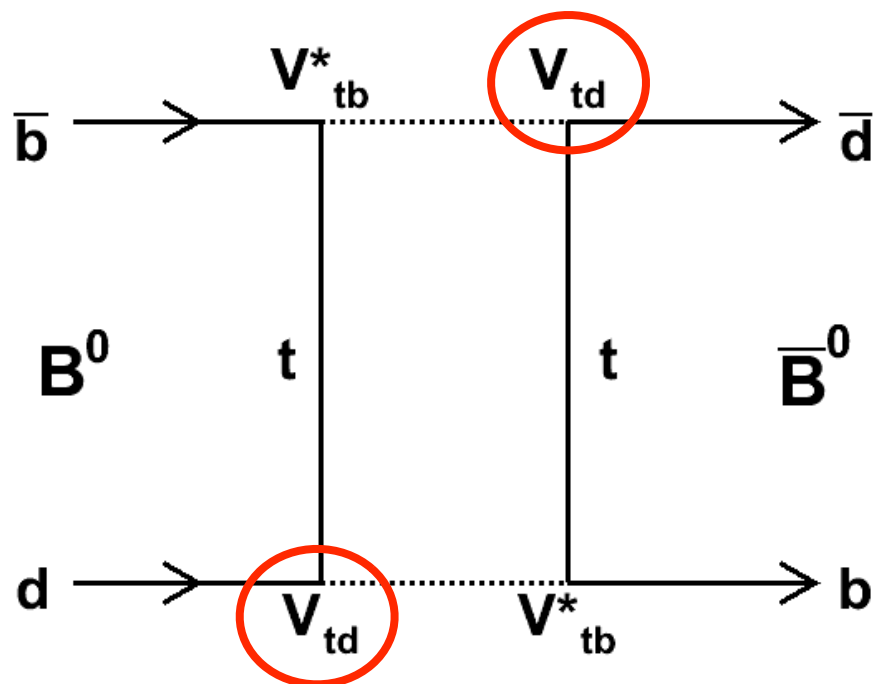
Similar result from \bar{p} BaBar, too. [hep-ex/0607105](#)₂₃

ϕ_1 from $b \rightarrow s\bar{s}s$

e.g. $B^0 \rightarrow \phi K_S$



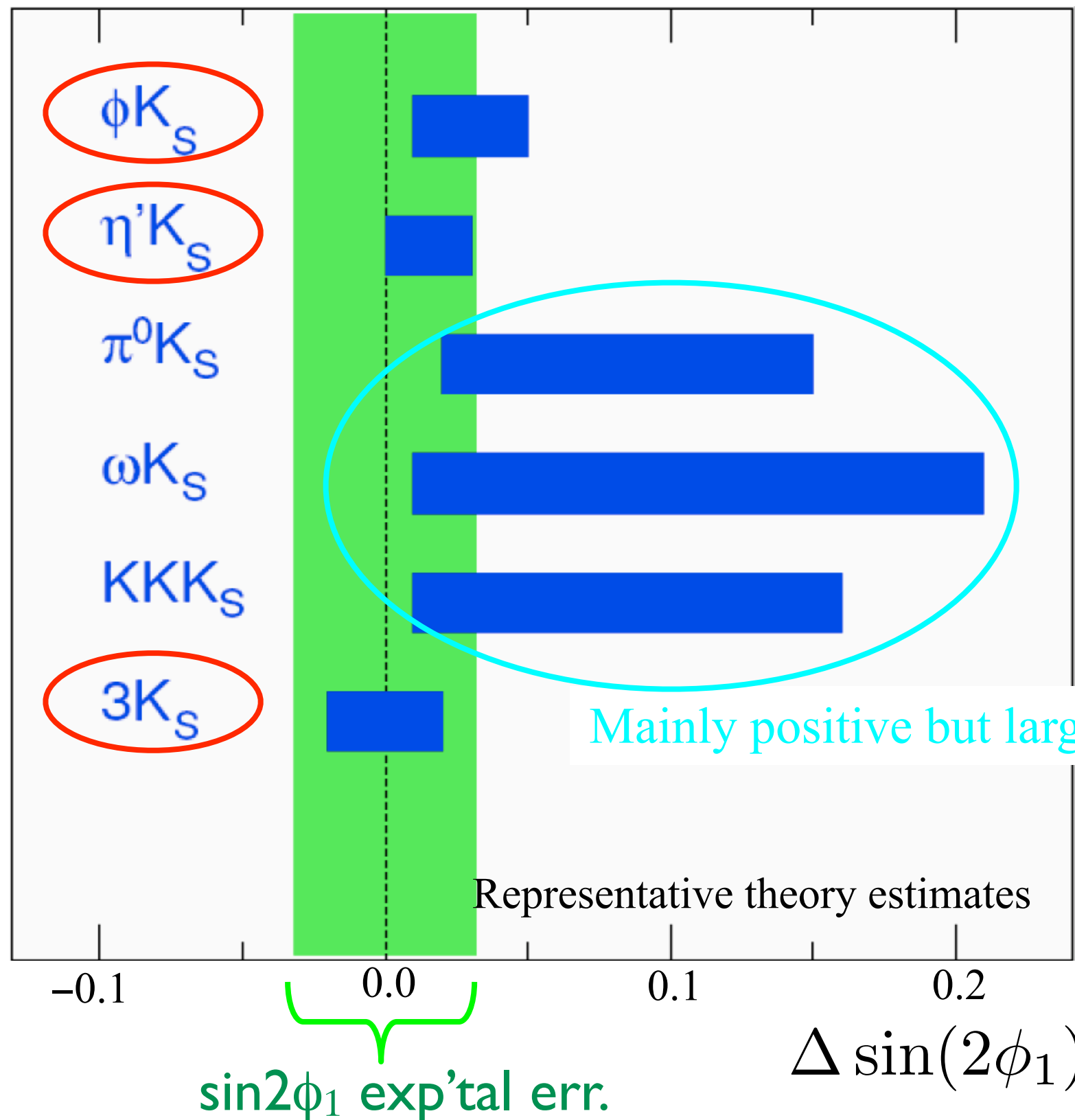
Two V_{td} vertices $e^{-i(2\phi_1)}$



Relative phase = $e^{i2(\phi_1 + \Delta)} \neq e^{i2\phi_1}$

$\sin(2\phi_{1\text{eff}}) \neq \sin(2\phi_1)$

Theoretical estimates of SM offsets



Lazzaro, ICHEP06

Short distance effect:

QCDF:

Beneke, PLB 620, 143 (2005)

Cheng, Chua, Yang, PRD 73, 014017 (2006)

pQCD:

Mishima, Sanda, PRD 72, 114005 (2005)

SCET:

Williamson, Zupan, PRD 74, 014003 (2006)

Long distance effect:

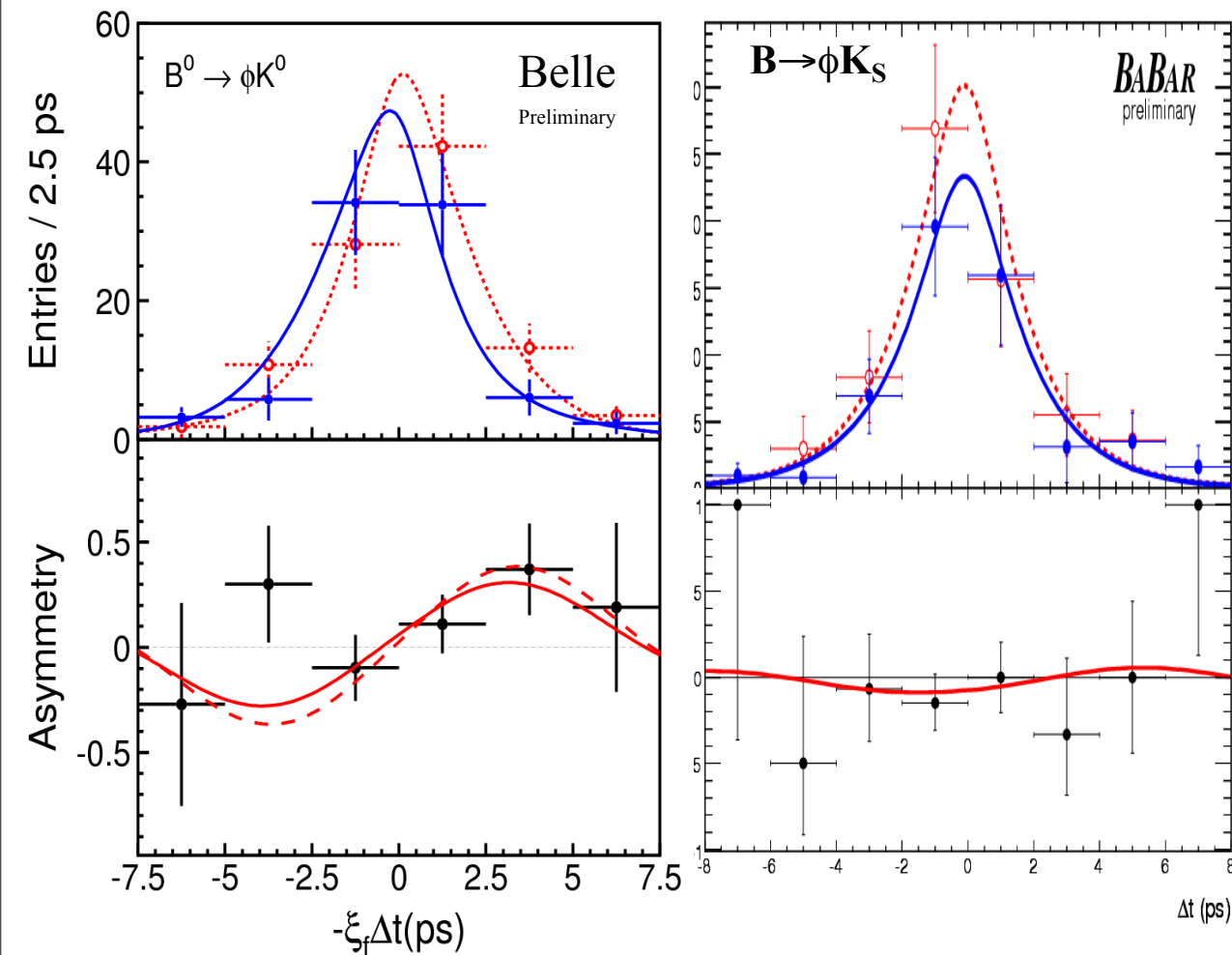
Cheng, Chua, Soni, PRD 72, 014006 (2005)

$B \rightarrow \phi K_S$ and $B \rightarrow \eta' K_S$

$$B \rightarrow \phi K_S: \text{B.F.} = (4.3 \pm 0.6) \times 10^{-6}$$

Even with small rates,
clear CP violation observed
in penguin decays

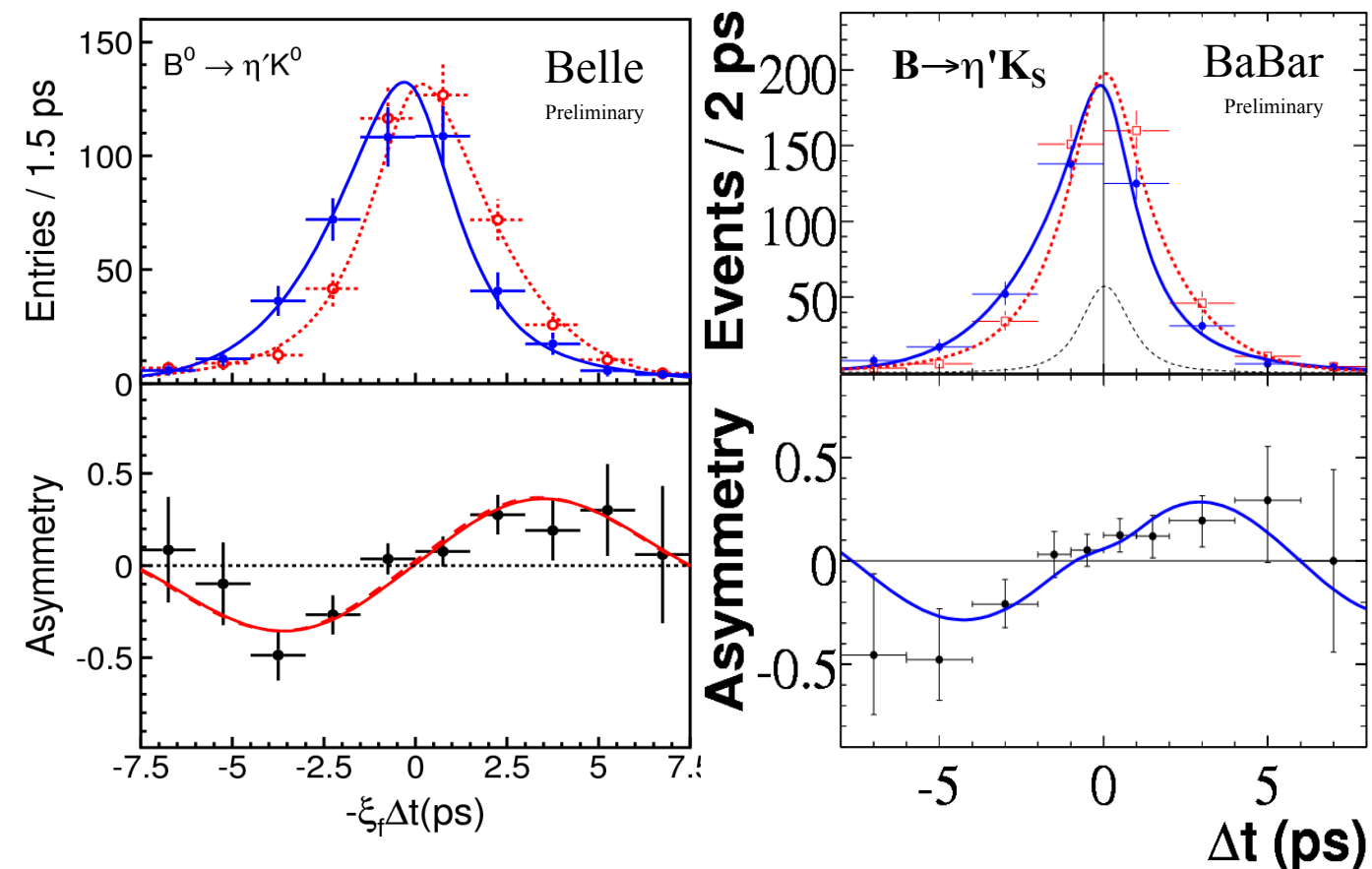
$$B \rightarrow \eta' K_S: \text{B.F.} = (3.4 \pm 0.2) \times 10^{-5}$$



Belle preliminary, hep-ex/0608039

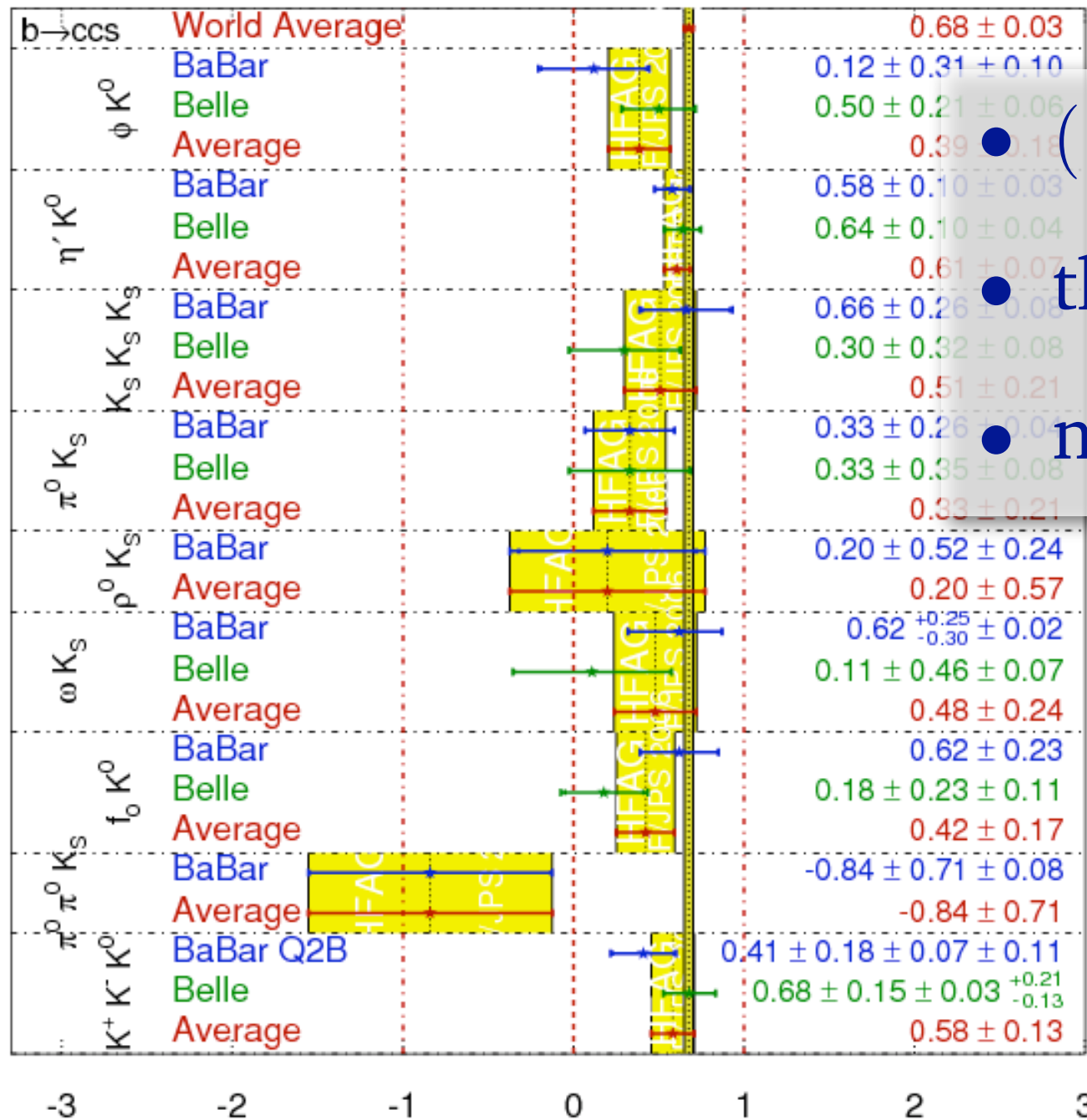
BaBar preliminary, hep-ex/0607112

BaBar preliminary, hep-ex/0609052



$$\sin(2\beta^{\text{eff}}) \equiv \sin(2\phi_1^{\text{eff}})$$

HFAG
DPF/JPS 2006
PRELIMINARY



• (−) in all 9 modes

• theory tends to predict (+) shifts

• naive avg. $\sin 2\phi_1^{\text{eff}} = 0.52 \pm 0.05$

$\Rightarrow 2.6 \sigma$ difference

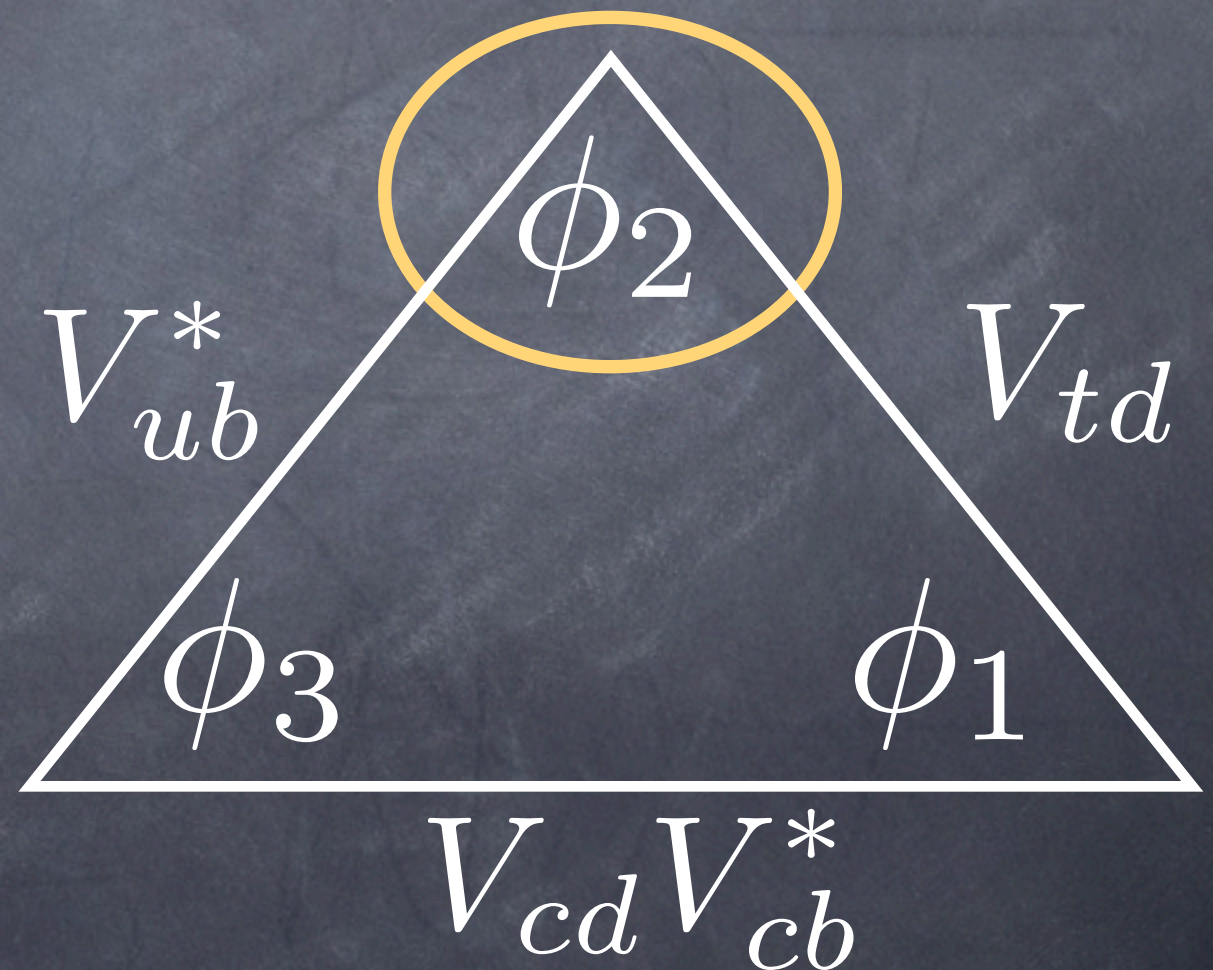
*More stat. is crucial for
mode-by-mode comparison!*

Other angles?



Unitarity triangle angles

BABAR:	β	α	γ
BELLE:	ϕ_1	ϕ_2	ϕ_3
	易	難	魔



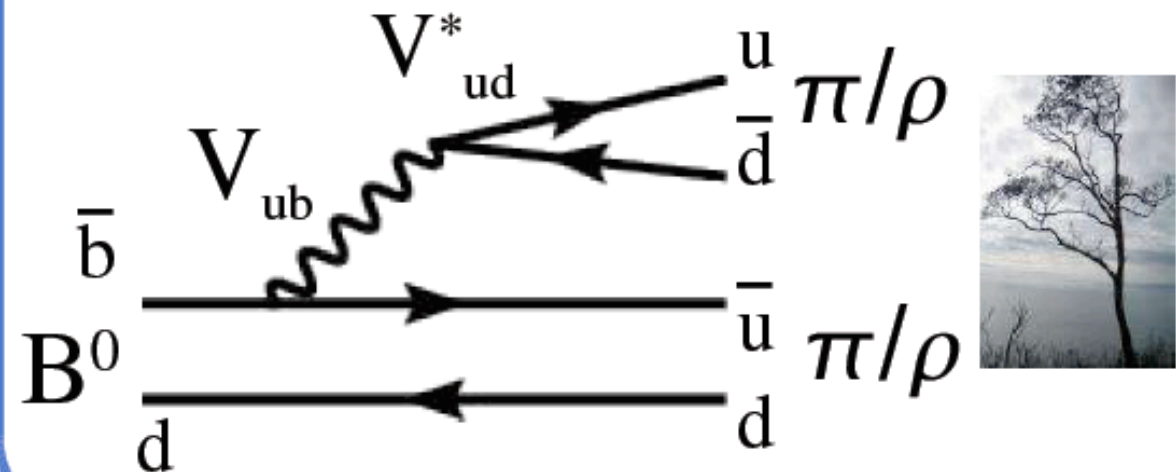
penguin's shaking a tree...

$$B^0 \rightarrow \rho^+ \rho^-$$

$$B^0 \rightarrow \rho^\pm \pi^\mp$$

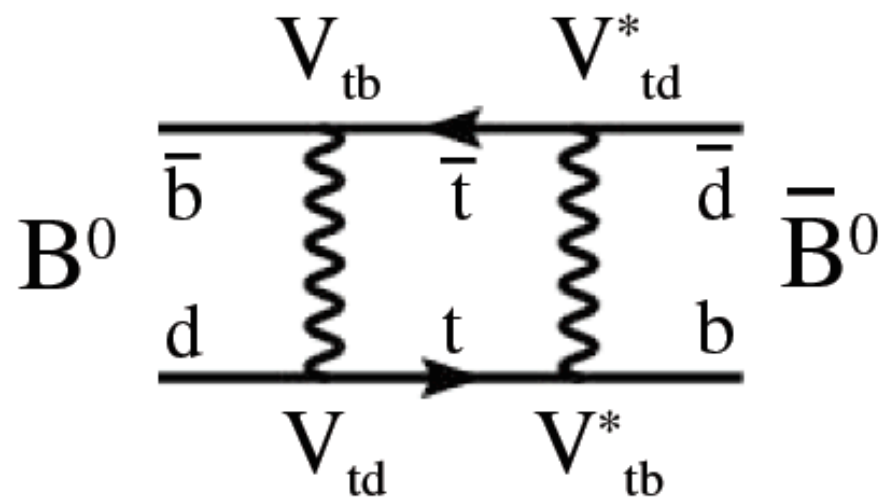
$$B^0 \rightarrow \pi^+ \pi^-$$

Tree diagram

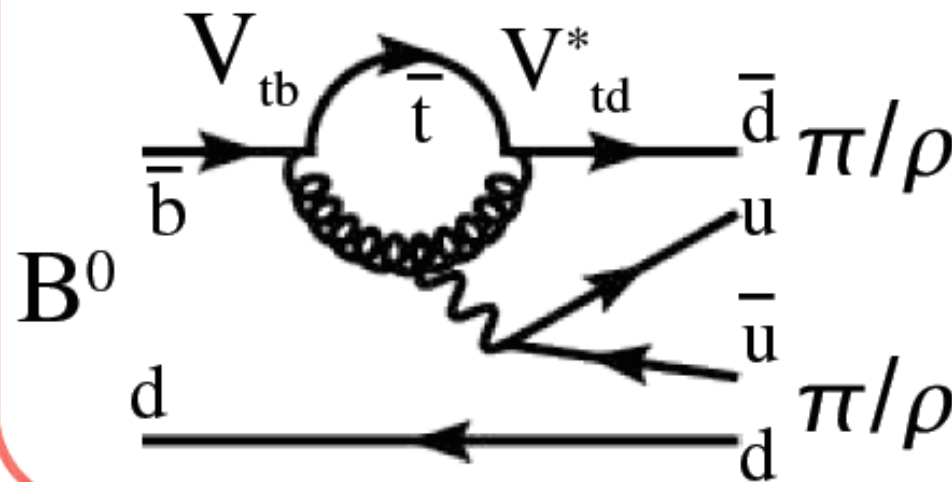


ϕ_2

Mixing



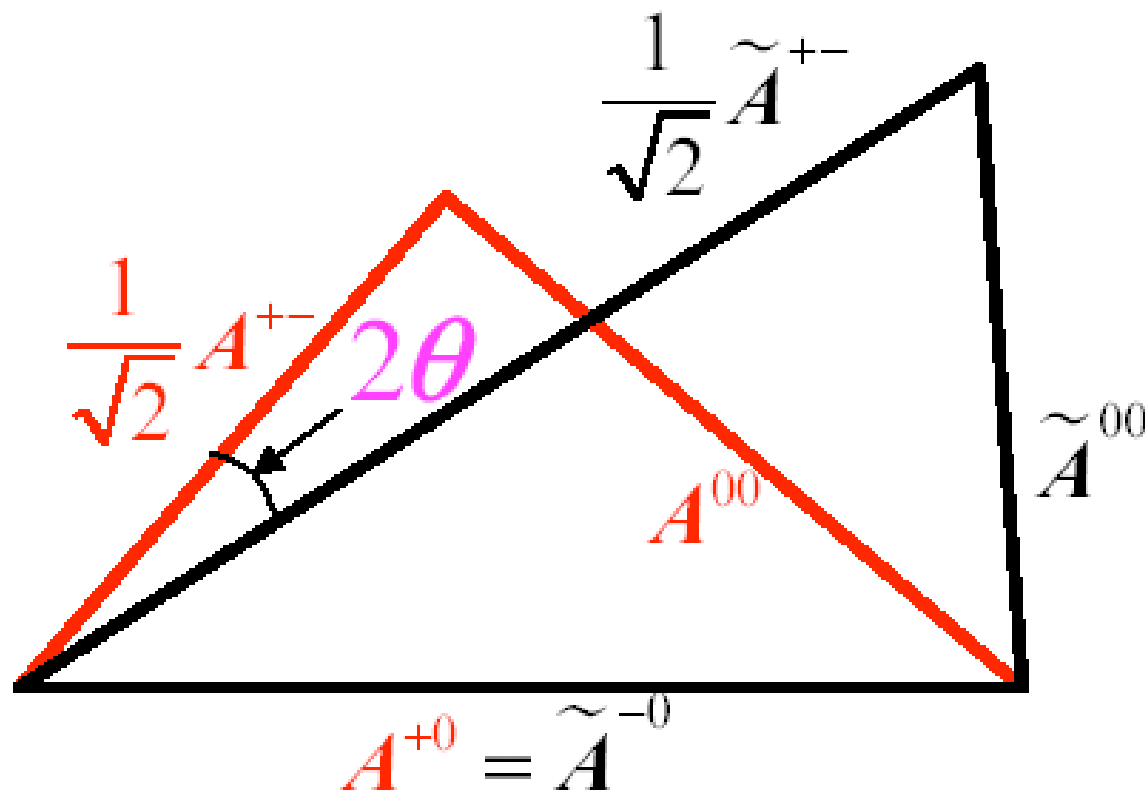
Penguin diagram



0

Isospin analysis: flavor SU(2)

Gronau-London (1990)



	<i>Amplitude for</i>
$A^{+-}(\bar{A}^{+-})$	$B^0(\bar{B}^0) \rightarrow \pi^+ \pi^-$
$A^{00}(\bar{A}^{00})$	$B^0(\bar{B}^0) \rightarrow \pi^0 \pi^0$
$A^{+0}(\bar{A}^{-0})$	$B^+(B^-) \rightarrow \pi^+ \pi^0 (\pi^- \pi^0)$

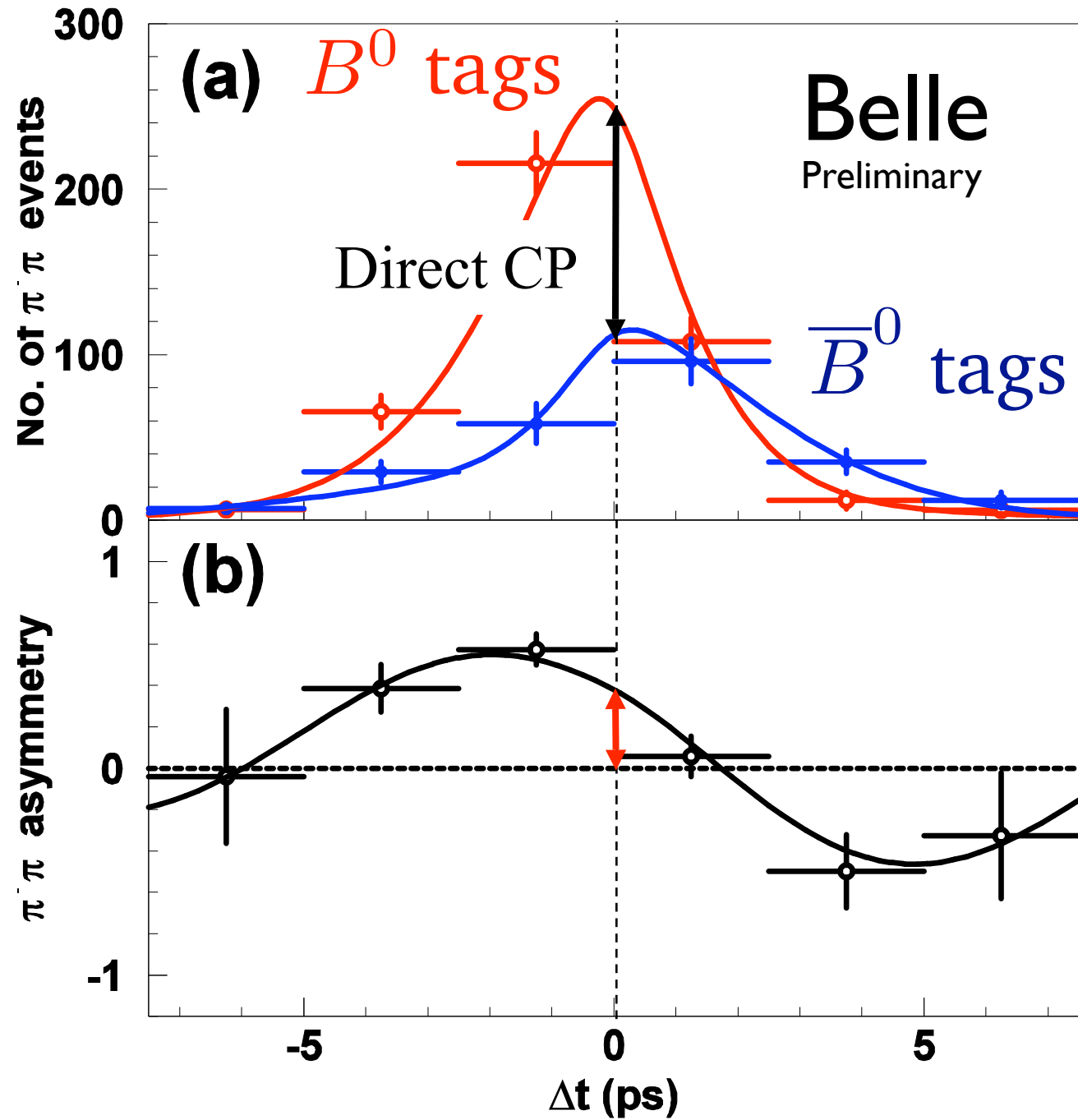
$$\tilde{A}^{ij} = e^{2\phi_3} \bar{A}^{ij}$$

- Model-independent (symmetry-dependent) method
- SU(2) breaking effect well below present statistical errors

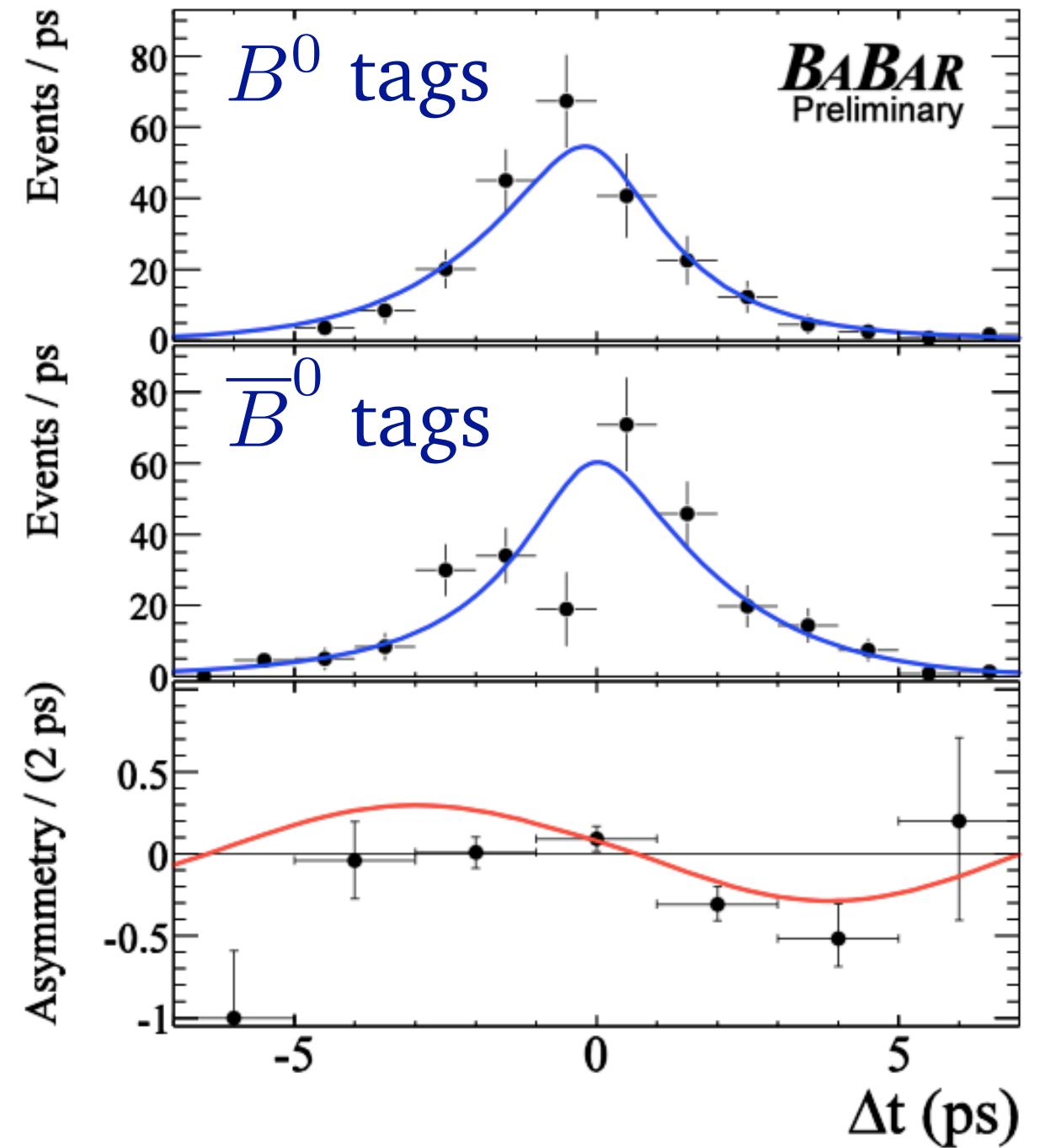
“Penguin pollution” can be removed

$B^0 \rightarrow \pi^+ \pi^-$

Belle preliminary; hep-ex/0608035



BaBar preliminary; hep-ex/0607106

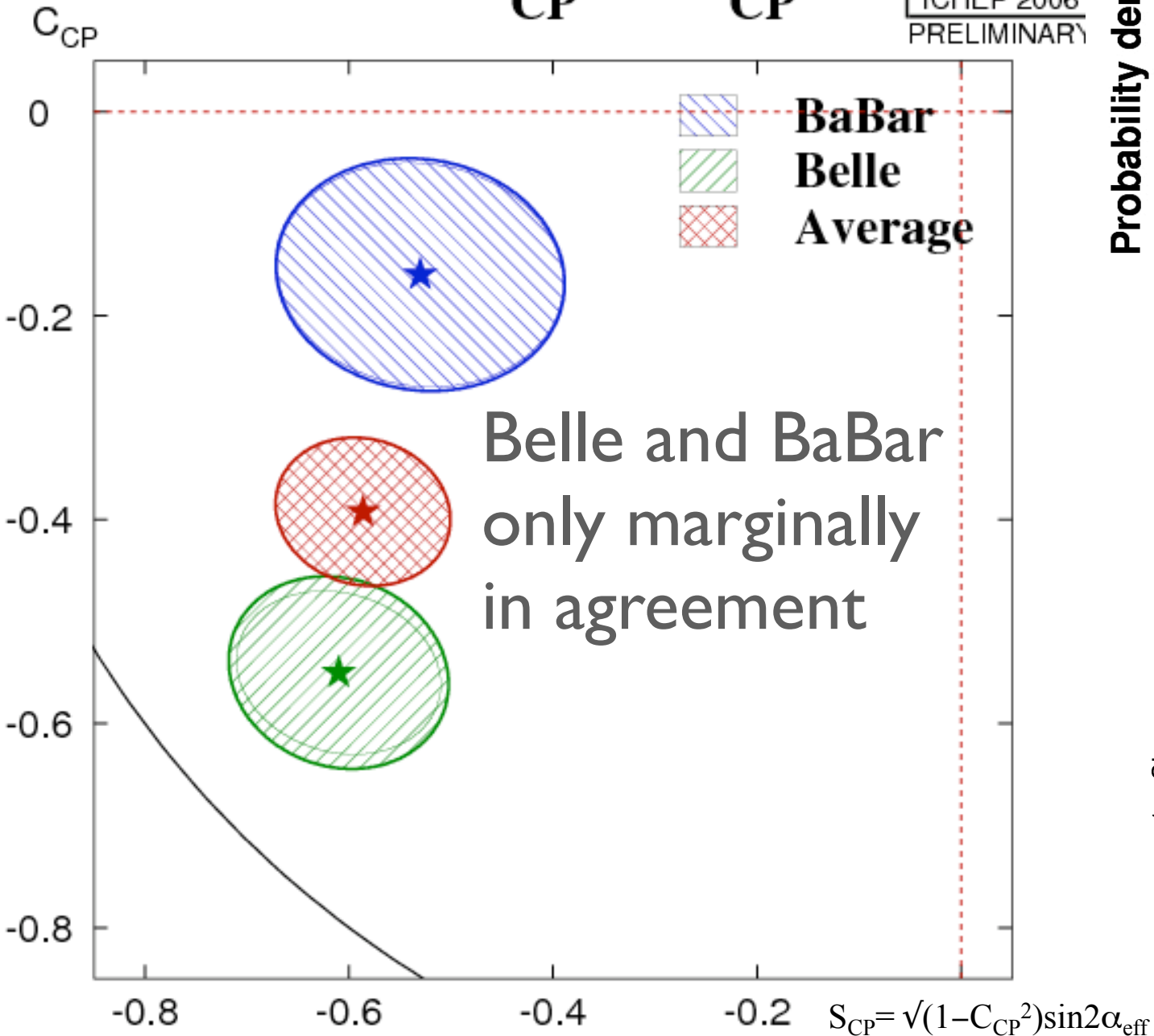


$$C_{CP} \neq 0 \rightarrow S_{CP} = \sqrt{(1 - C_{CP}^2)} \sin 2\alpha_{\text{eff}}$$

$B^0 \rightarrow \pi^+ \pi^-$

$\pi^+ \pi^- S_{CP}$ vs C_{CP}

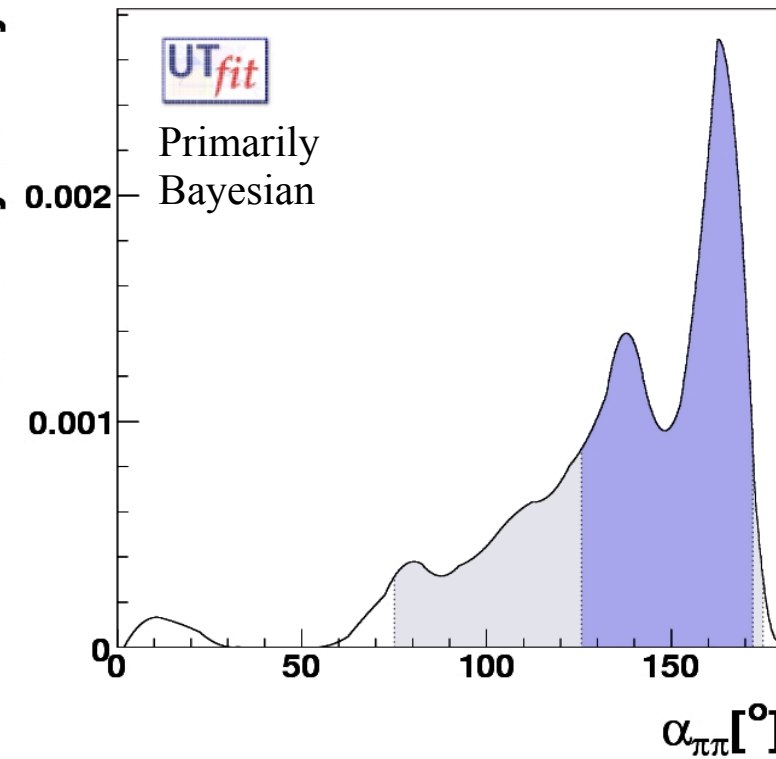
HFAG
ICHEP 2006
PRELIMINARY



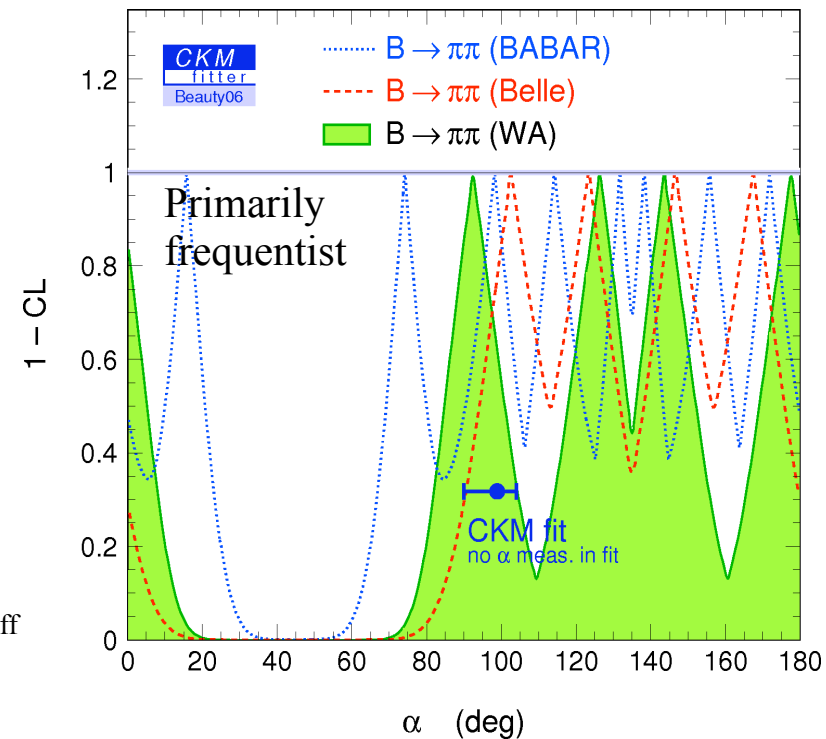
Contours give $-2\Delta(\ln L) = \Delta\chi^2 = 1$, corresponding to 60.7% CL for 2 dof

Including $\pi^+ \pi^0, \pi^0 \pi^0$ isospin analysis

Probability density



UTFit,
M.Bona et al.,
hep-ph/0606167



CKMFitter,
J.Charles et al.,
Eur.Phys.J.C41,
1 (2005)

Bayesian and frequentist interpretations
give quantitatively **different** conclusions

$$\phi_2 \text{ from } B^0 \rightarrow \rho^+ \rho^-$$

Advantages of $\rho^+ \rho^-$:

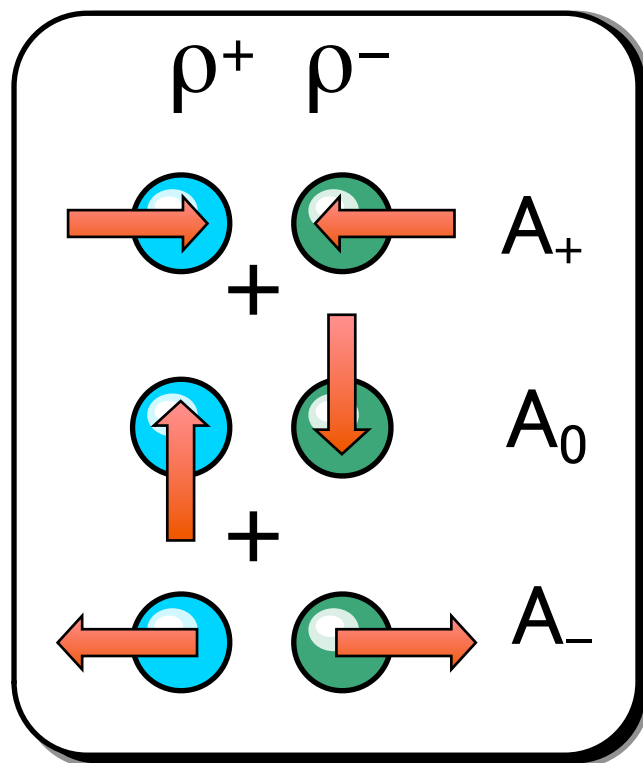
- small penguin contribution due to relatively small

BaBar hep-ex/0612021
subm'd to PRL

$$\text{BF}(B^0 \rightarrow \rho^0 \rho^0) = (1.07 \pm 0.33 \pm 0.19) \times 10^{-6}$$

- relatively large measured branching fraction $\frac{\text{Br}(B^0 \rightarrow \rho^+ \rho^-)}{\text{Br}(B^0 \rightarrow \pi^+ \pi^-)} \sim 4.4$

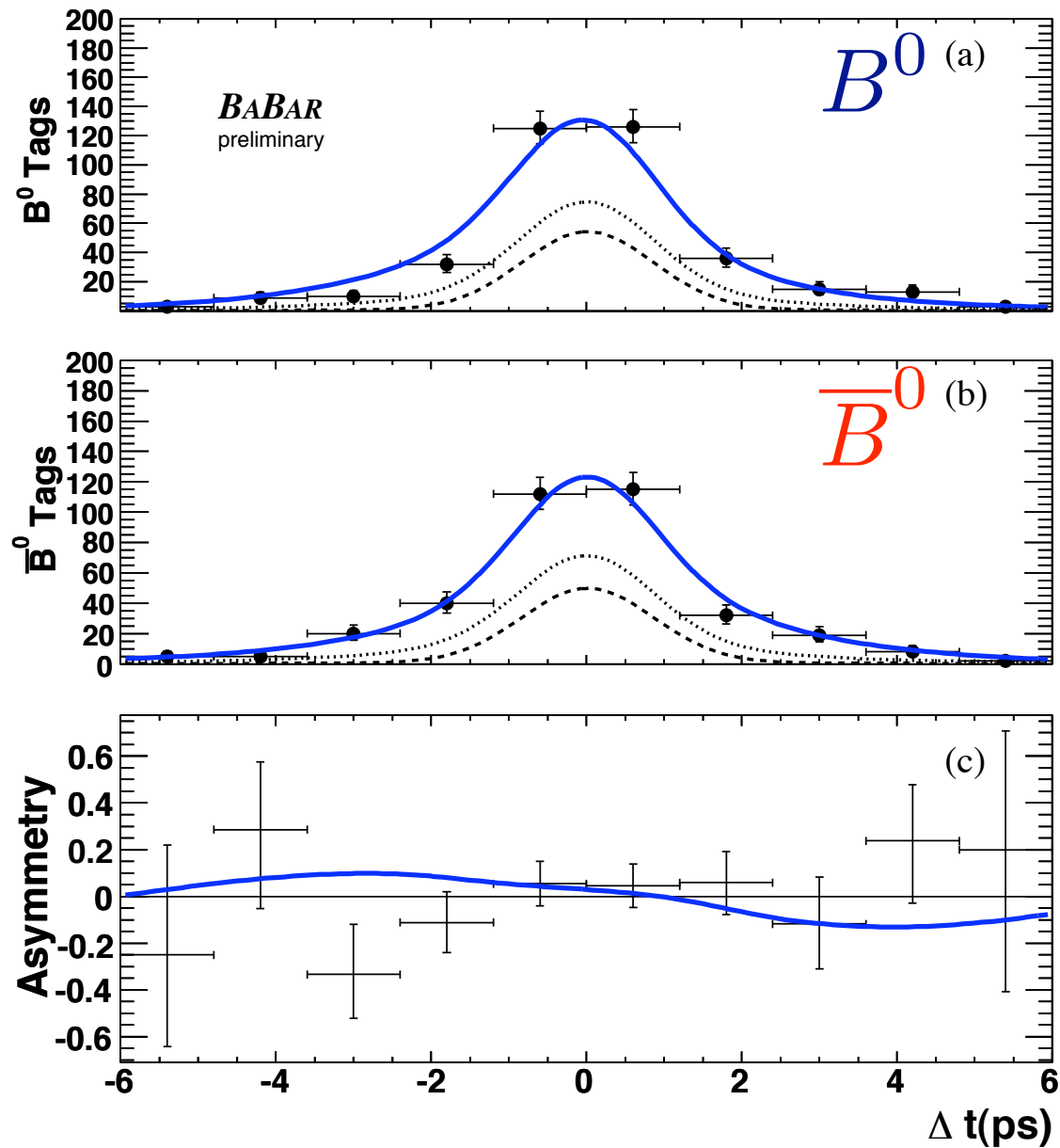
☺ Fortunately, longitudinal polarization dominates



$$\left. \begin{array}{l} \boxed{(A_+ + A_-)/\sqrt{2}} \quad \boxed{A_{\parallel}} \quad +1 \\ \boxed{A_0} \quad +1 \\ \boxed{(A_+ - A_-)/\sqrt{2}} \quad \boxed{A_{\perp}} \quad -1 \end{array} \right\} \text{CP}$$

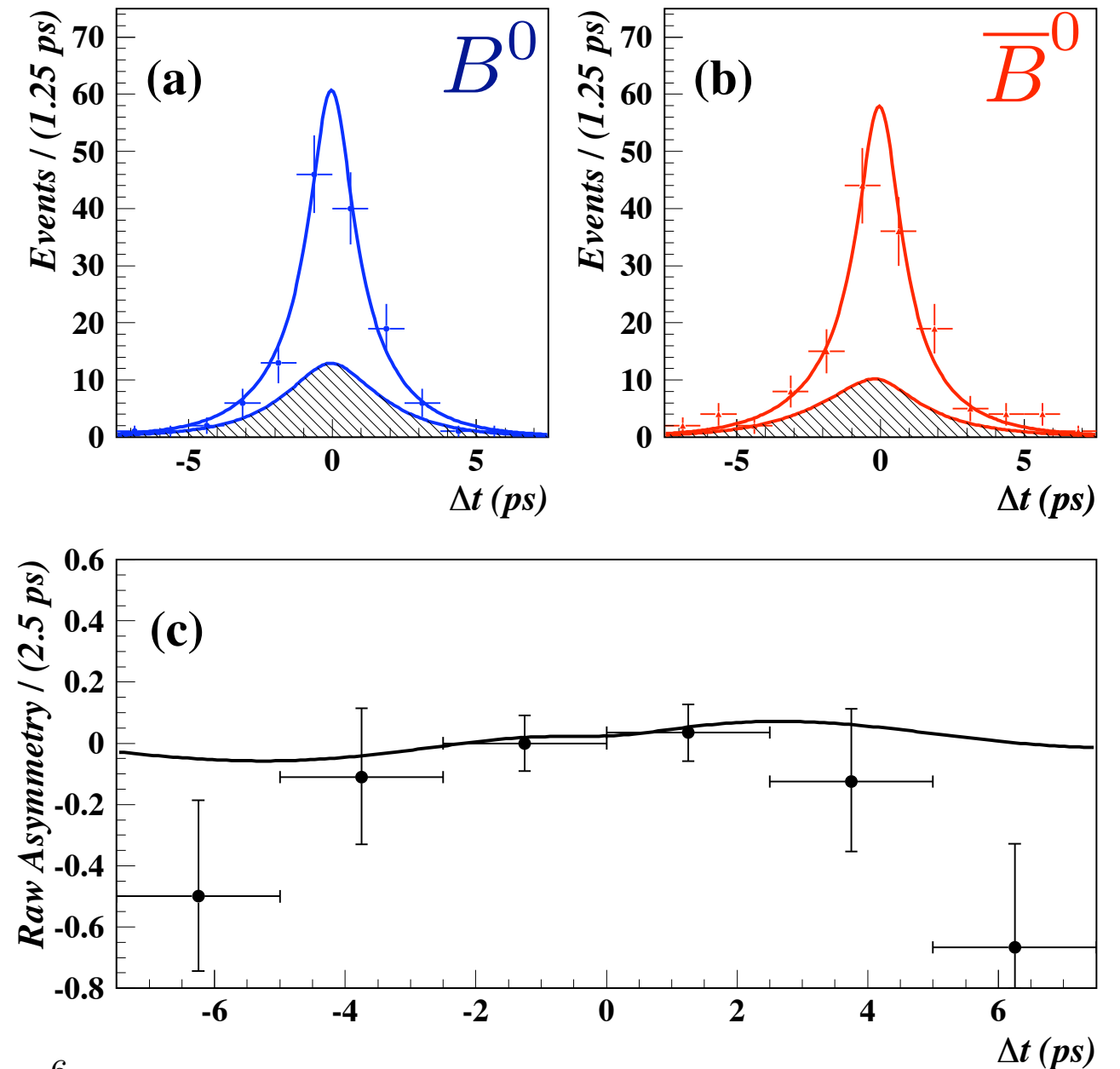
ϕ_2 from $B^0 \rightarrow \rho^+ \rho^-$

BaBar preliminary; hep-ex/0607098



$$\begin{aligned}
 \mathcal{B}(B^0 \rightarrow \rho^+ \rho^-) &= (23.5 \pm 2.2(\text{stat}) \pm 4.1(\text{syst})) \times 10^{-6}, \\
 f_L &= 0.977 \pm 0.024(\text{stat})^{+0.015}_{-0.013}(\text{syst}), \\
 S_{\text{long}} &= -0.19 \pm 0.21(\text{stat})^{+0.05}_{-0.07}(\text{syst}), \\
 C_{\text{long}} &= -0.07 \pm 0.15(\text{stat}) \pm 0.06(\text{syst}).
 \end{aligned}$$

Belle, preliminary; hep-ex/0702009



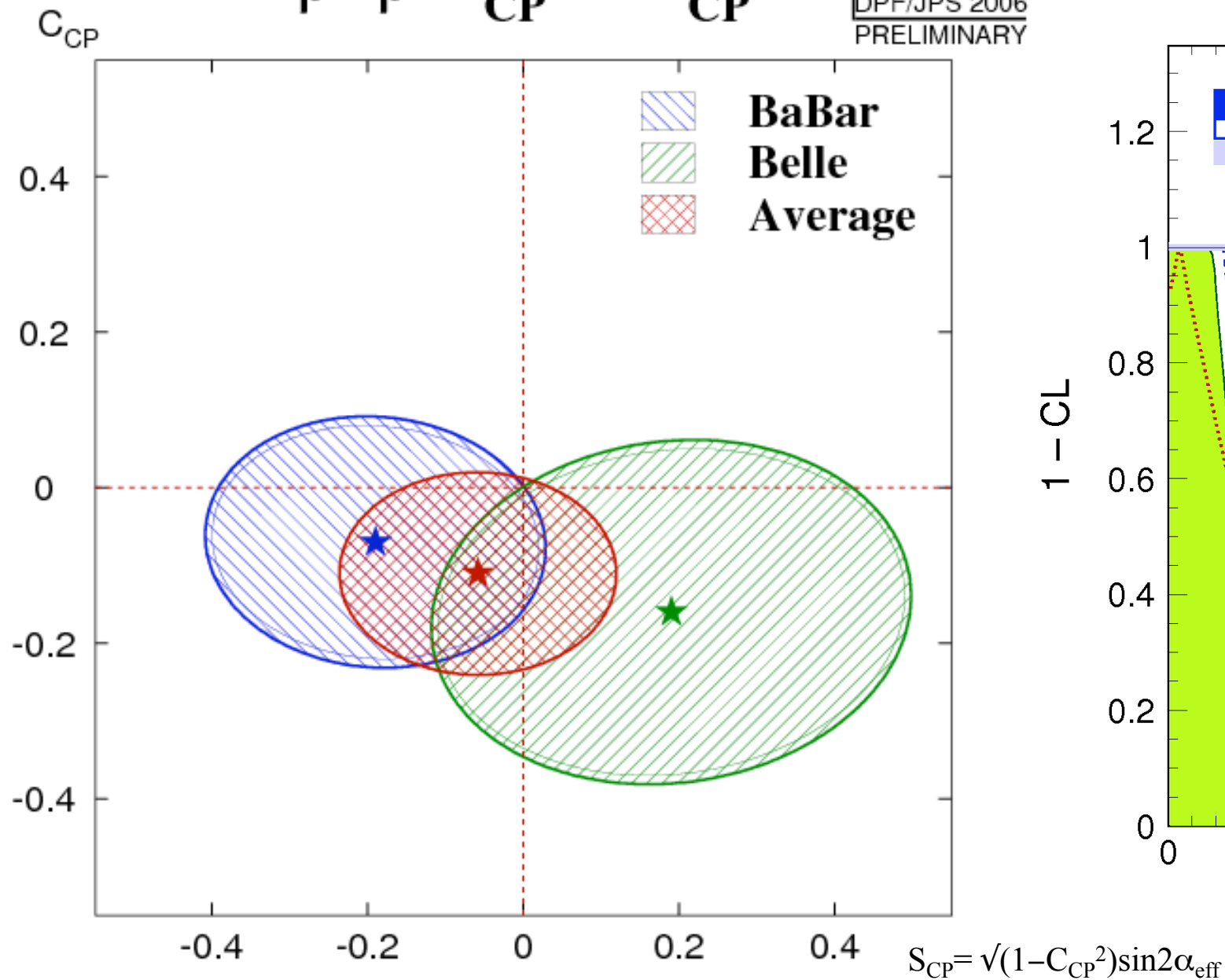
$$\begin{aligned}
 \mathcal{A}_L &= 0.16 \pm 0.21(\text{stat}) \pm 0.07(\text{syst}) \\
 \mathcal{S}_L &= 0.19 \pm 0.30(\text{stat}) \pm 0.07(\text{syst})
 \end{aligned}$$

A, S : both consistent with 0

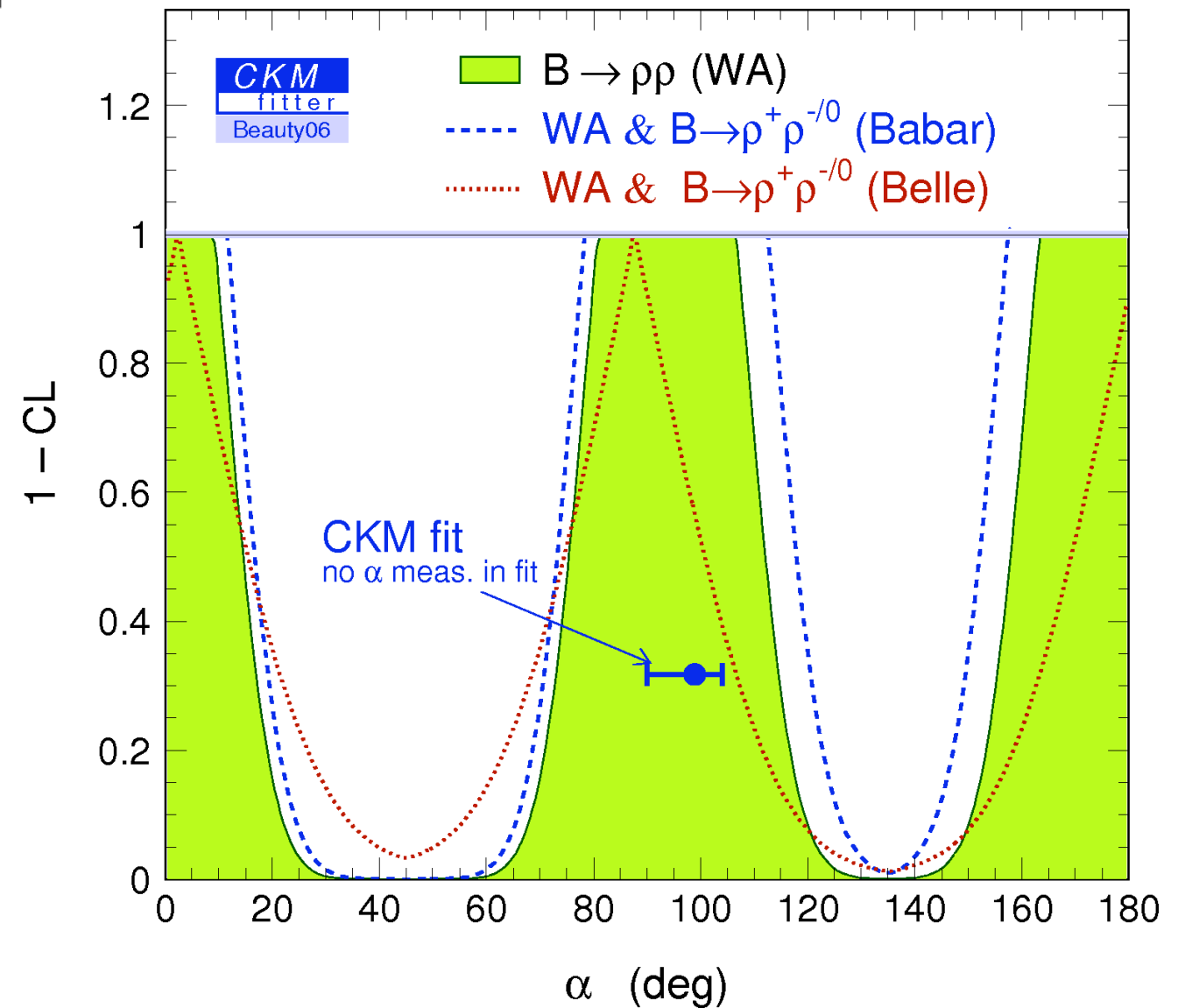
ϕ_2 from $B^0 \rightarrow \rho^+ \rho^-$

$\rho^+ \rho^- S_{CP}$ vs C_{CP}

HFAG
DPF/JPS 2006
PRELIMINARY

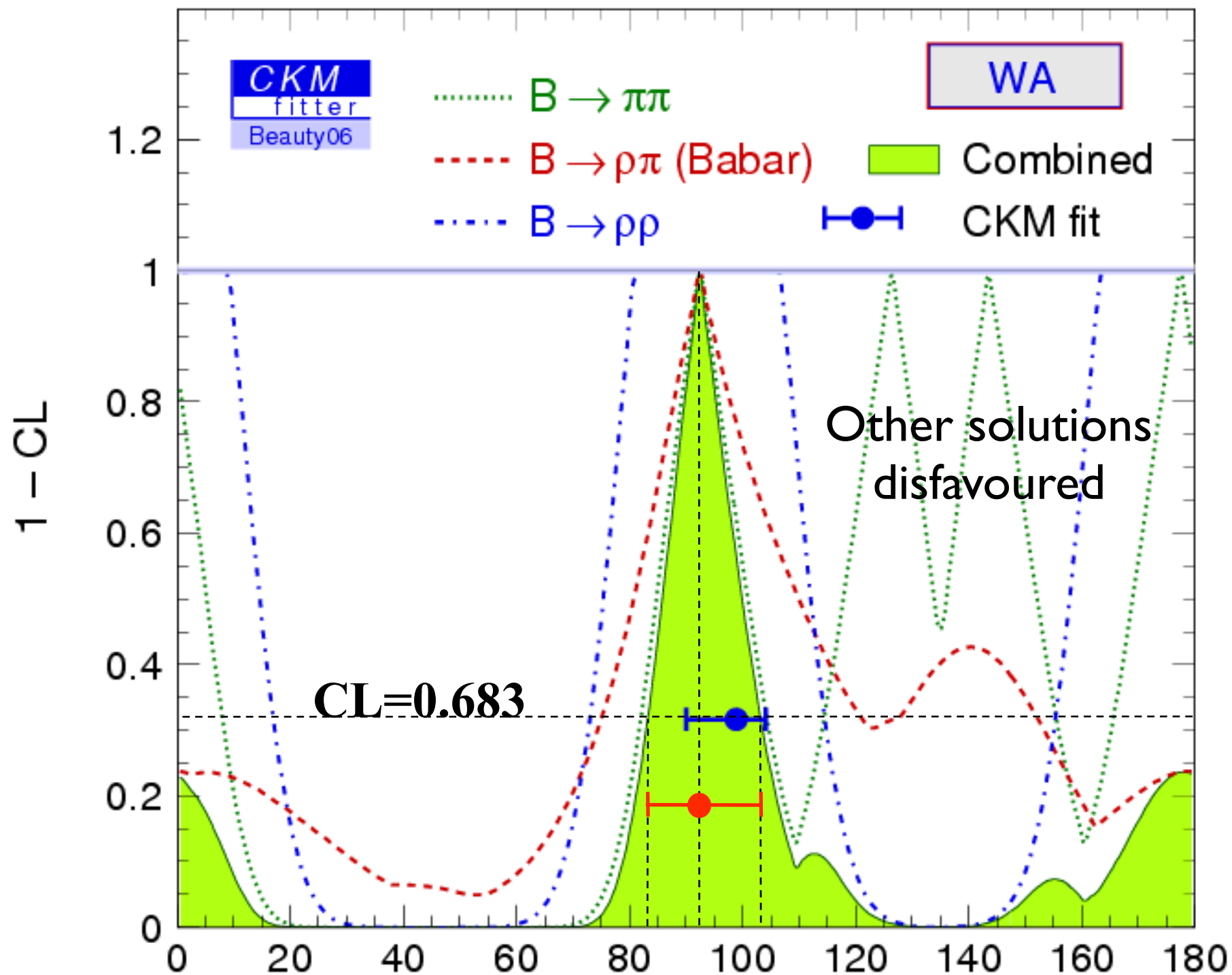


Including $\rho^+ \rho^0, \rho^0 \rho^0$ isospin analysis



Small CP violation: $\alpha_{eff} \sim 90^\circ$ or $0/180^\circ$

ϕ_2 combined (including $\rho\pi$)



Indirect: $\phi_2 = 100^{+5}_{-7}(\text{deg})$ Combined: $\phi_2 = 93^{+11}_{-9}(\text{deg})$

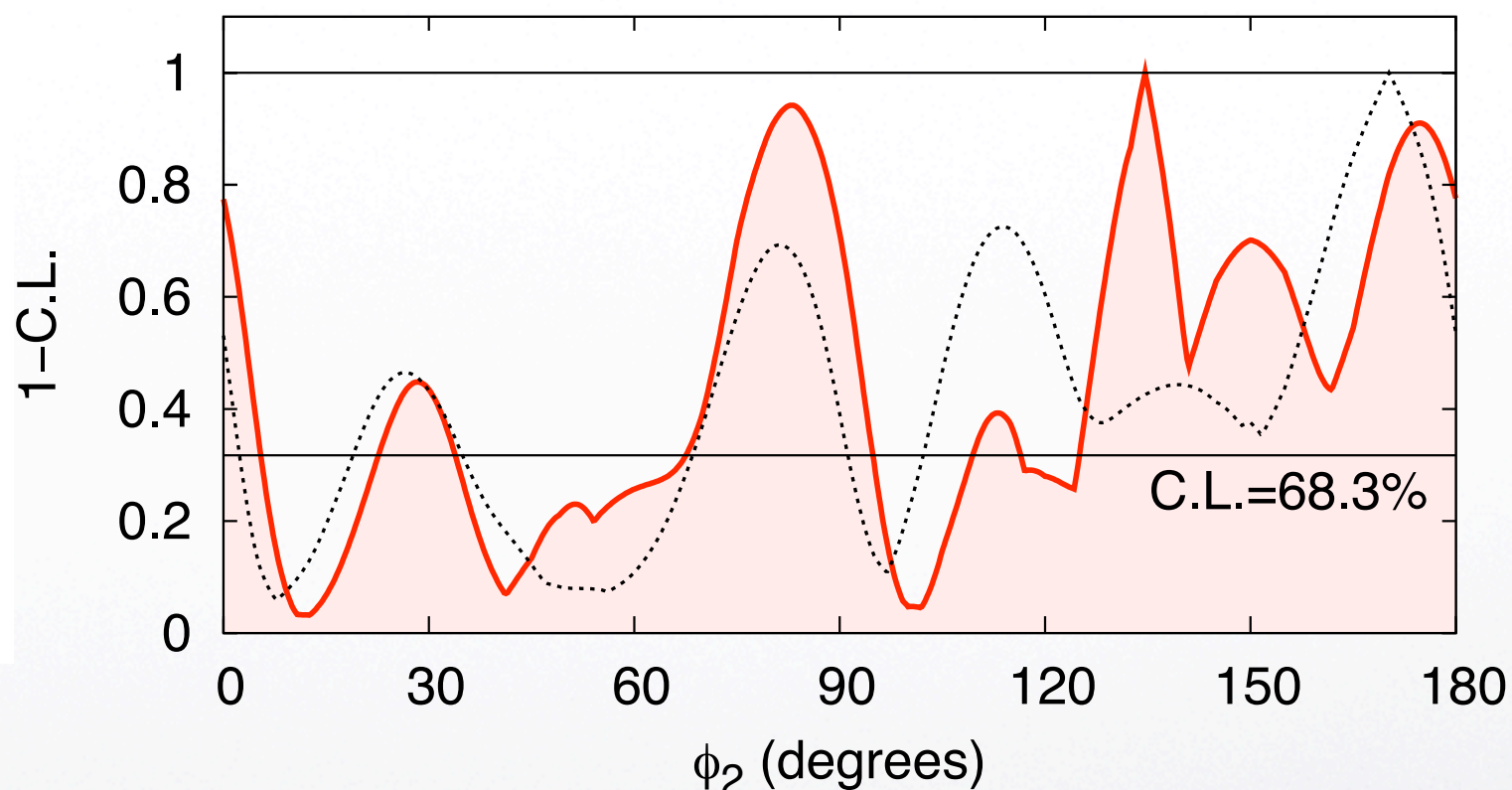


$B^0 \rightarrow (\rho\pi)^0$ from Belle

New!

hep-ex/0701015

$$\begin{aligned} |\rho^+\pi^0\rangle &= \sqrt{\frac{1}{2}} |2,1\rangle + \sqrt{\frac{1}{2}} |1,1\rangle, \\ |\rho^0\pi^+\rangle &= \sqrt{\frac{1}{2}} |2,1\rangle - \sqrt{\frac{1}{2}} |1,1\rangle, \\ |\rho^+\pi^-\rangle &= \sqrt{\frac{1}{6}} |2,0\rangle + \sqrt{\frac{1}{2}} |1,0\rangle + \sqrt{\frac{1}{3}} |0,0\rangle, \\ |\rho^-\pi^+\rangle &= \sqrt{\frac{1}{6}} |2,0\rangle - \sqrt{\frac{1}{2}} |1,0\rangle + \sqrt{\frac{1}{3}} |0,0\rangle, \\ |\rho^0\pi^0\rangle &= \sqrt{\frac{2}{3}} |2,0\rangle - \sqrt{\frac{1}{3}} |0,0\rangle. \end{aligned}$$

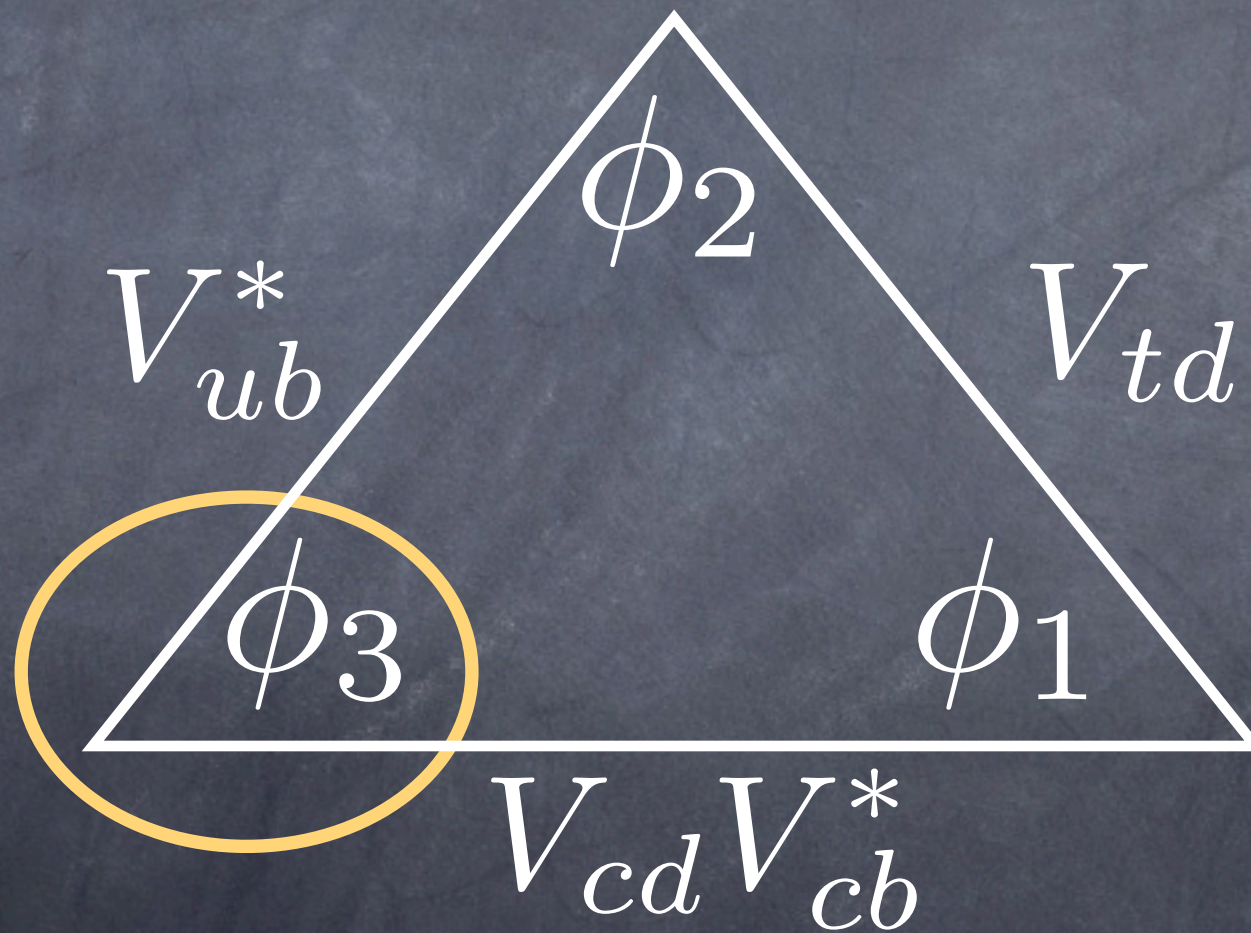


- Time-dependent Dalitz + isospin pentagon (1st time!)
- subm'ed to PRL & to be included in the world average

Other angles?

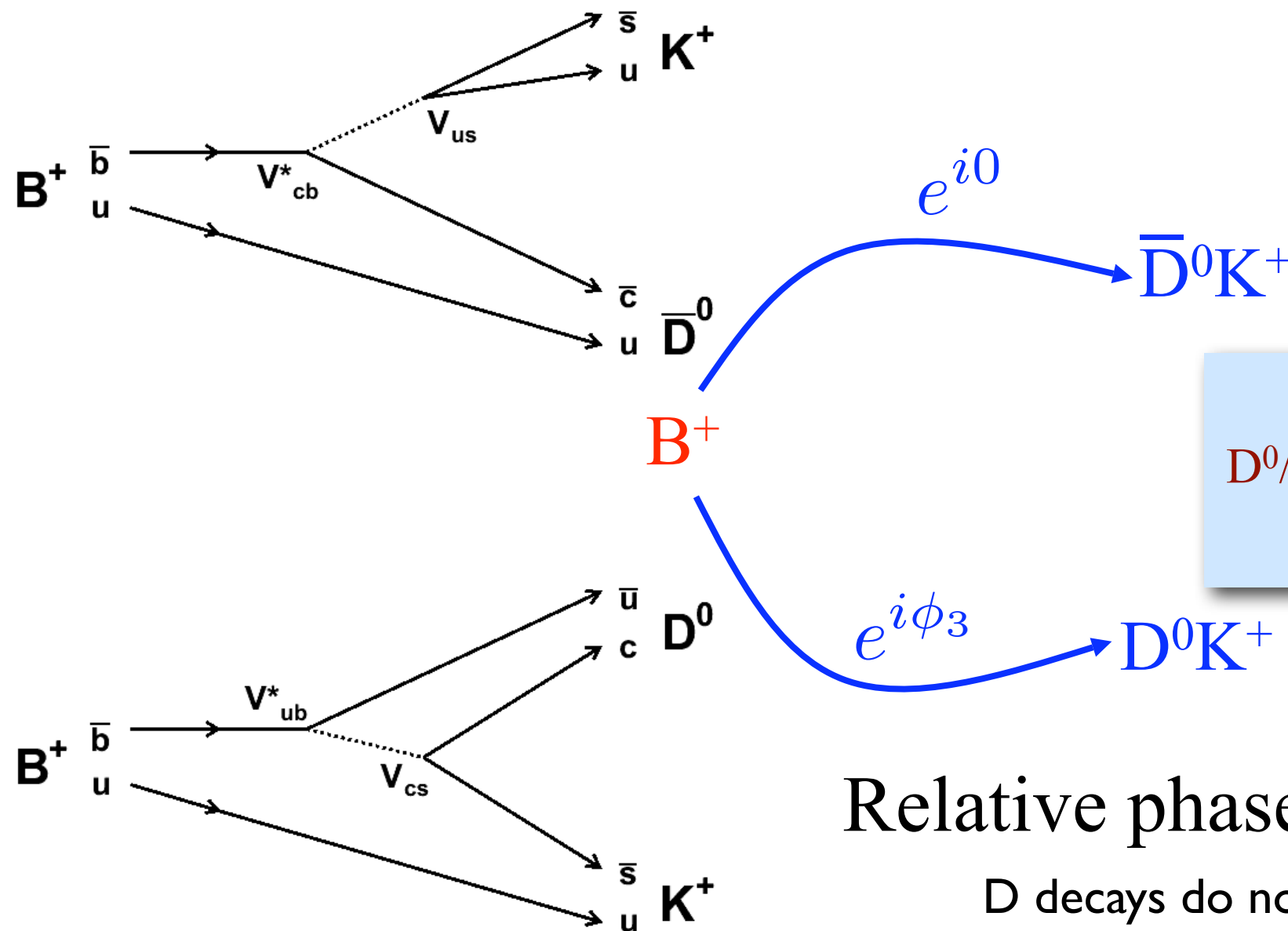
Unitarity triangle angles

BABAR:	β	α	γ
BELLE:	ϕ_1	ϕ_2	ϕ_3
	易	難	魔



ϕ_3 from CPV in DK modes

e.g. $B^+ \rightarrow D^0/\bar{D}^0 K^+$



GLW: Gronau, London, Wyler (2001)
 ADS: Atwood, Dunietz, Soni (1997)
 GGSZ: Giri, Grossman, Soffer, Zupan (2003)

\rightarrow CP state (GLW)
 $D^0/\bar{D}^0 \rightarrow K^-\pi^+/K^+\pi^-$, CA/DCS (ADS)
 $\rightarrow K_S\pi^+\pi^-$, Dalitz (GGSZ)

Relative phase = $e^{-i\phi_3}$

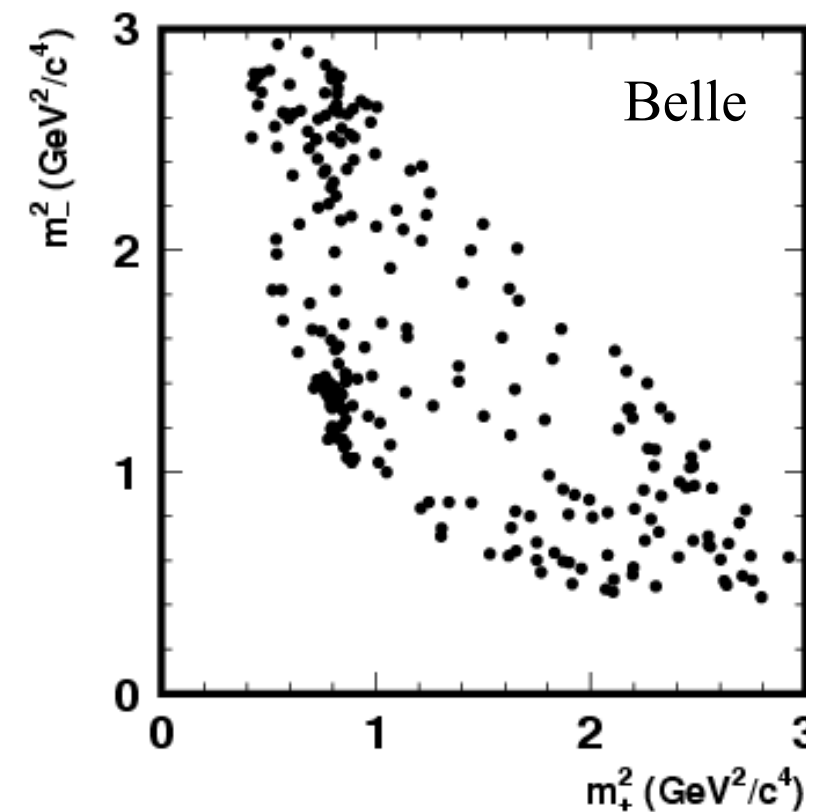
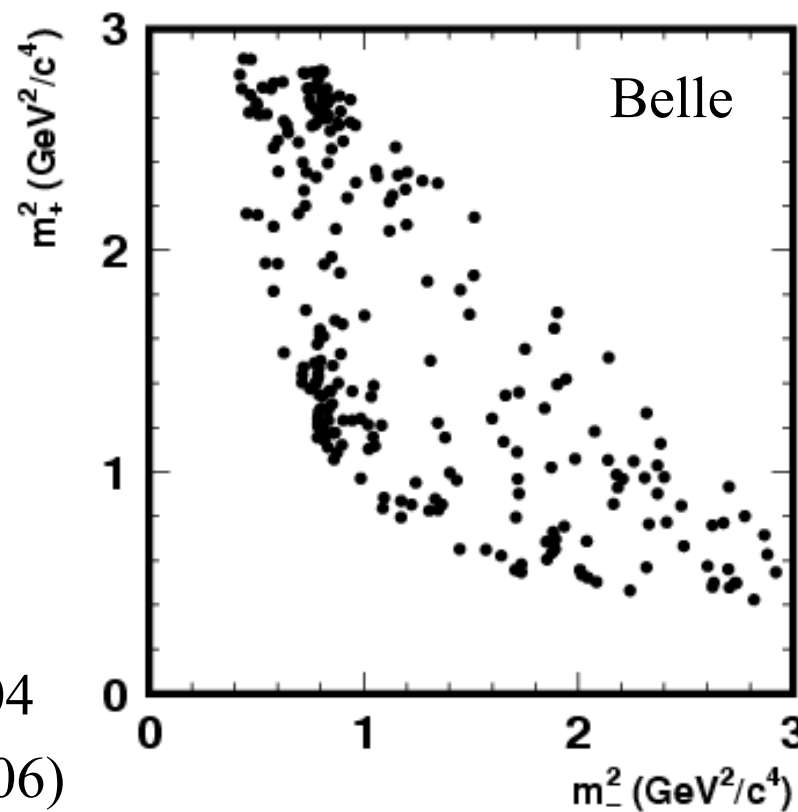
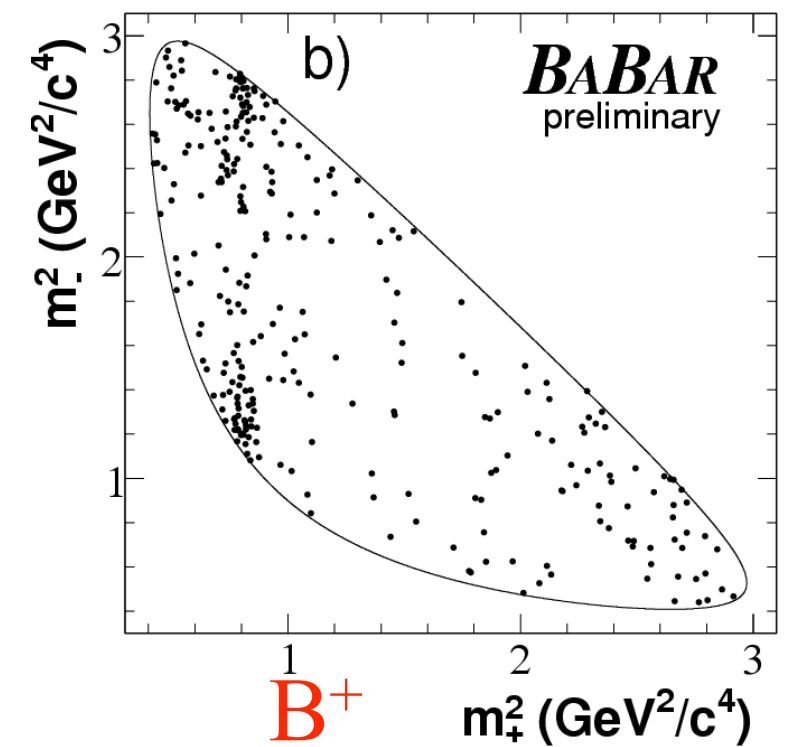
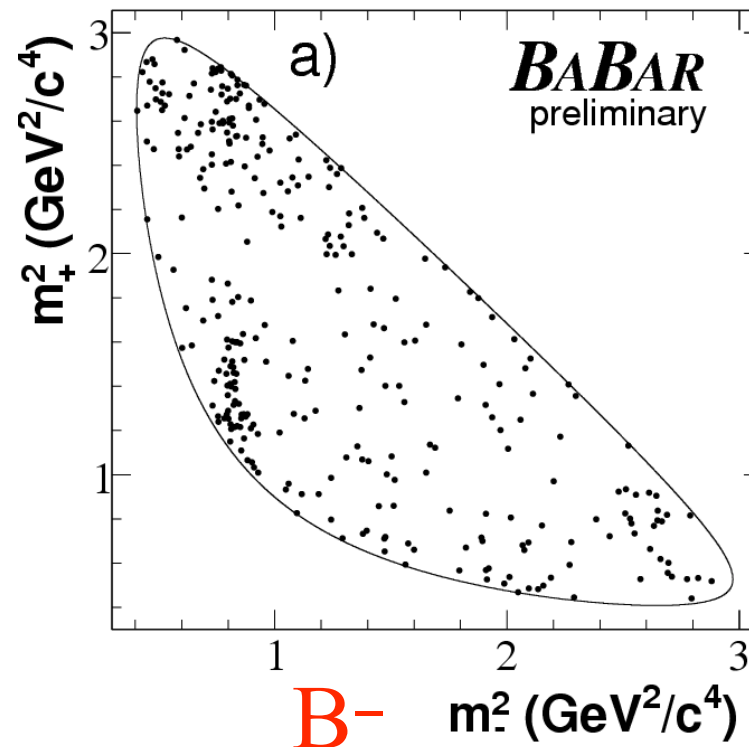
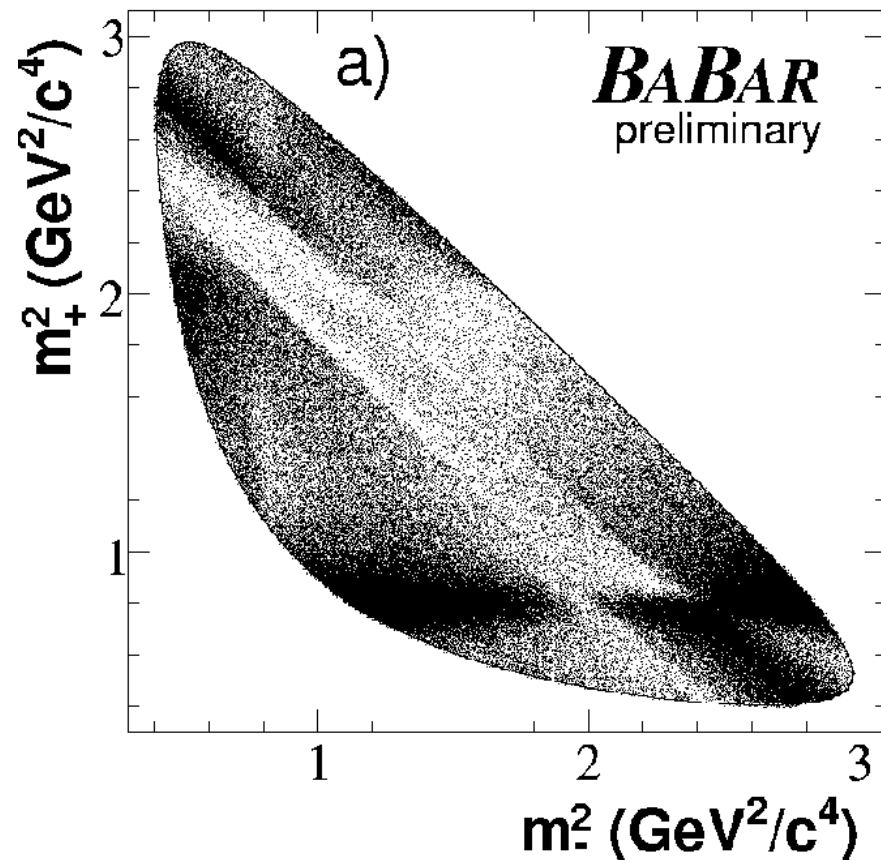
D decays do not involve V_{ub} or V_{td}
 \rightarrow no contribution to phase

$B^\pm \rightarrow$ no mixing, no t -dependence

The GGSZ method

Look for deviations in $B^\pm \rightarrow D^0 K^\pm$ plots

Map out Dalitz plot
from **all** D^0 decays



BaBar preliminary; hep-ex/0607104
Belle; Phys.Rev.D 73, 172009 (2006)

Results from GGSZ method

Express in terms of
measurables from B^\pm

$$x_\pm = r_B \cos(\delta_B \pm \phi_3)$$

$$y_\pm = r_B \sin(\delta_B \pm \phi_3)$$

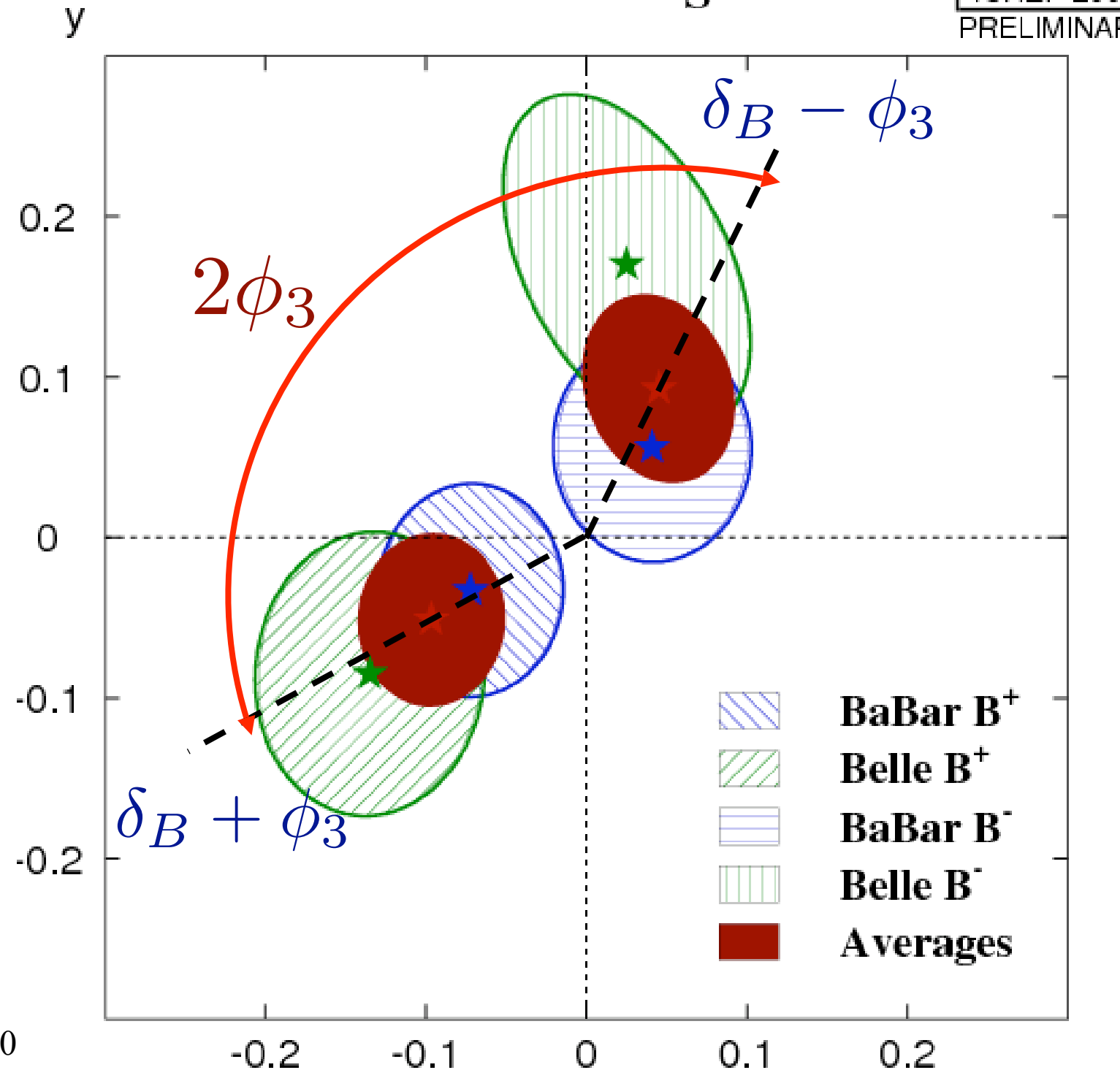
r_B : ratio of D/\bar{D} ampl.
 $= 0.16 \pm 0.07$

δ_B : D/\bar{D} relative phase

Different r_B , δ_B for each
mode $D^{(*)}K^{(*)}$

$DK^+, D \rightarrow K_S \pi^+ \pi^-$

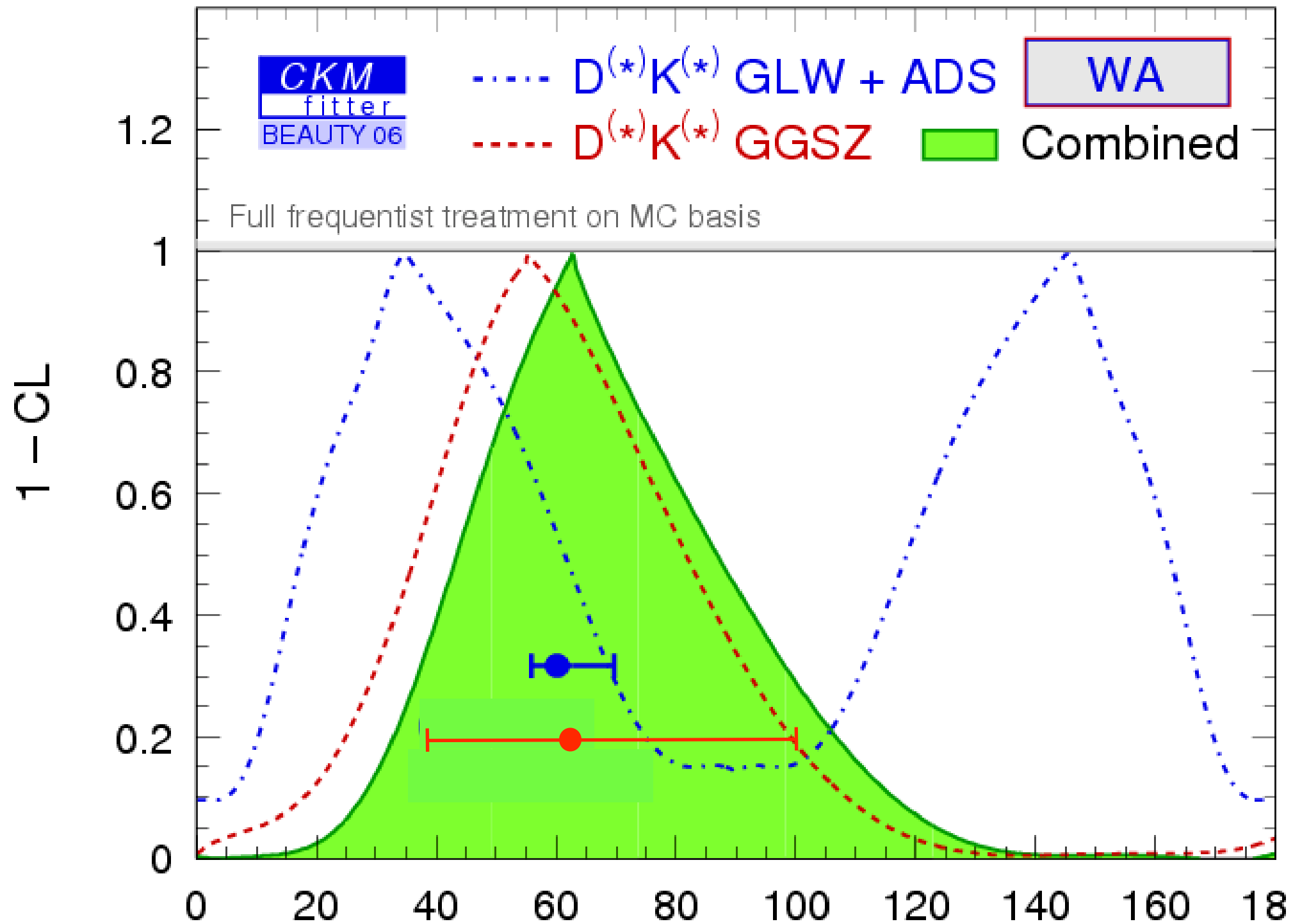
HFAG
ICHEP 2006
PRELIMINARY



BaBar preliminary; hep-ex/060710
Belle; PRD 73, 112009 (2006)

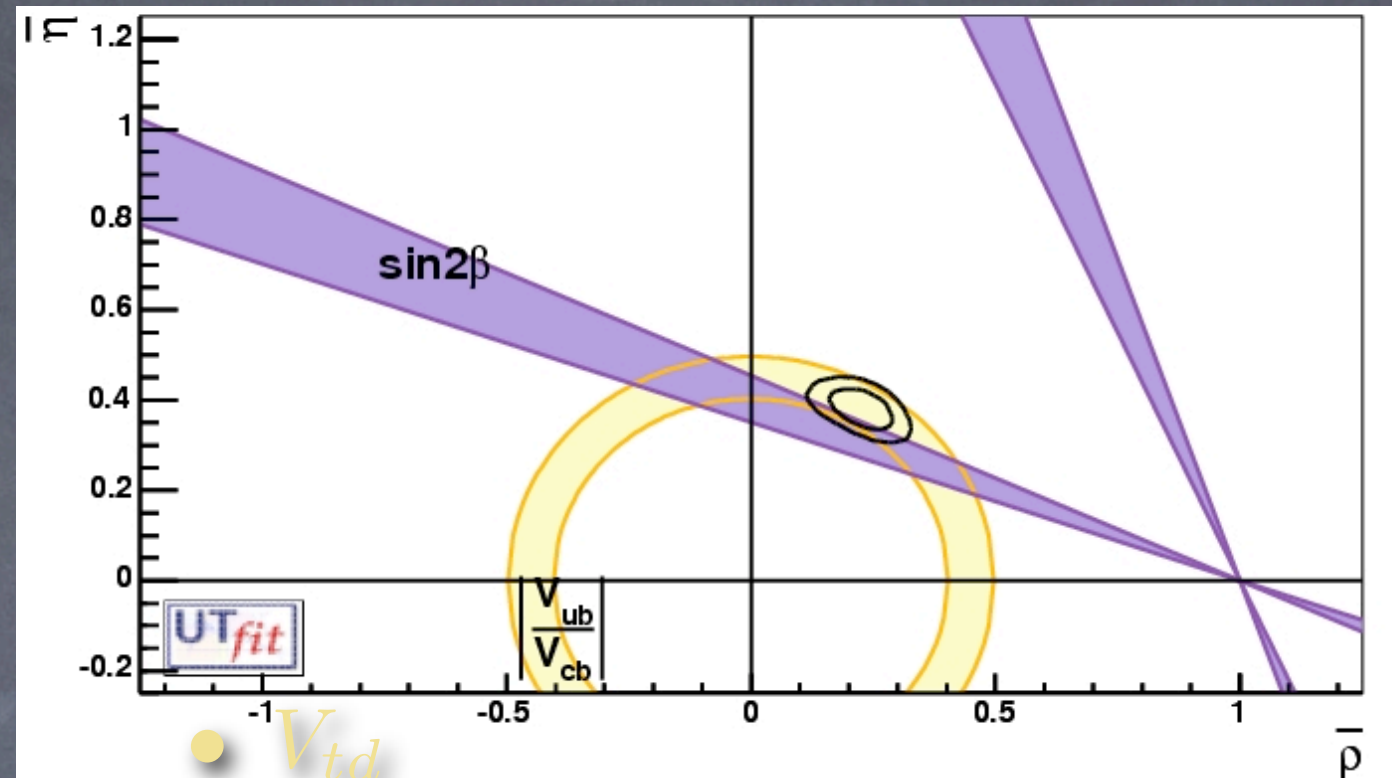
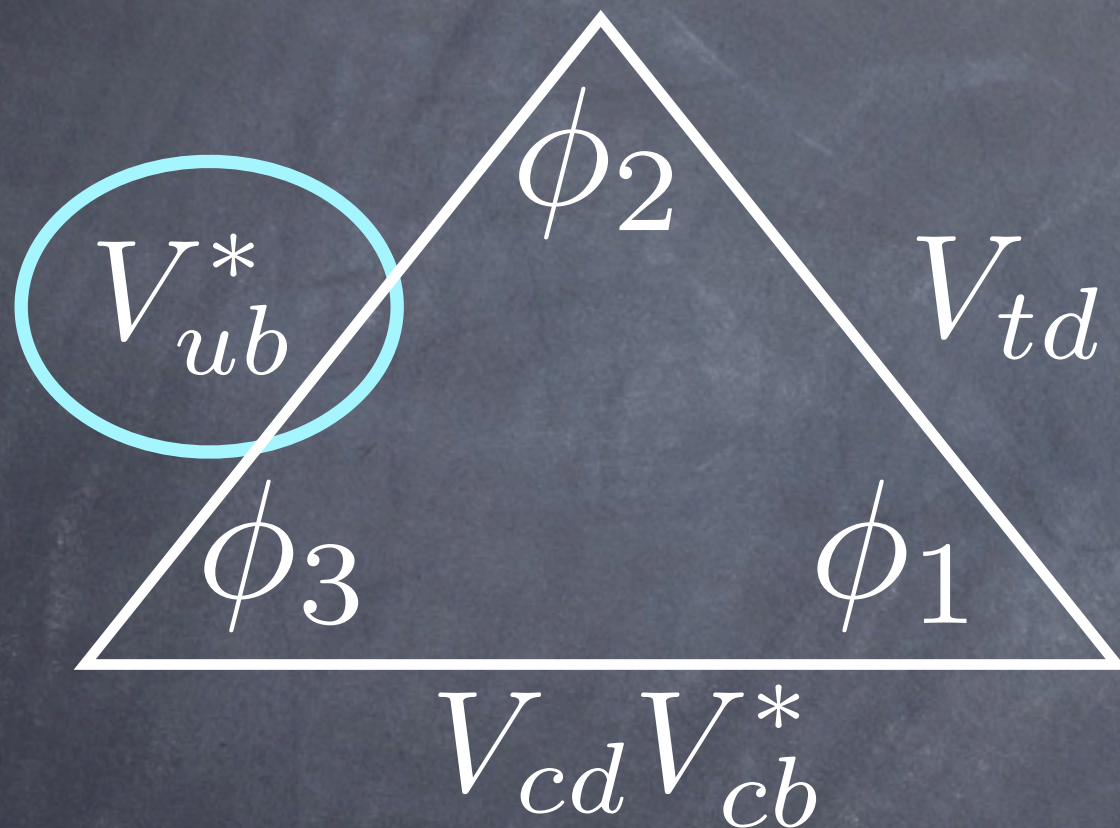
Contours give $-2\Delta(\ln L) = \Delta\chi^2 = 1$, corresponding to 60.7% CL for 2 dof

Overall result



Indirect: $\phi_3 = 59^{+9}_{-4}(\text{deg})$ γ (deg) Combined: $\phi_3 = 62^{+38}_{-24}(\text{deg})$

The Sides



- $B \rightarrow X_d \gamma$ (A. Bevan)
- B_s mixing (K. Pitts)

• V_{cb}

- $\mathcal{O}(1\%)$ precision

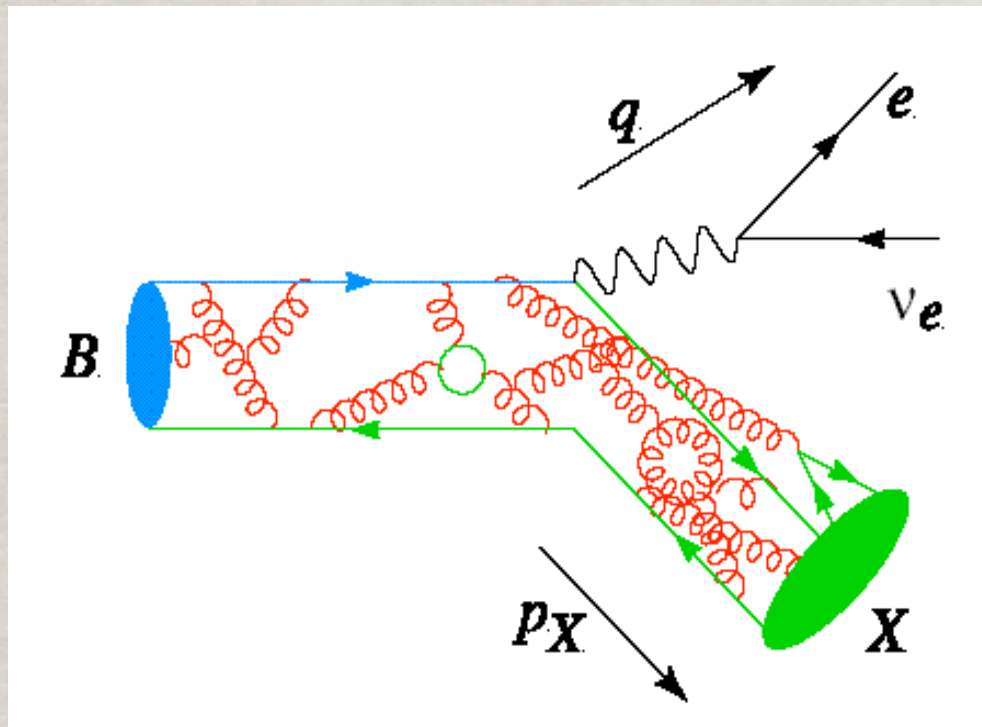
■ $V = |V| \exp(i\phi)$

- $|V|$ from semi-leptonic decay rates
- ϕ from CP asymmetries

just overly simplified guidelines

$$\Gamma_{X\ell\nu} \propto |V_{ij}|^2$$

Semileptonic B Decays

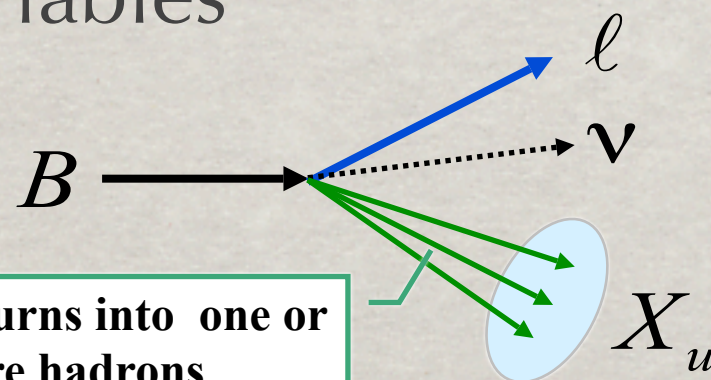


- * Semileptonic B decays allow measurement of $|V_{cb}|$ and $|V_{ub}|$ from tree level processes.
- * Presence of a single hadronic current allows control of theoretical uncertainties.

$$|V_{cb}| \gg |V_{ub}|$$

$$\frac{\Gamma(b \rightarrow u\ell^-\bar{\nu})}{\Gamma(b \rightarrow c\ell^-\bar{\nu})} \approx \frac{|V_{ub}|^2}{|V_{cb}|^2} \approx \frac{1}{50}$$

kinematic variables
for $B \rightarrow X\ell\nu$

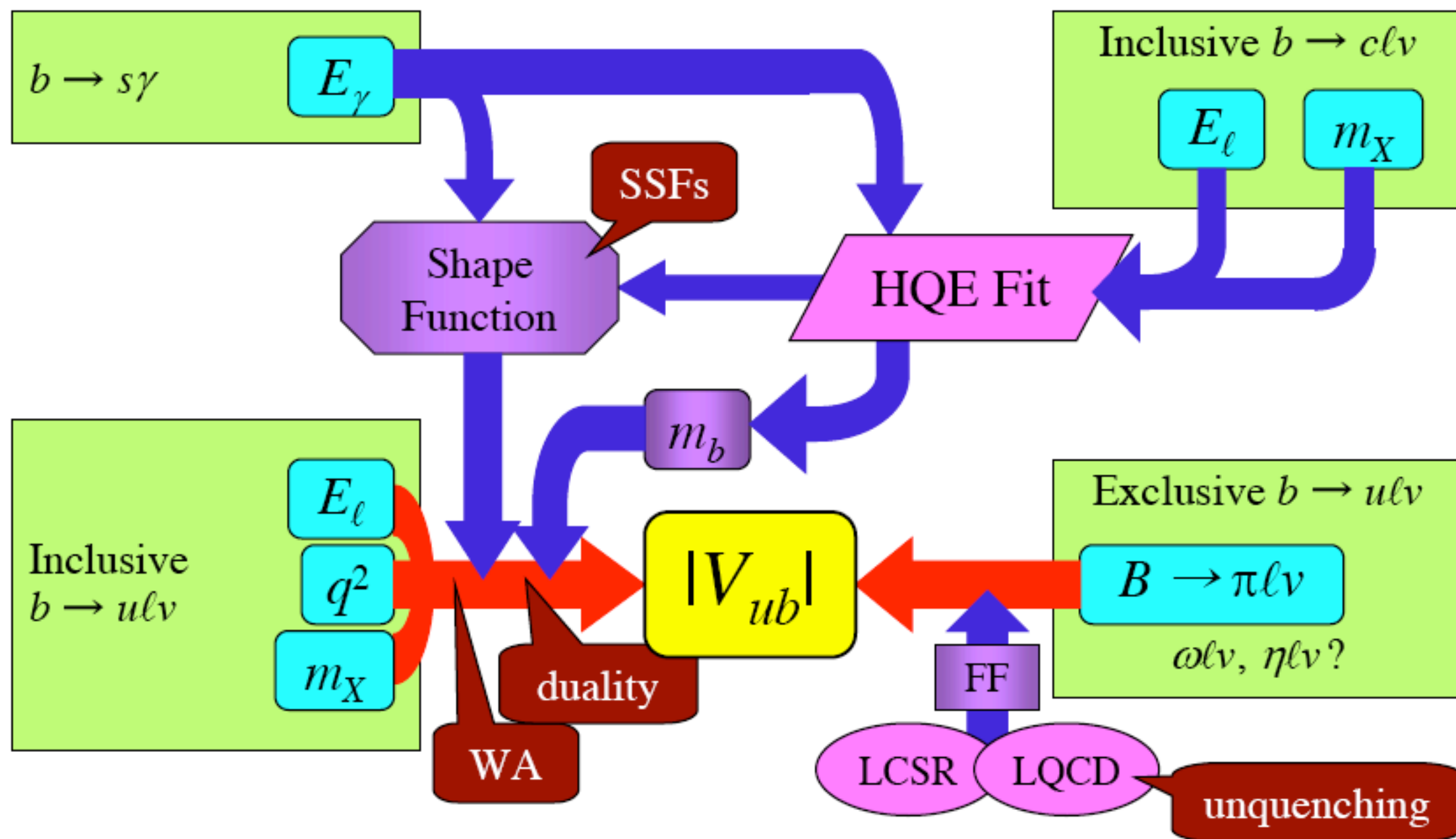


E_ℓ = lepton energy

$$q^2 = (p_\ell + p_\nu)^2$$

m_X = mass of the hadronic part

Morii's chart for V_{ub}



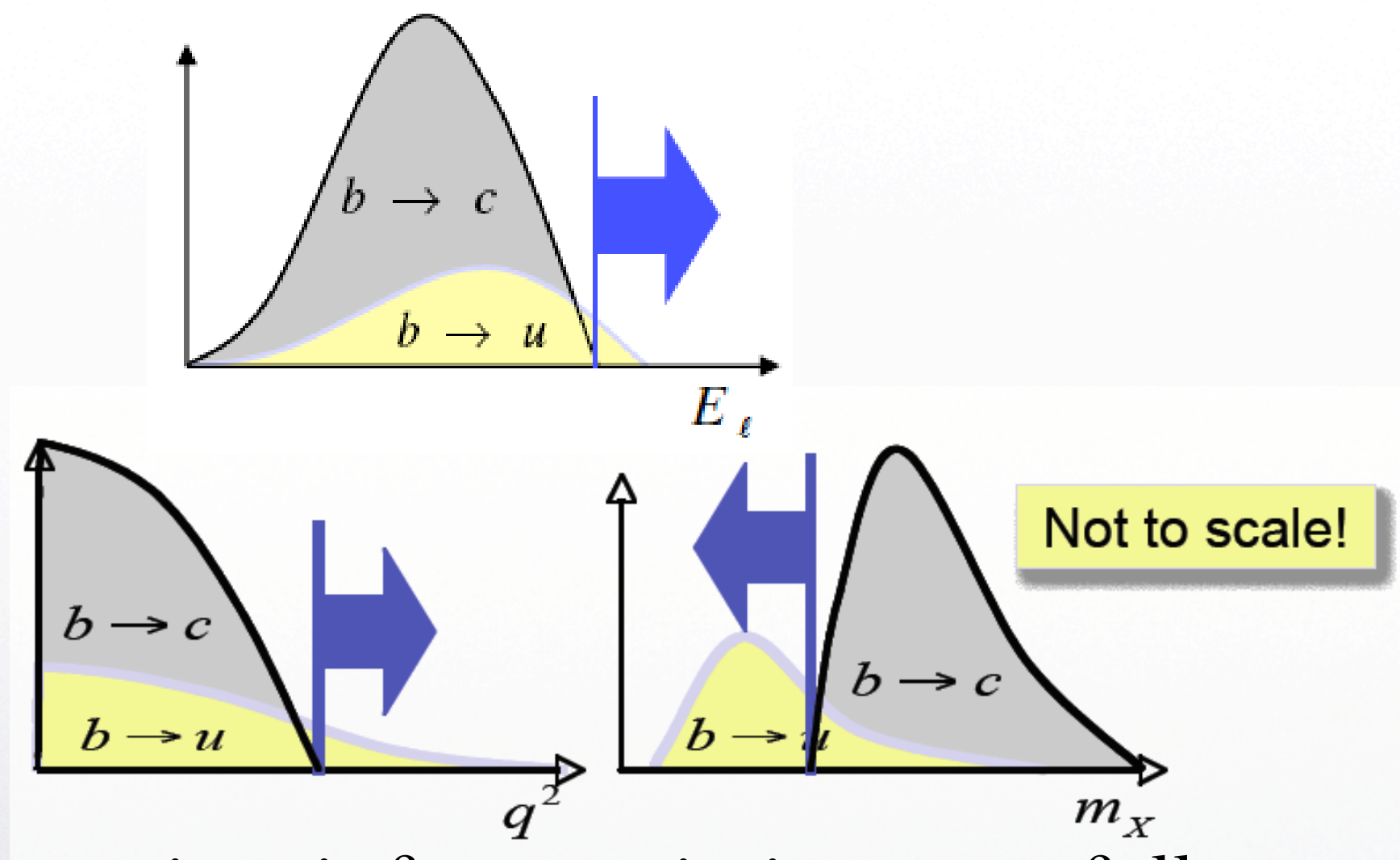
17 January 2006

M. Morii, Harvard

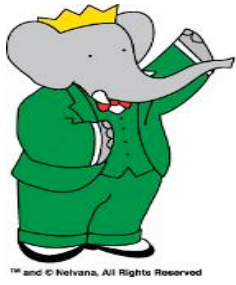
36



Inclusive $B \rightarrow X_u \ell \nu$



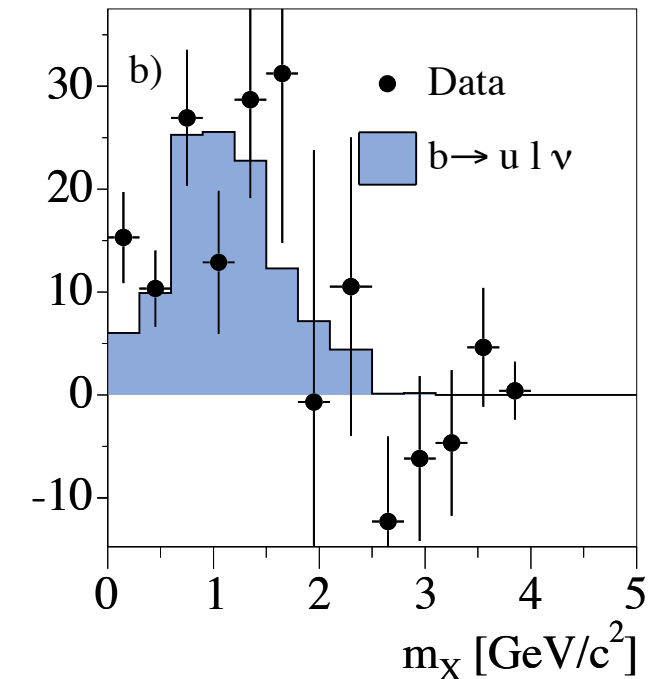
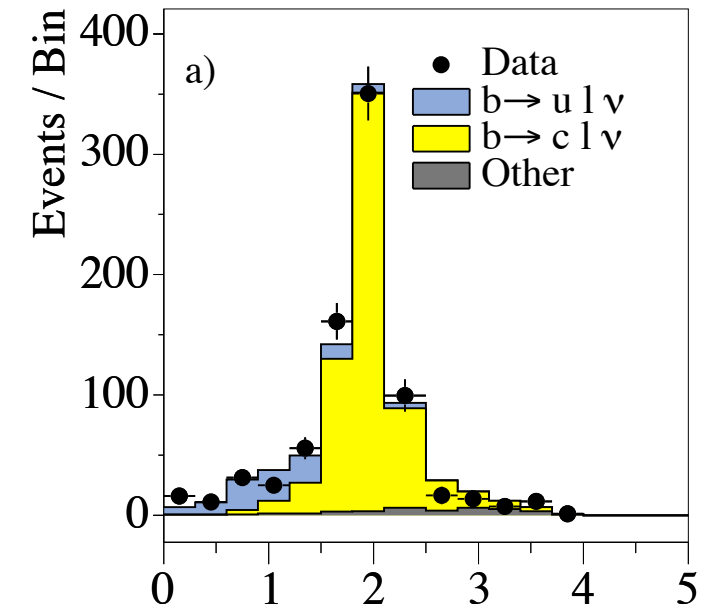
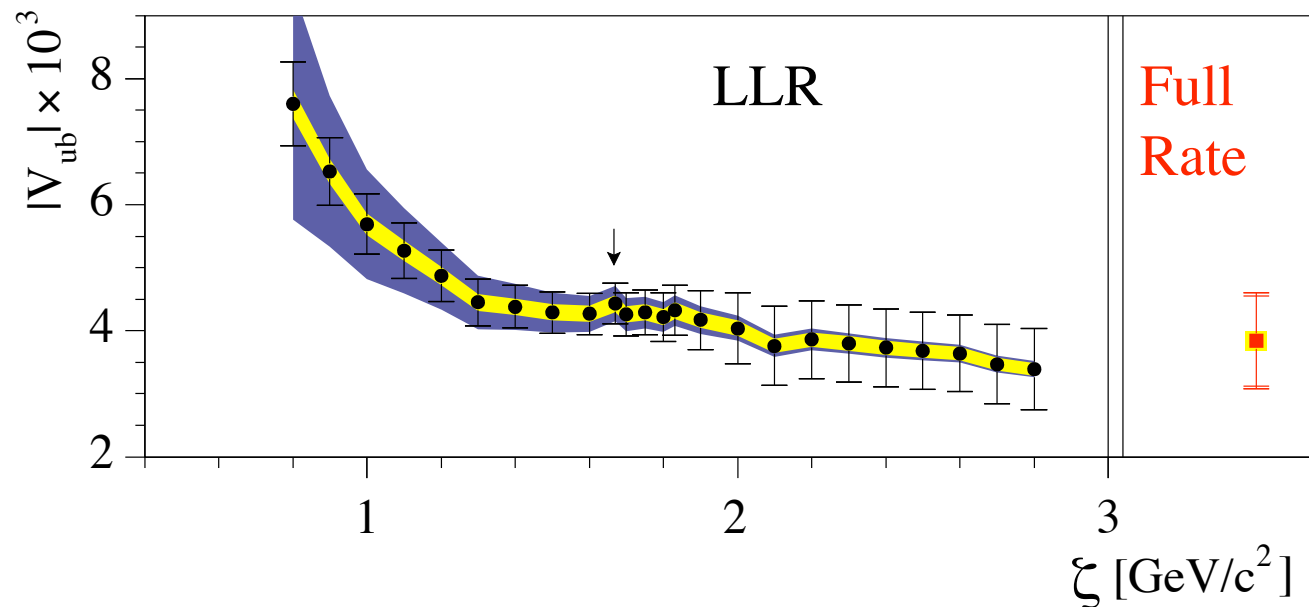
q^2 and M_X requires info. on missing ν --> full-recon.



V_{ub} (LLR method)

$$\Gamma(B \rightarrow X_u \ell \nu) = \frac{|V_{ub}|^2}{|V_{ts}|^2} \int W(E_\gamma) \frac{d\Gamma(B \rightarrow X_s \gamma)}{dE_\gamma} dE_\gamma$$

- m_{X_u} ($B \rightarrow X_u \ell \nu$) and E_γ ($B \rightarrow X_s \gamma$)
- two methods
 - ★ $m_{X_u} < \zeta (< 1.67 \text{ GeV})$ (LLR)
 - ★ m_{X_u} in full range (U, HLM)



m_X acceptance

LLR : $M_X < 1.67 \text{ GeV}$: $|V_{ub}| = (4.43 \pm 0.38_{\text{stat}} \pm 0.25_{\text{syst}} \pm 0.29_{\text{theo}}) 10^{-3}$ **12%** **72%**

OPE: $M_X < 2.50 \text{ GeV}$: $|V_{ub}| = (3.84 \pm 0.70_{\text{stat}} \pm 0.30_{\text{syst}} \pm 0.10_{\text{theo}}) 10^{-3}$ **20%** **98%**

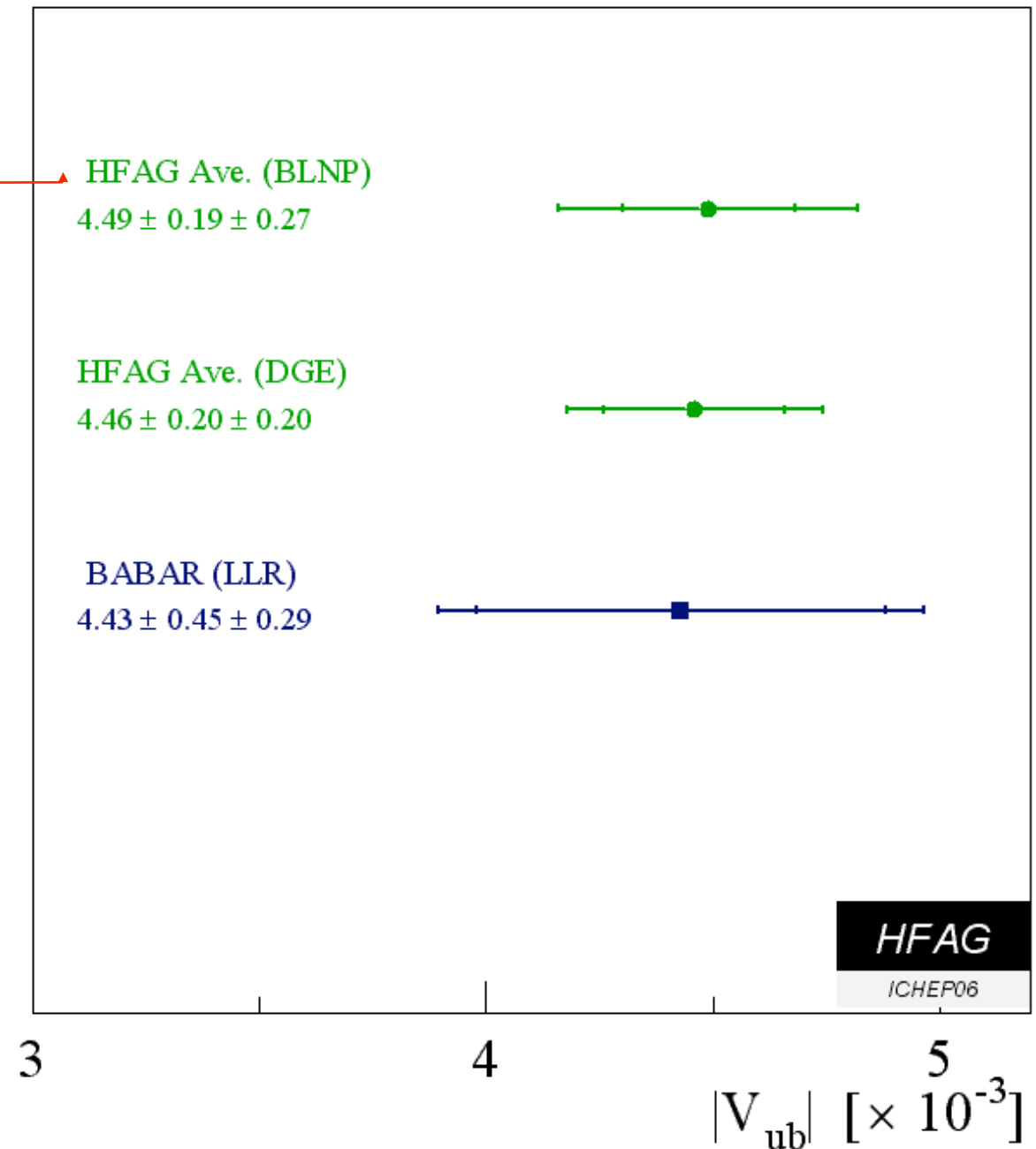
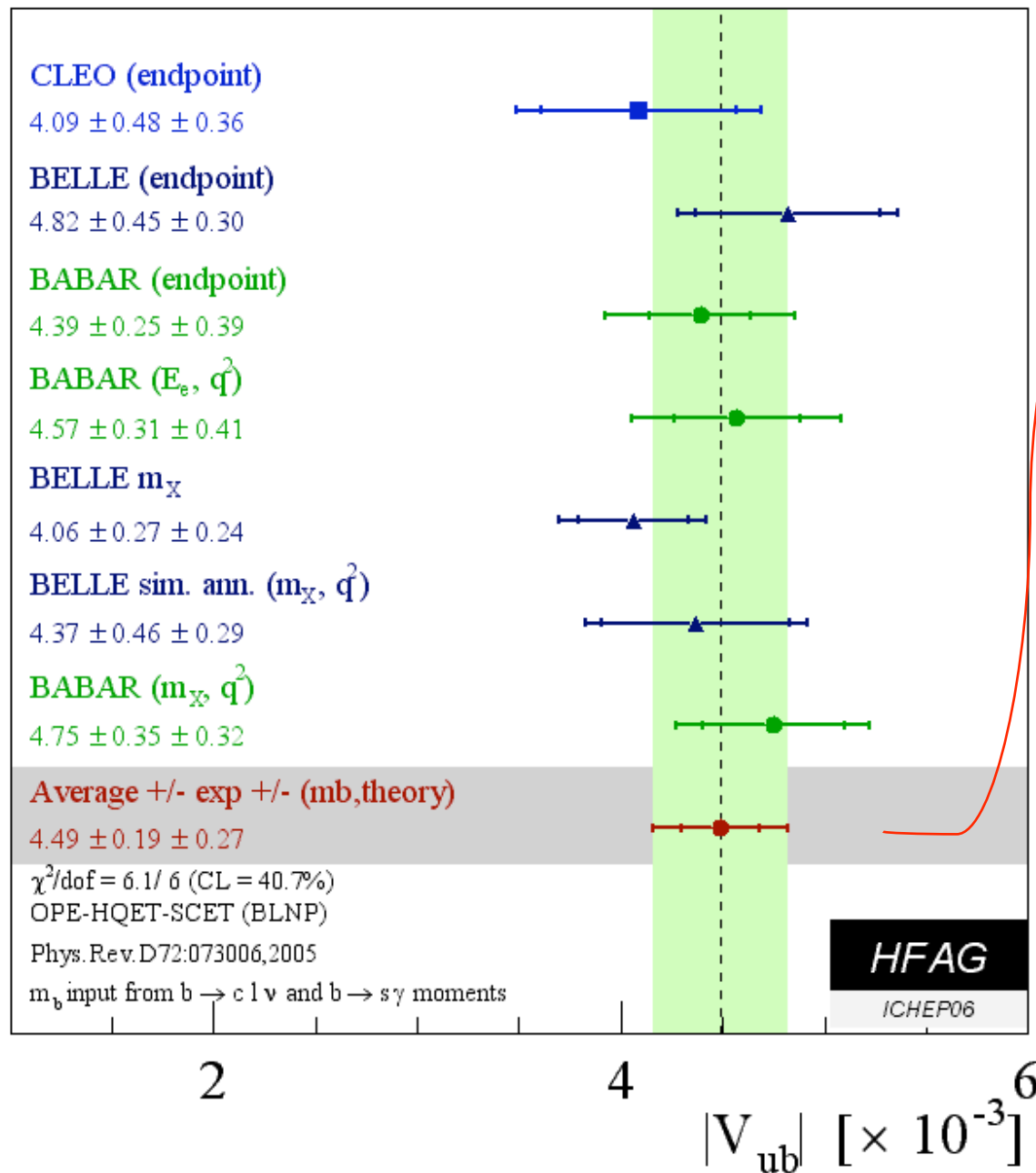
V_{ub} inclusive summary

BLNP: Bosch, Lange, Neubert, Paz (2005)

DGE: Anderson, Gardi (2006)

LLR: Leibovich, Low, Rothstein (2006)

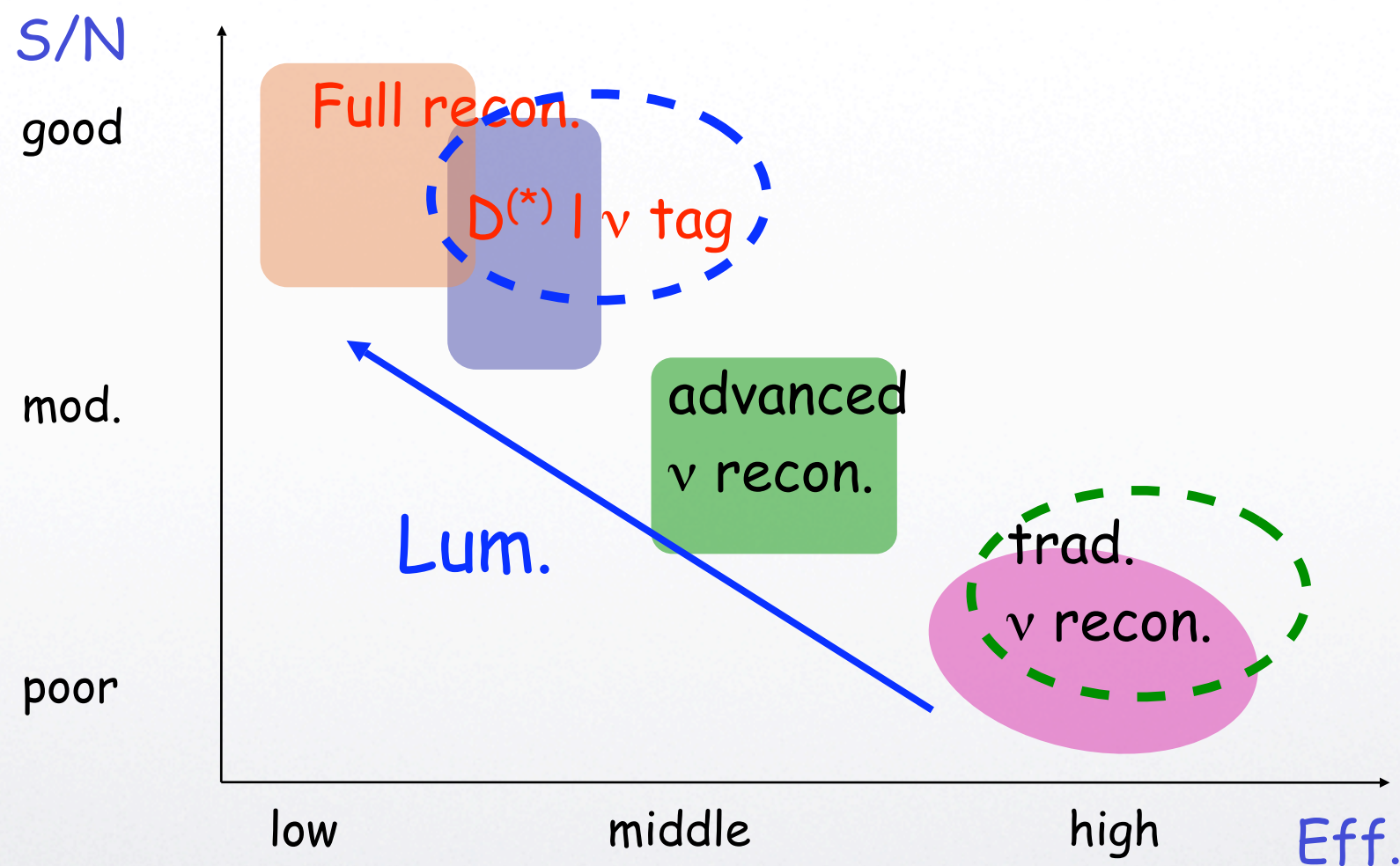
Representative theory example (BLNP)



Room for some experimental statistical improvement

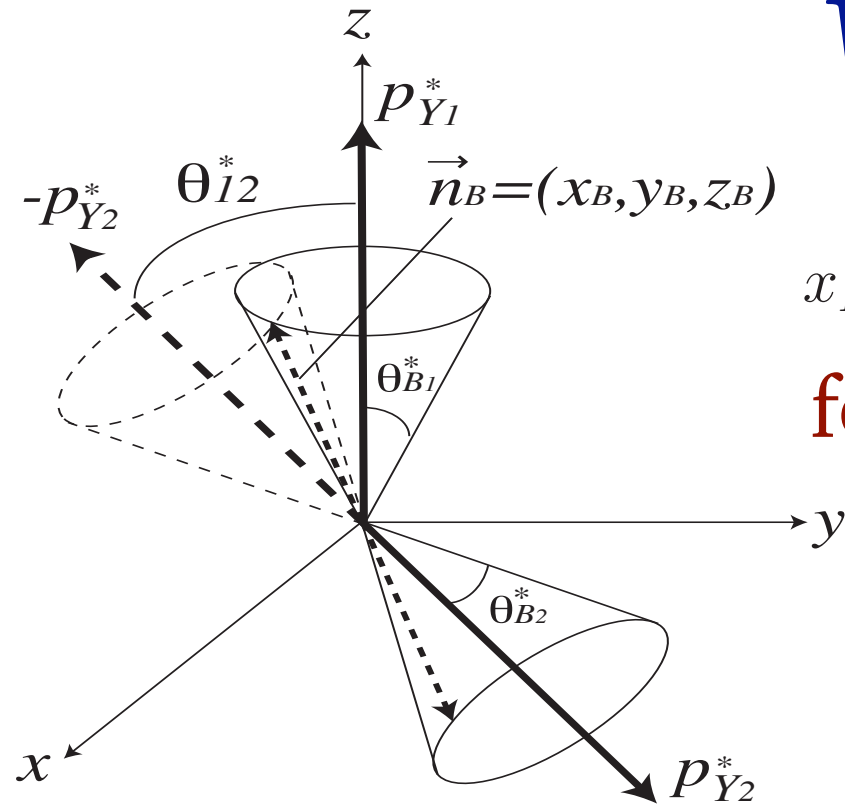


Exclusive $B \rightarrow X_u \ell \nu$



How well can we measure the q^2 dist. for $B \rightarrow X_u \ell \nu$?

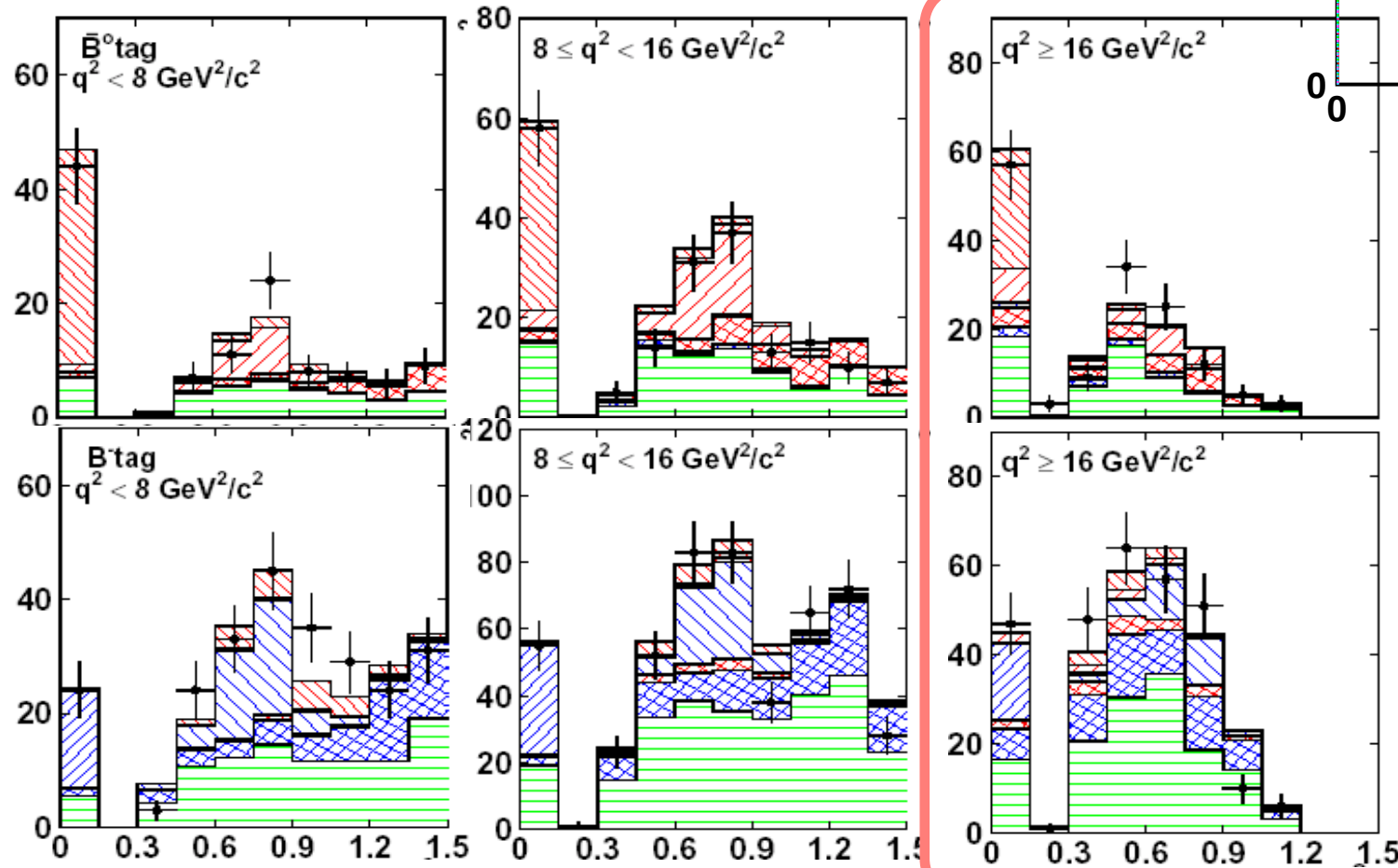
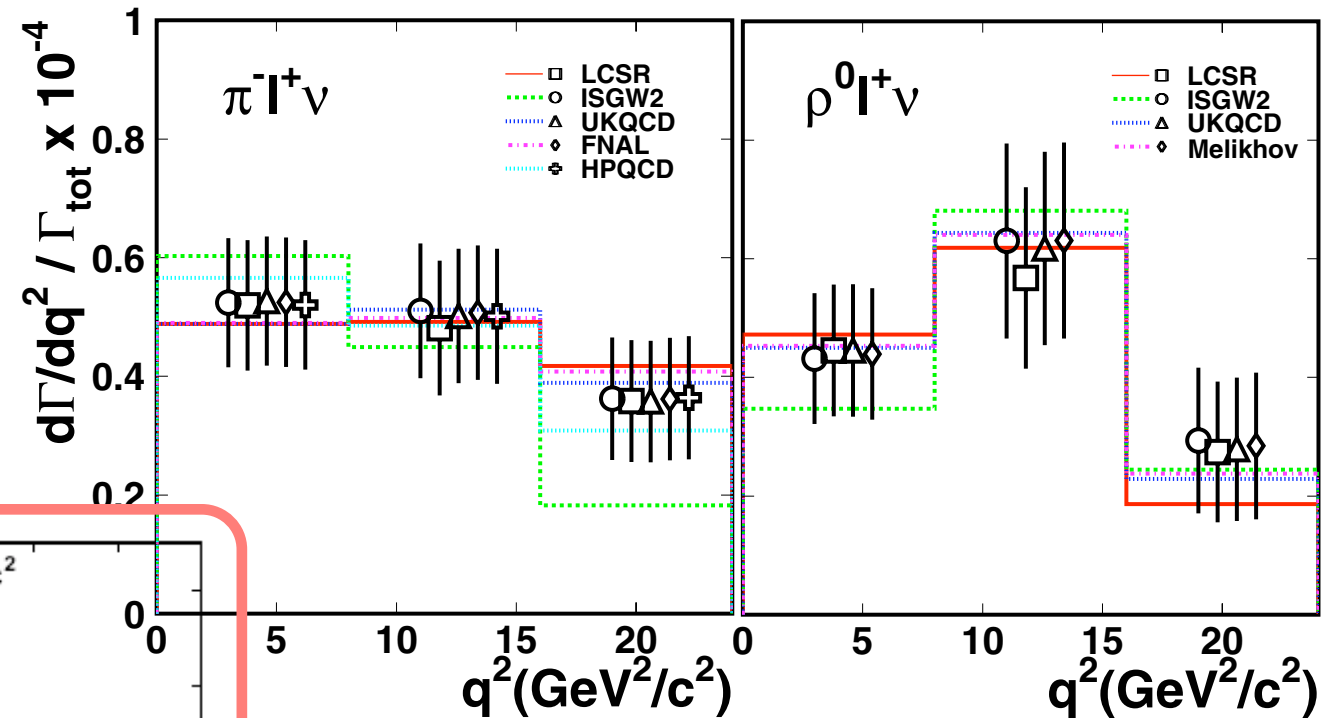
V_{ub} by $D^* \ell \nu$ tagging



$$x_B^2 = 1 - \frac{1}{\sin^2 \theta_{12}^*} (\cos^2 \theta_{B1}^* + \cos^2 \theta_{B2}^* - 2 \cos \theta_{B1}^* \cos \theta_{B2}^* \cos \theta_{12}^*)$$

for true signal, $0 < x_B^2 < 1$

Fig. 1. Kinematics of the double semileptonic decay.



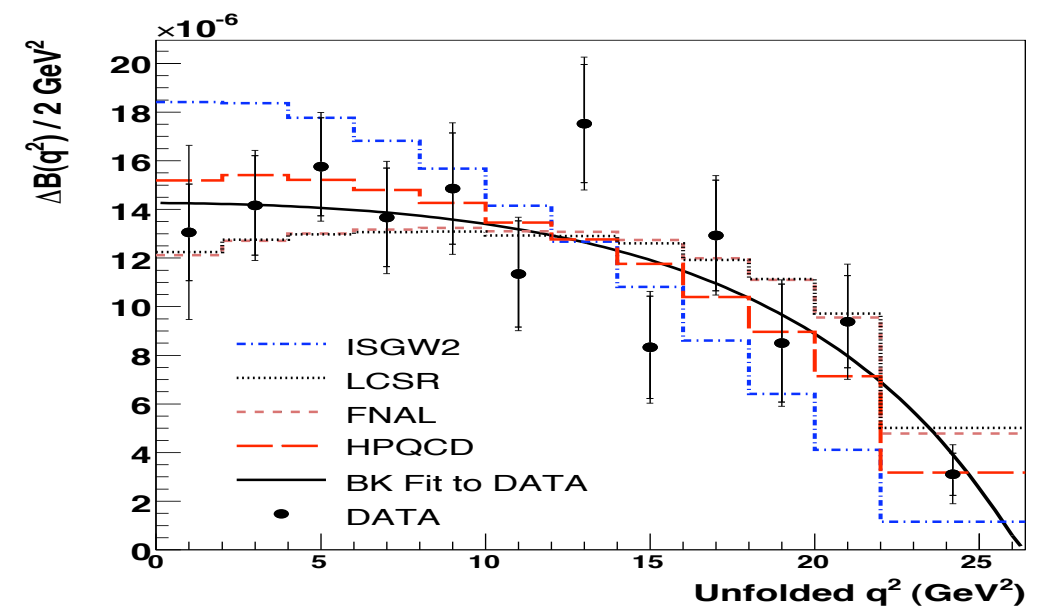
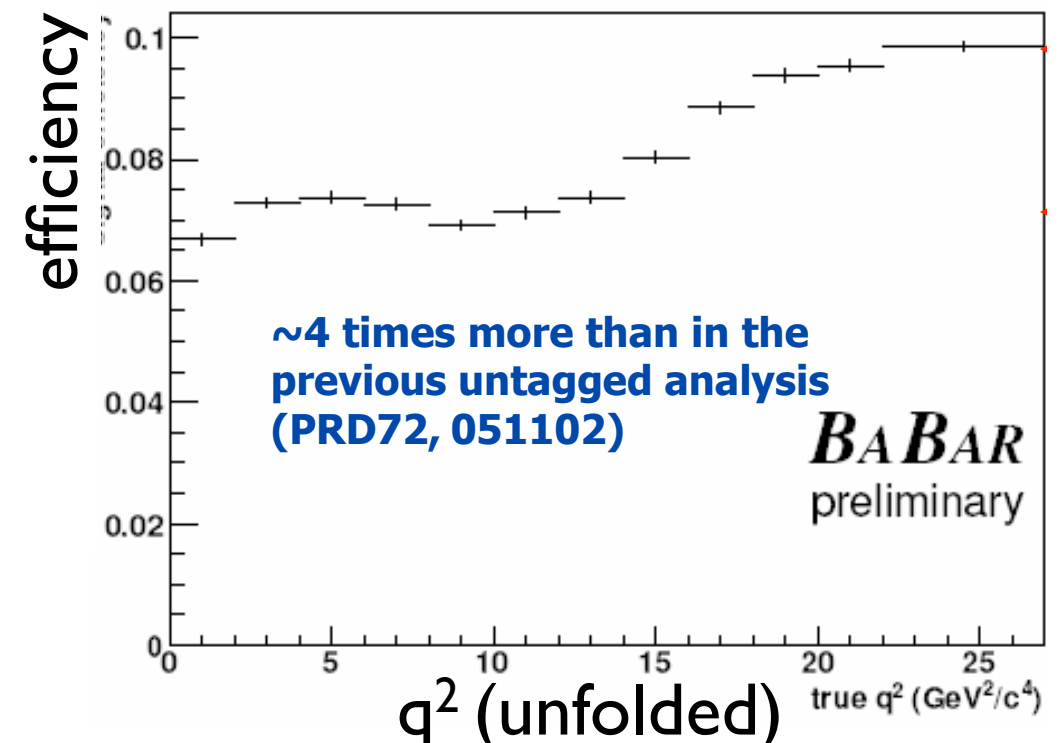
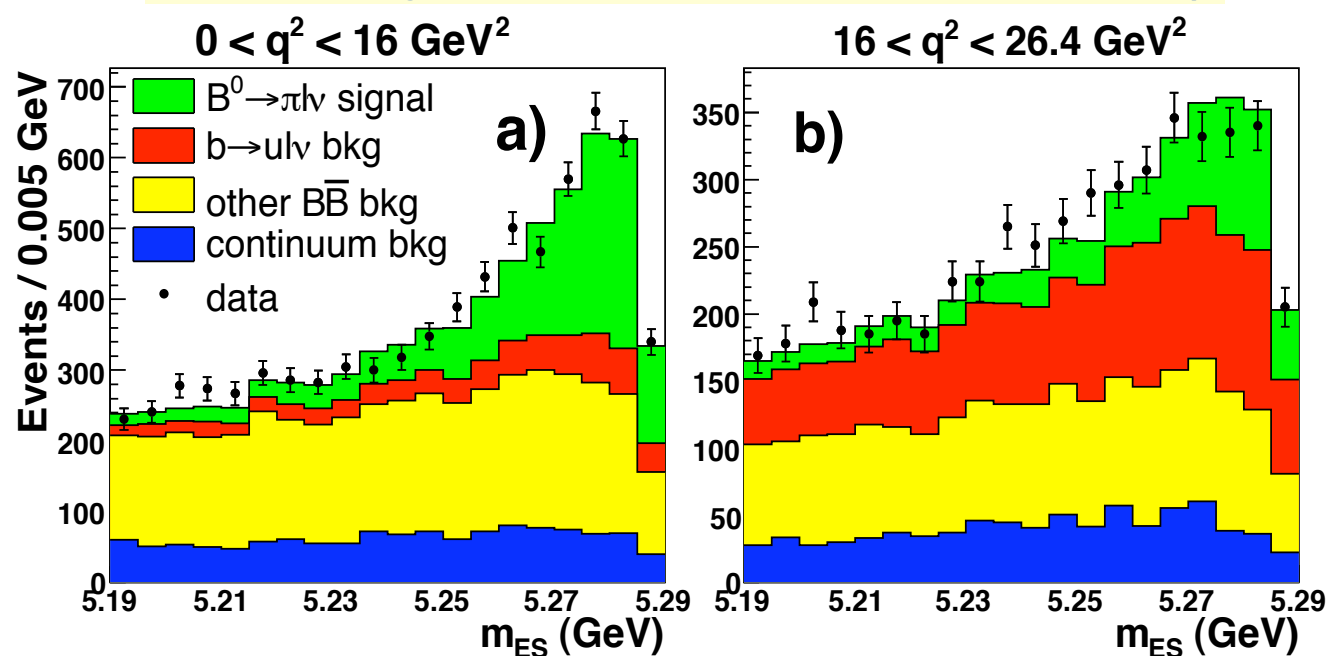
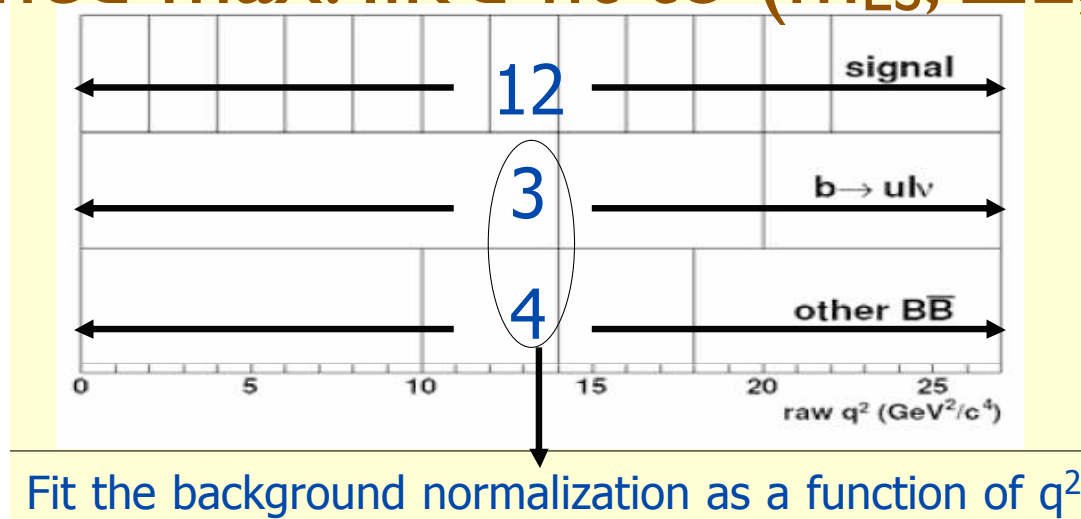
Mode	$ V_{ub} (\times 10^{-3})$
$\pi^- \ell^+ \nu$	$3.59 \pm 0.51 \pm 0.10$ FNAL
$\pi^0 \ell^+ \nu$	$3.63 \pm 0.70 \pm 0.20^{+0.63}_{-0.41}$
$\pi^- \ell^+ \nu + \pi^0 \ell^+ \nu$	$3.60 \pm 0.41 \pm 0.20^{+0.62}_{-0.41}$
$\pi^- \ell^+ \nu$	$4.02 \pm 0.57 \pm 0.22^{+0.62}_{-0.41}$ HPQCD
$\pi^0 \ell^+ \nu$	$4.06 \pm 0.78 \pm 0.22^{+0.60}_{-0.41}$
$\pi^- \ell^+ \nu + \pi^0 \ell^+ \nu$	$4.03 \pm 0.46 \pm 0.22^{+0.59}_{-0.41}$

Measurement of the $B^0 \rightarrow \pi^- \ell^+ \nu$ Form-Factor Shape and Branching Fraction, and Determination of $|V_{ub}|$ with a Loose Neutrino Reconstruction Technique



- loose requirement on $\pi^- \ell^+$
- cuts optimized as a ftn. of q^2
- eff. up by ~ 4 times

binned max. lik'd fit to $(m_{ES}, \Delta E, q^2)$

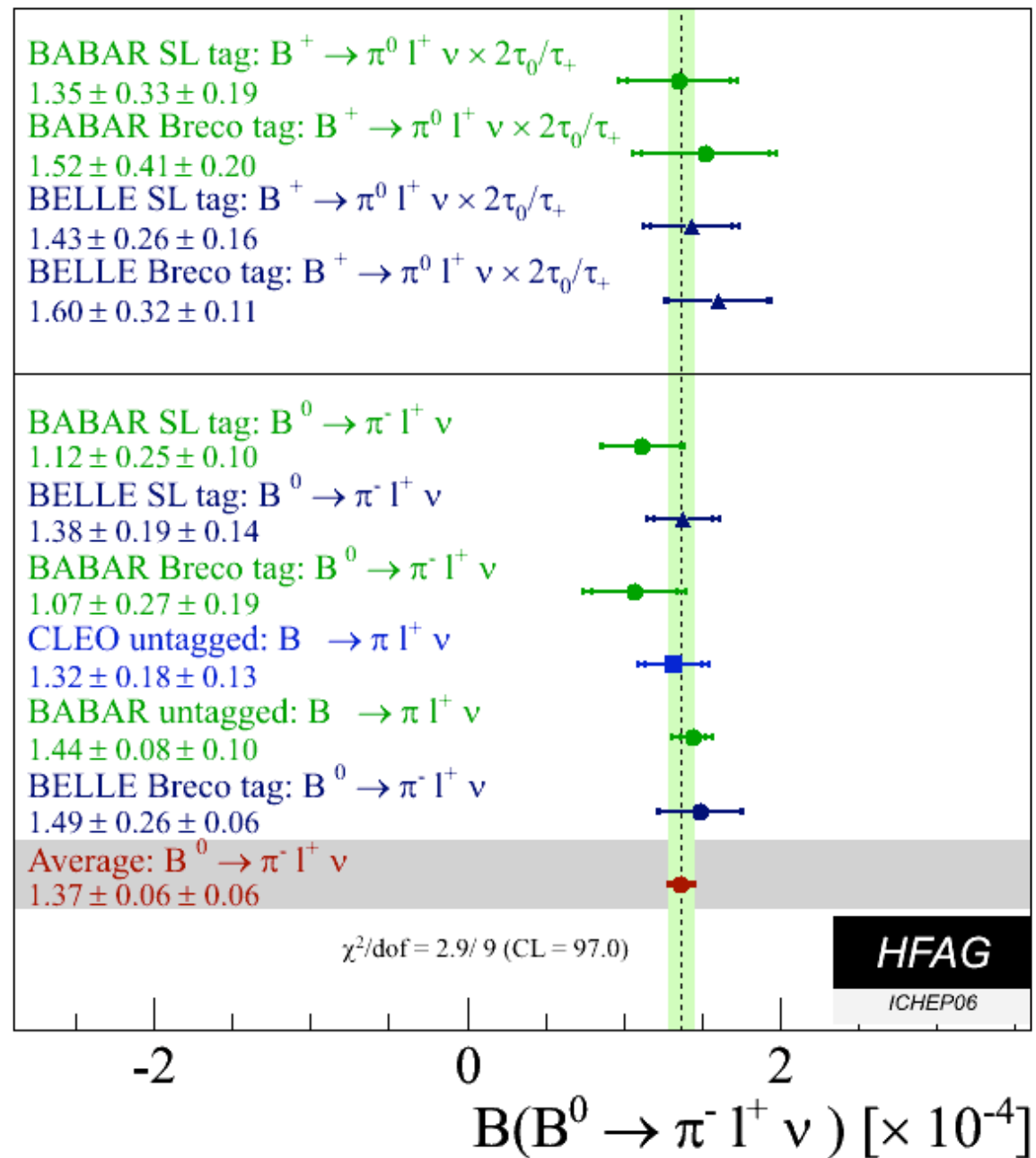


$$|V_{ub}| = (4.1 \pm 0.2 \pm 0.2^{+0.6}_{-0.4}) \times 10^{-3}$$

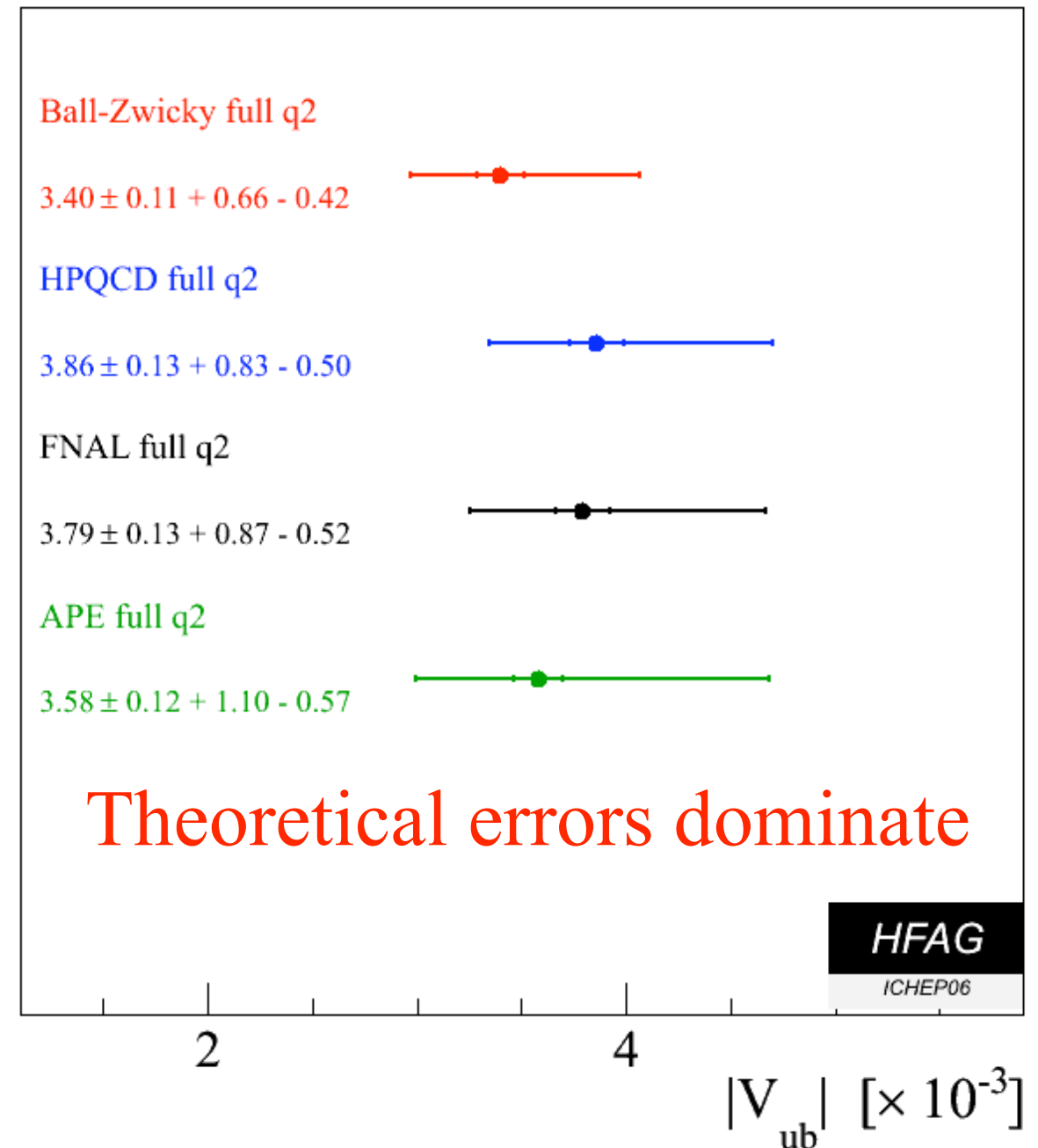
on $q^2 > 16$; by HPQCD

V_{ub} exclusive summary

Measured branching fractions (all 2006)

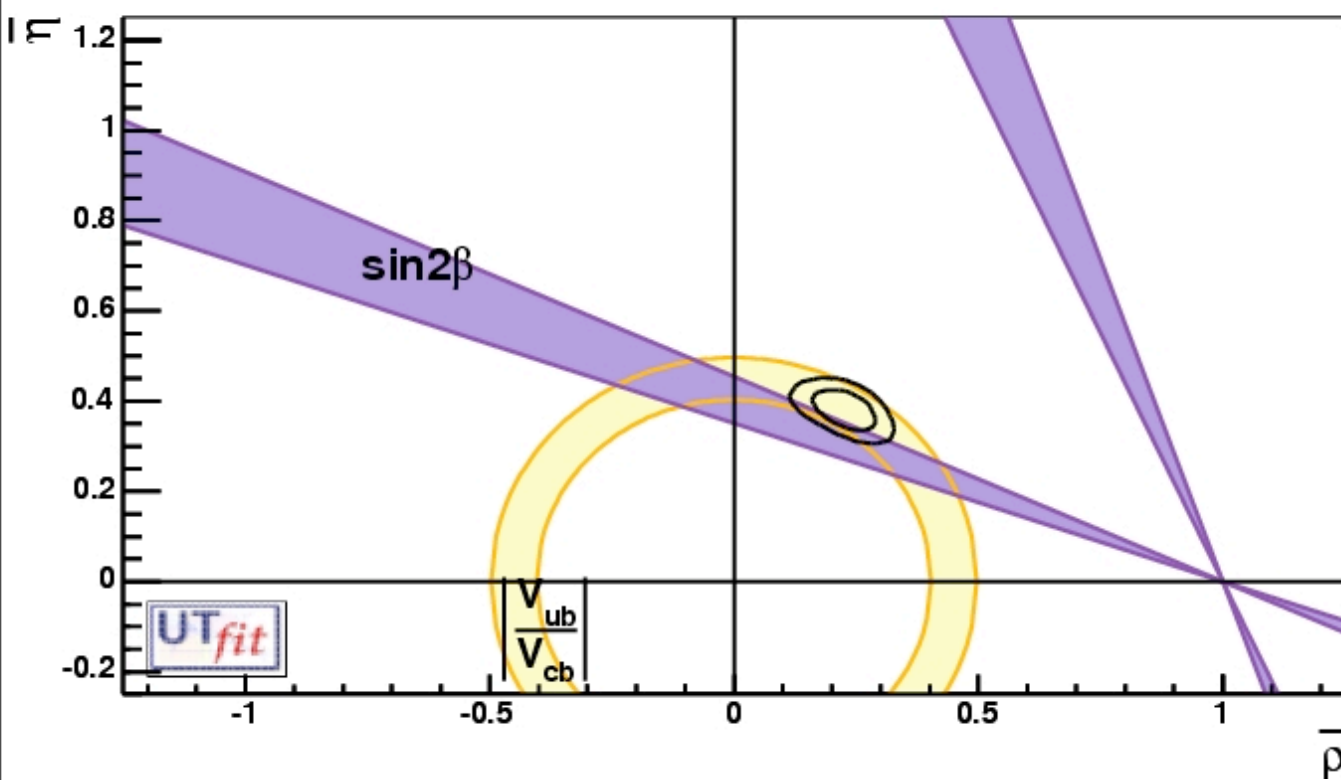


Using theoretical form factors



Experiments starting to measure form factor shape from data; allows elimination of some theory models

V_{ub} vs. $\sin 2\phi_1$

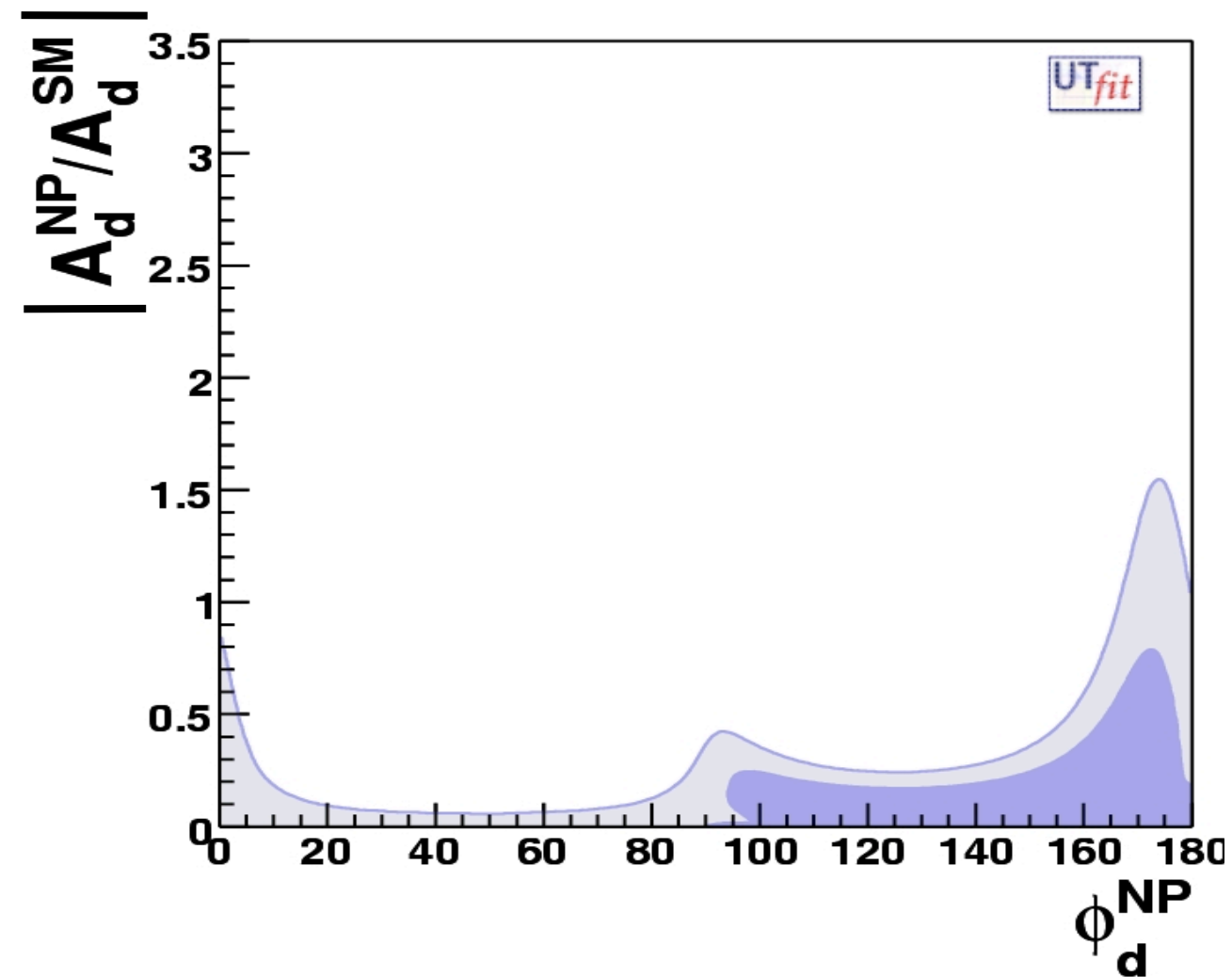


Direct: $\sin 2\phi_1 = 0.67 \pm 0.03$
 Indirect: $\sin 2\phi_1 = 0.76 \pm 0.04$
 Difference: $= 0.09 \pm 0.05$
 Not statistically significant, but...

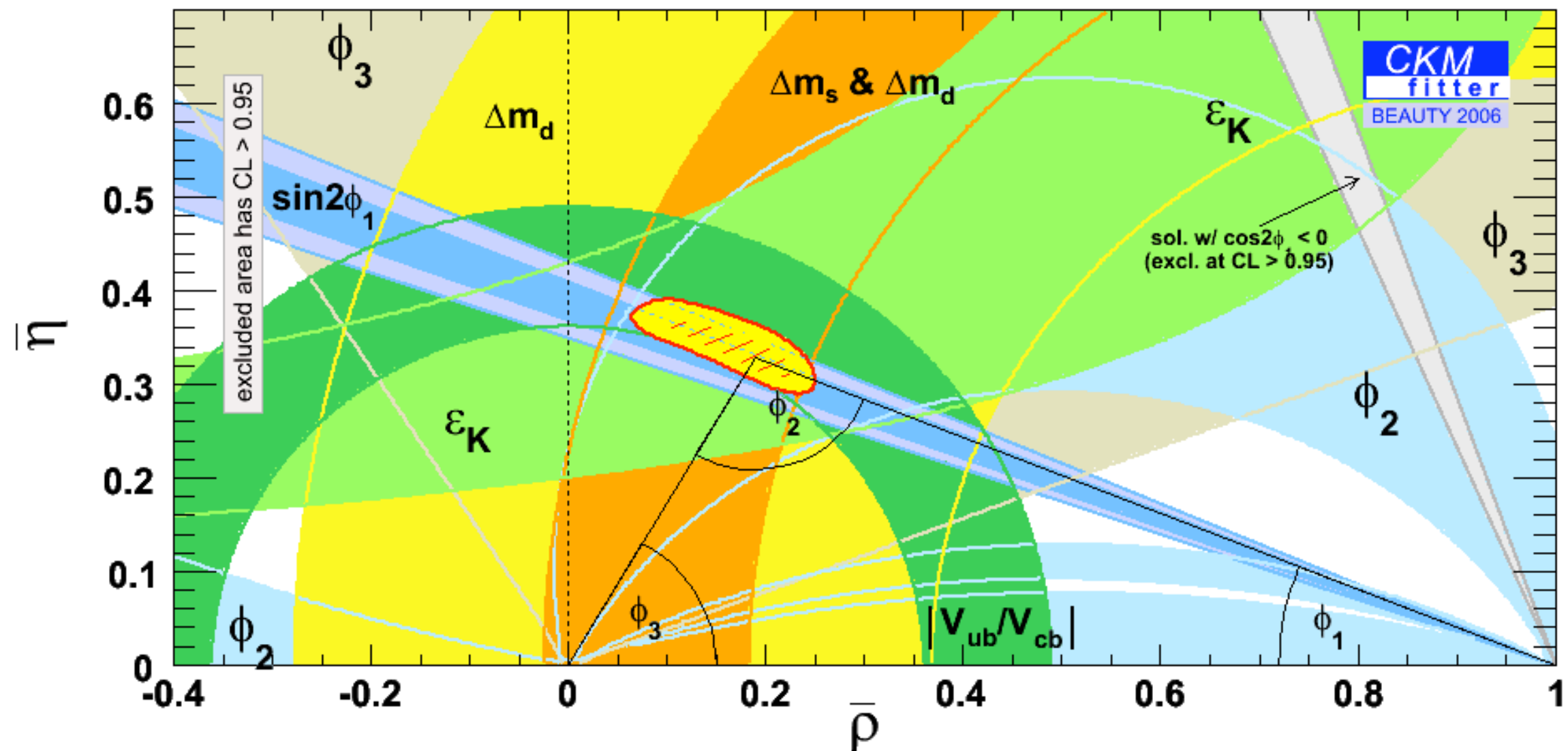
Model independent NP in B mixing
 Add new amplitude to SM

$$A_d = A_d^{\text{SM}} \left[1 + \left| A_d^{\text{NP}} / A_d^{\text{SM}} \right| e^{i2\phi_d^{\text{NP}}} \right]$$

→ modifies ϕ_1 to $\phi_1 + \phi_d^{\text{NP}}$



Summary



Good overall agreement. $O(1)$ new physics unlikely.

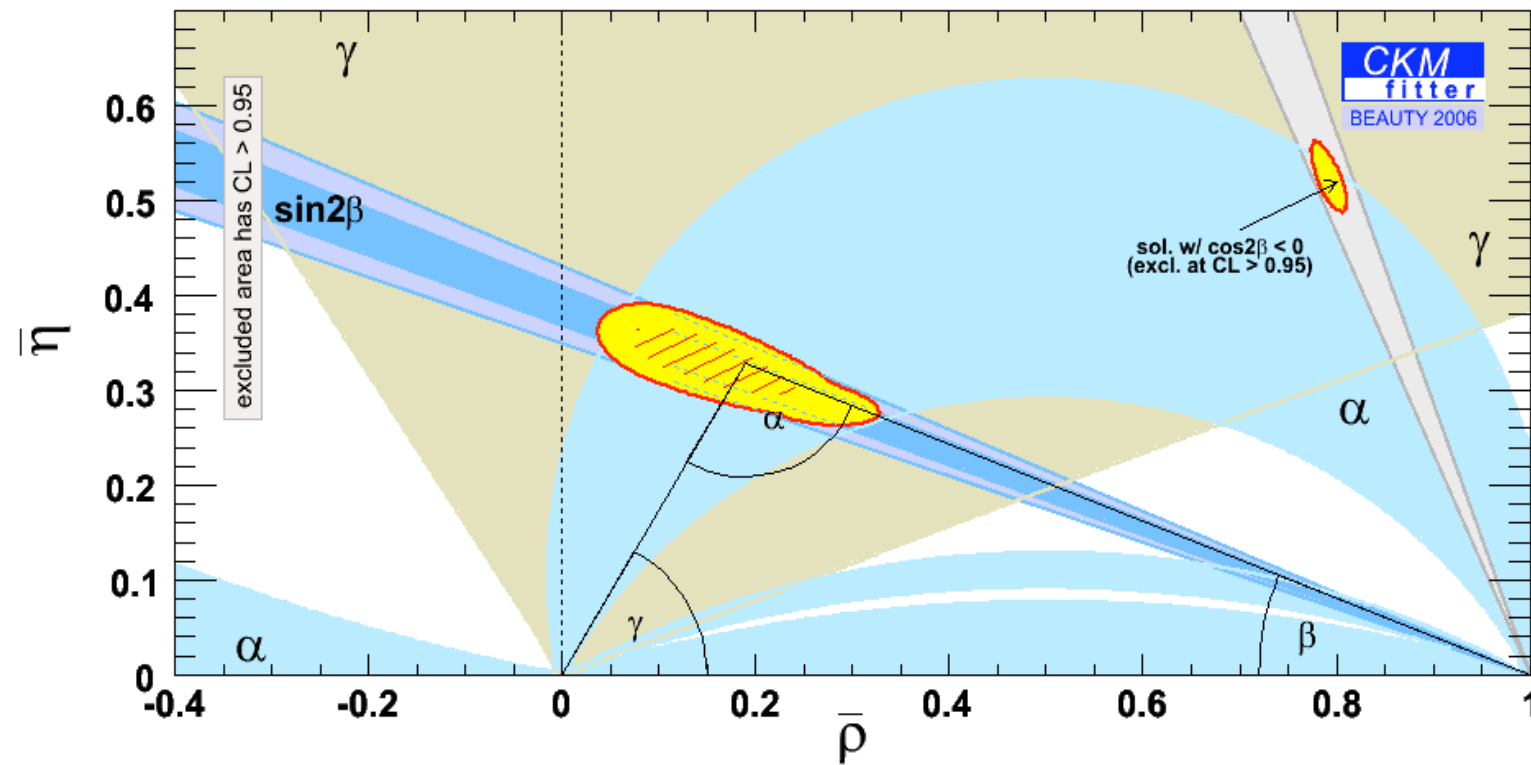
If \exists (new physics @ TeV), deviation might be $\sim [(M_{\text{top}}/M_{\text{NP}})^2, (M_{\text{top}}/M_{\text{NP}})]$,

Prospects

- e^+e^- B-factories are only half-way
 - each aiming at $\sim 1\text{ab}^{-1}$
- Many measurements are statistics-limited
 - $\sin 2\phi_1: b \rightarrow c\bar{c}s$ vs. $b \rightarrow s\bar{s}s$
 - CKM test: all-angles vs. all-sides
- V_{ub} is largely theory-limited; *but*
 - HQE parameters can be better determined experimentally by measuring moments of $B \rightarrow X_c\ell\nu$, $B \rightarrow X_s\gamma$
- More data also brings new decay modes & improved analysis
 - resulting in better than $1/\sqrt{N}$
- LHC, Super-B, ...

All-angles vs. All-sides

Angles



Sides

